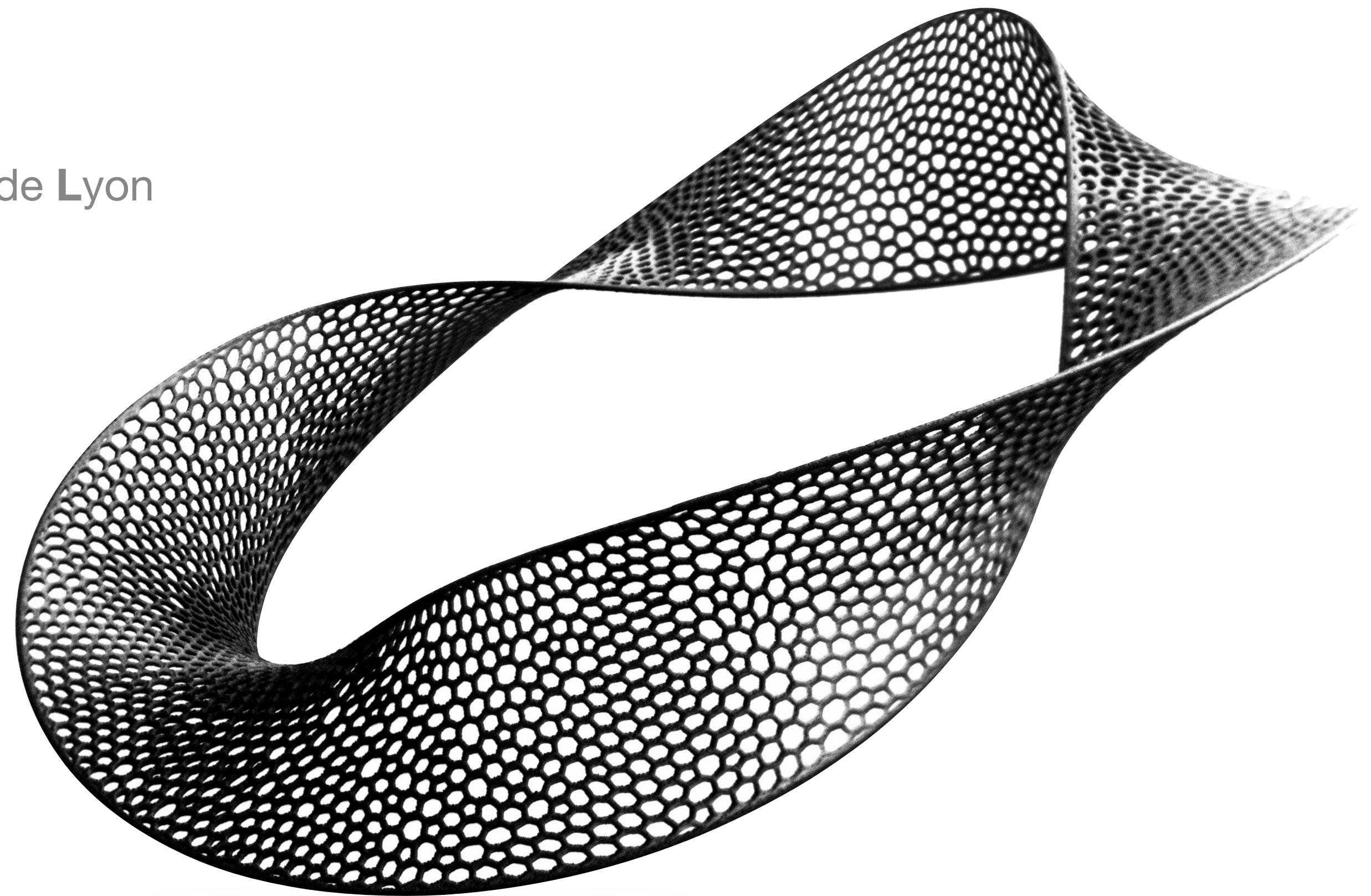


Twisted Matter

David Carpentier

CNRS and Ecole Normale Supérieure de Lyon



Properties of Matter

mechanical properties



Liquid



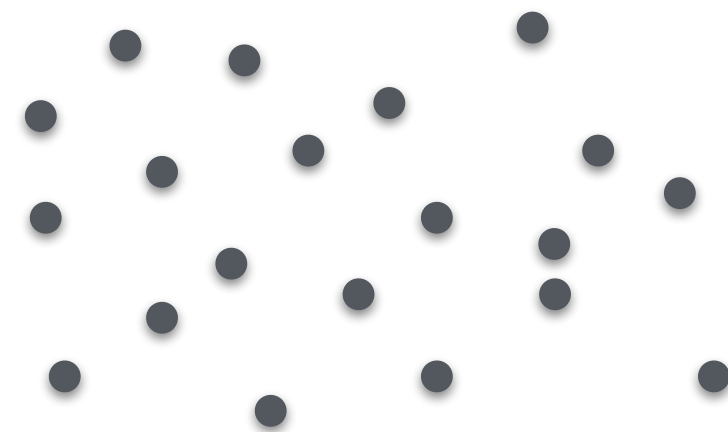
Solid

Properties of Matter

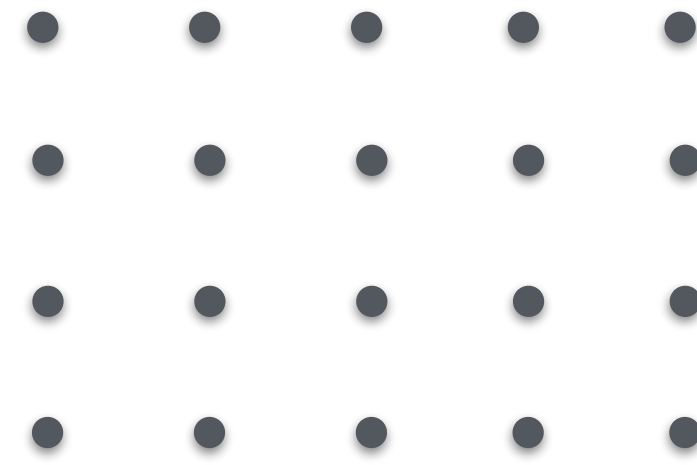
mechanical properties



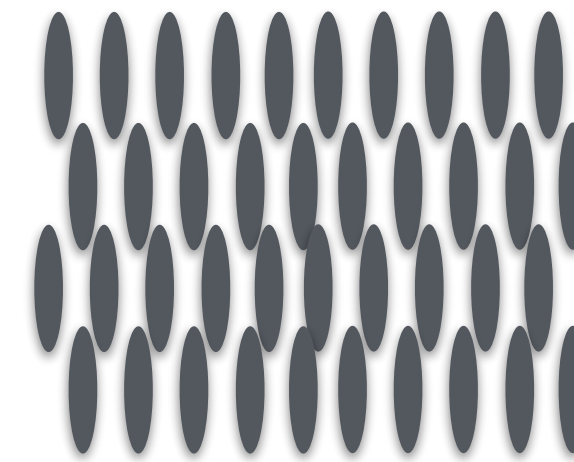
Liquid



Solid



nematic



Properties of Matter

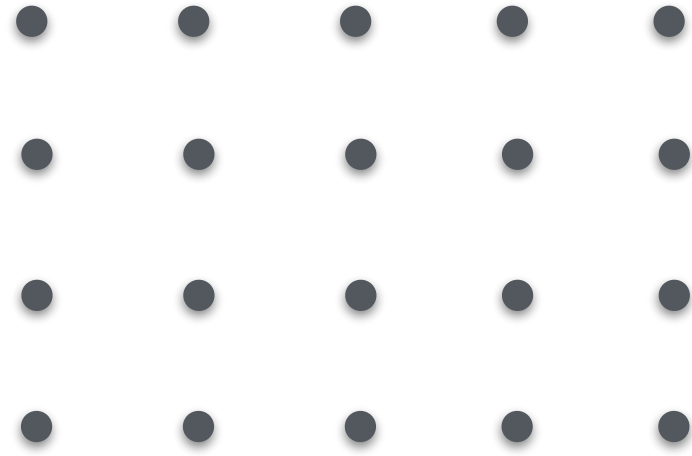
mechanical properties



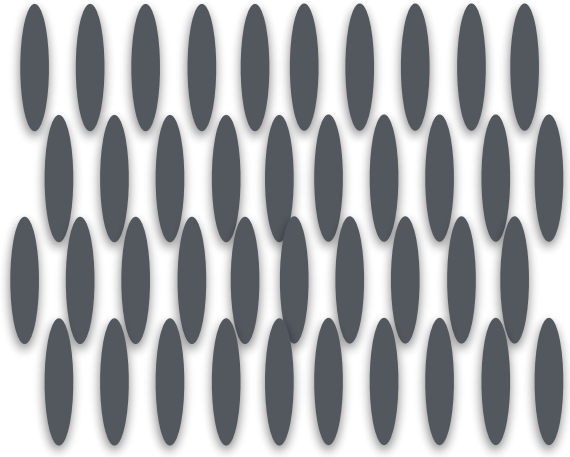
Liquid



Solid

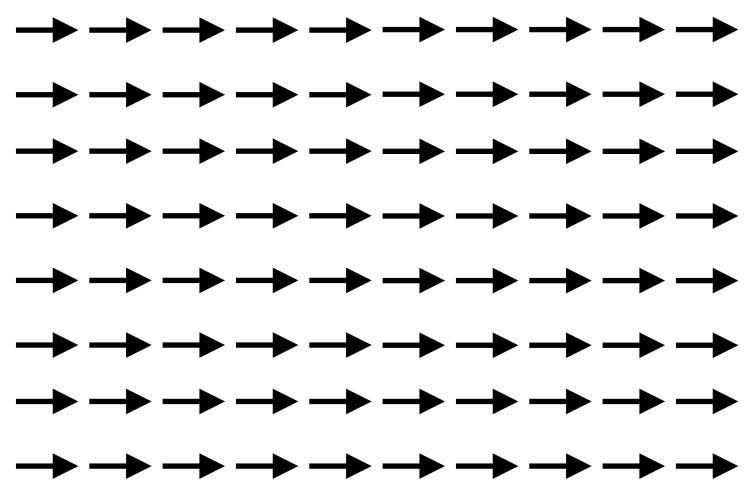


nematic

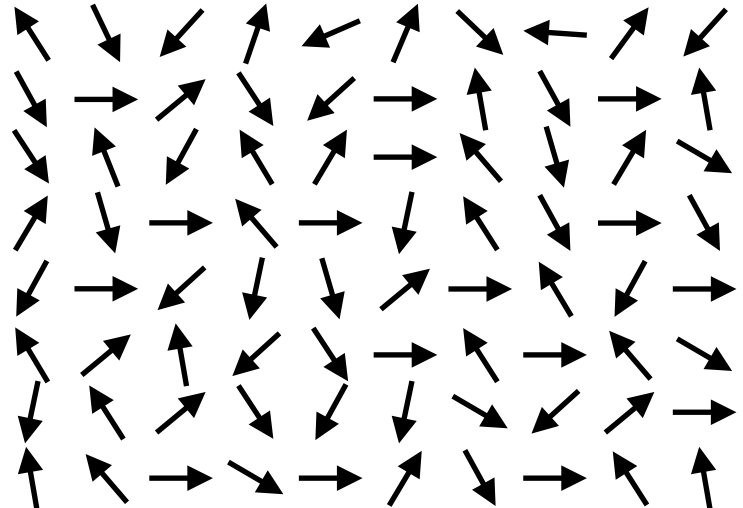


magnetic properties

Ferromagnet



Paramagnet



electronic properties

Metals



Insulators

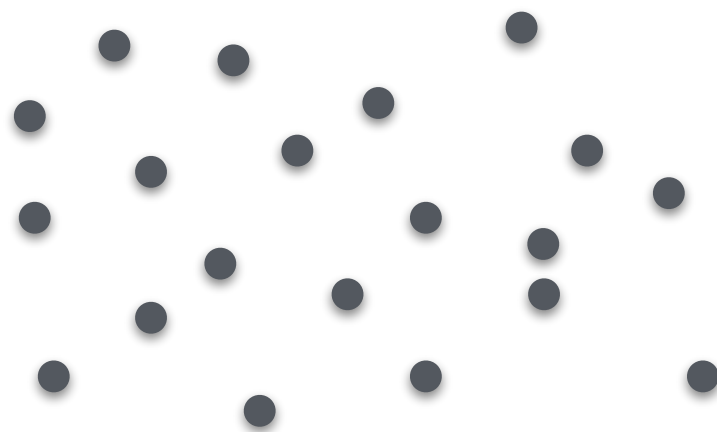


Properties of Matter

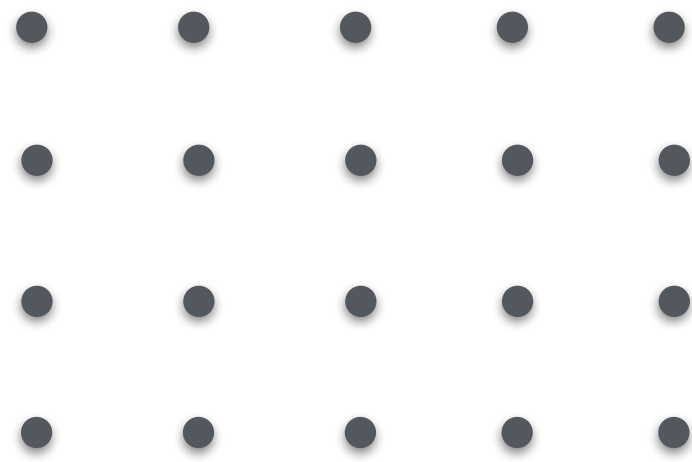
mechanical properties



Liquid

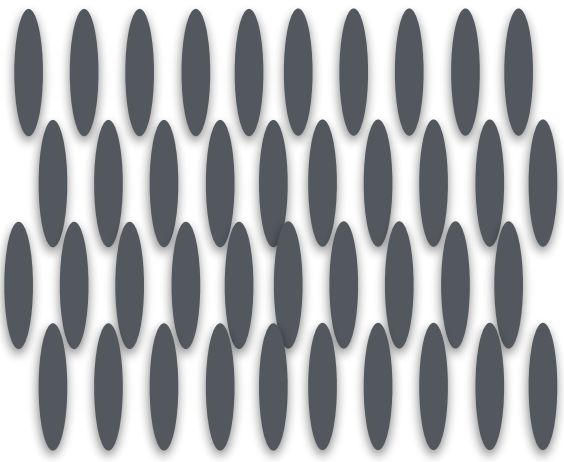


Solid



Symmetry : leaves the phase invariant

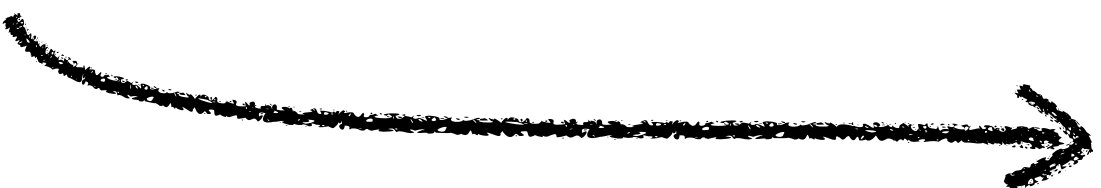
nematic



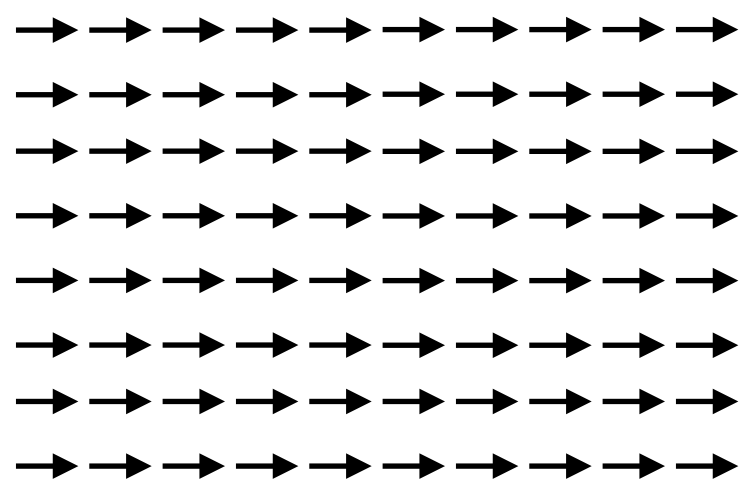
magnetic properties



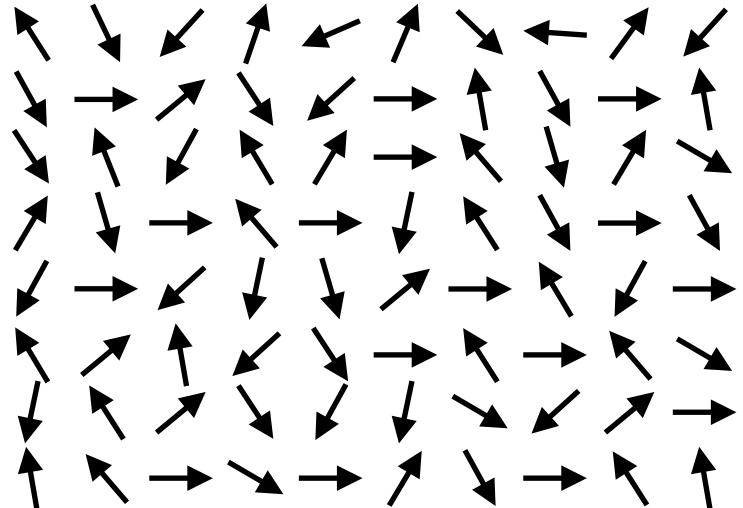
electronic properties



Ferromagnet



Paramagnet



Metals

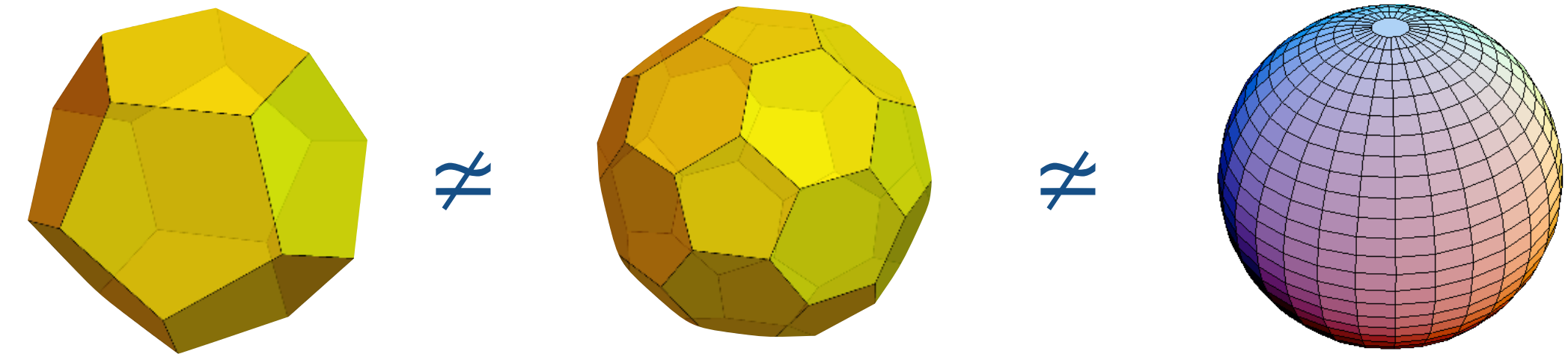


Insulators

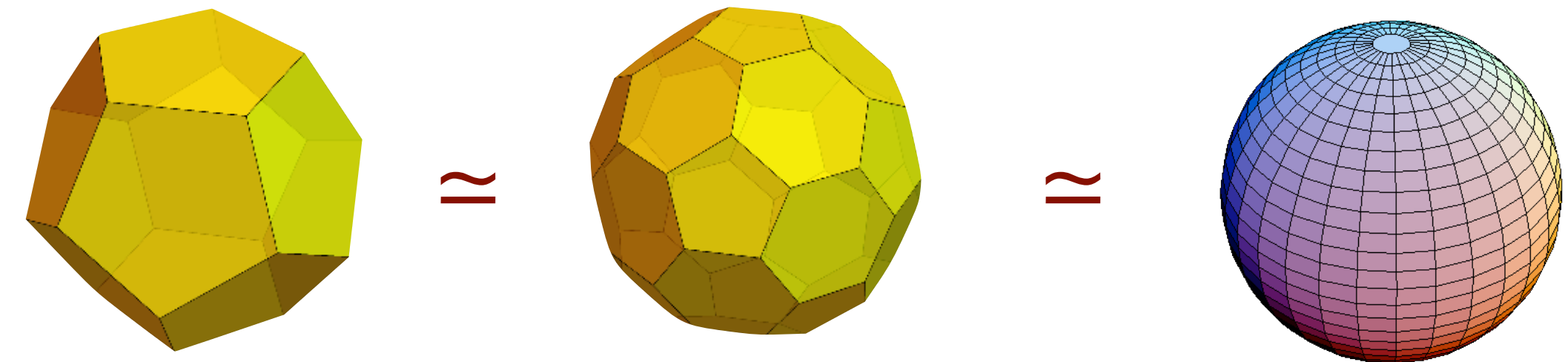


Topology versus Geometry

Symmetry : Leaves the
physikalisch invariant

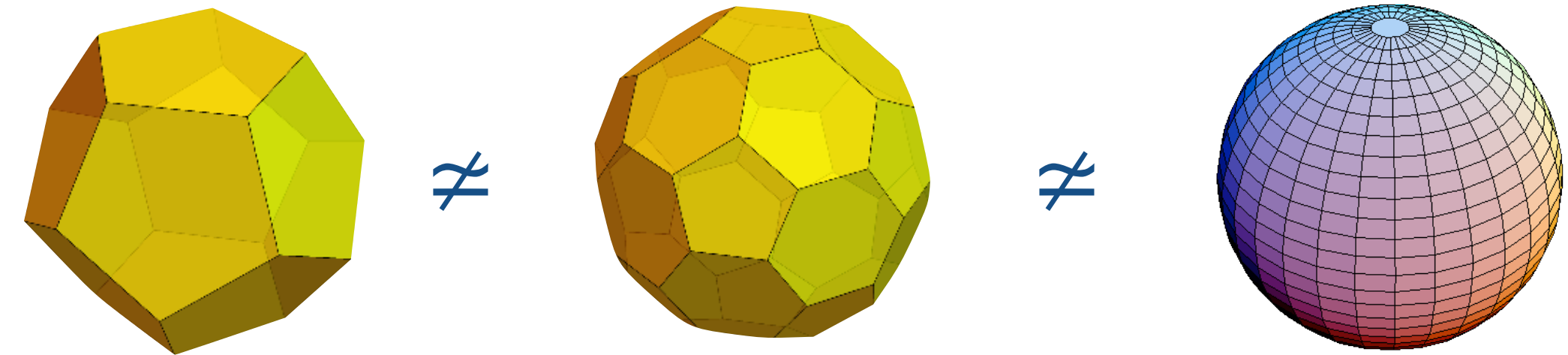


Topology : global shape
the study of properties unaffected
by the continuous change of shape
or size

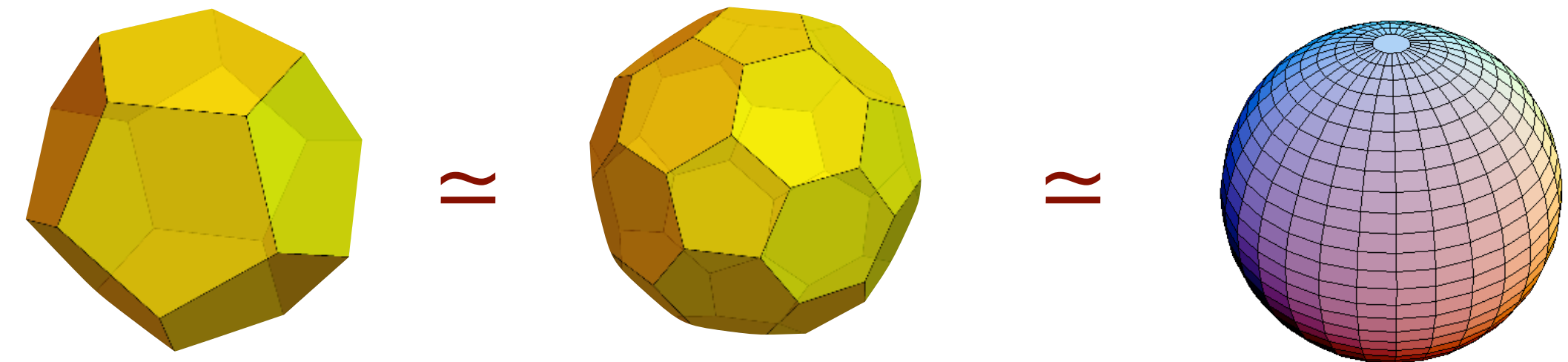


Topology versus Geometry

Symmetry : Leaves the object invariant



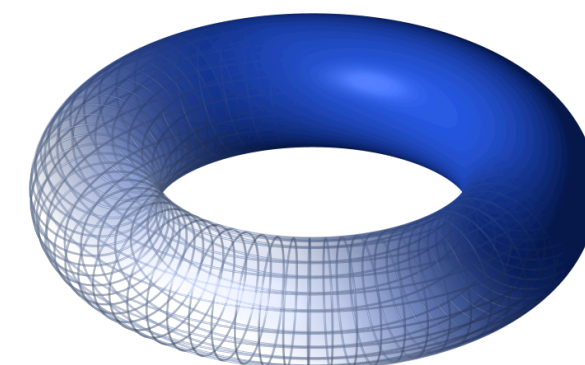
Topology : global shape
the study of properties unaffected by the continuous change of shape or size



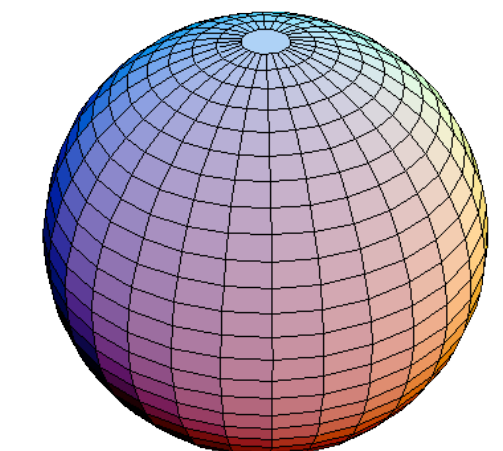
\cong



\cong



\neq



the famous «donut »

Twisted Matter : characterization via topology

Topology : global shape

the study of properties unaffected
by the continuous change of shape
or size

Electronic Properties of
Quantum Matter

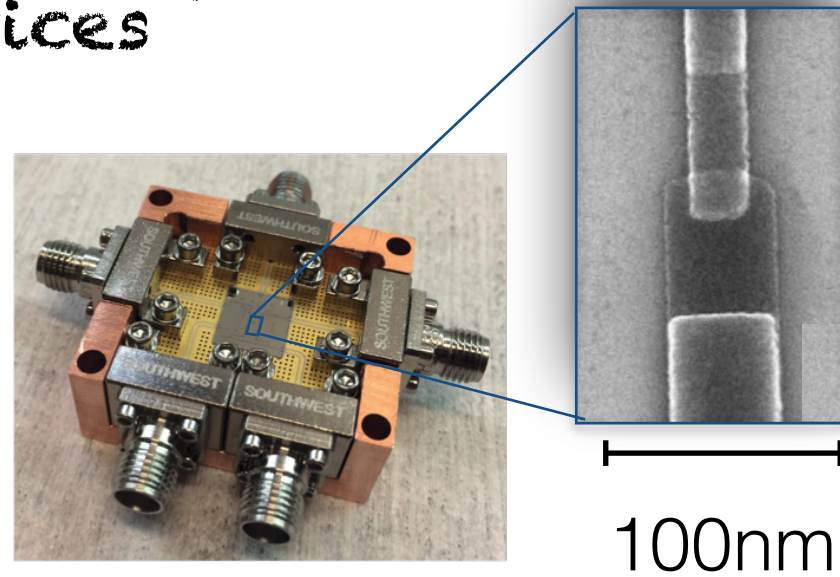
Topological Insulators



Twisted Matter : characterization via topology

Quantum Technologies

Robust quantum devices



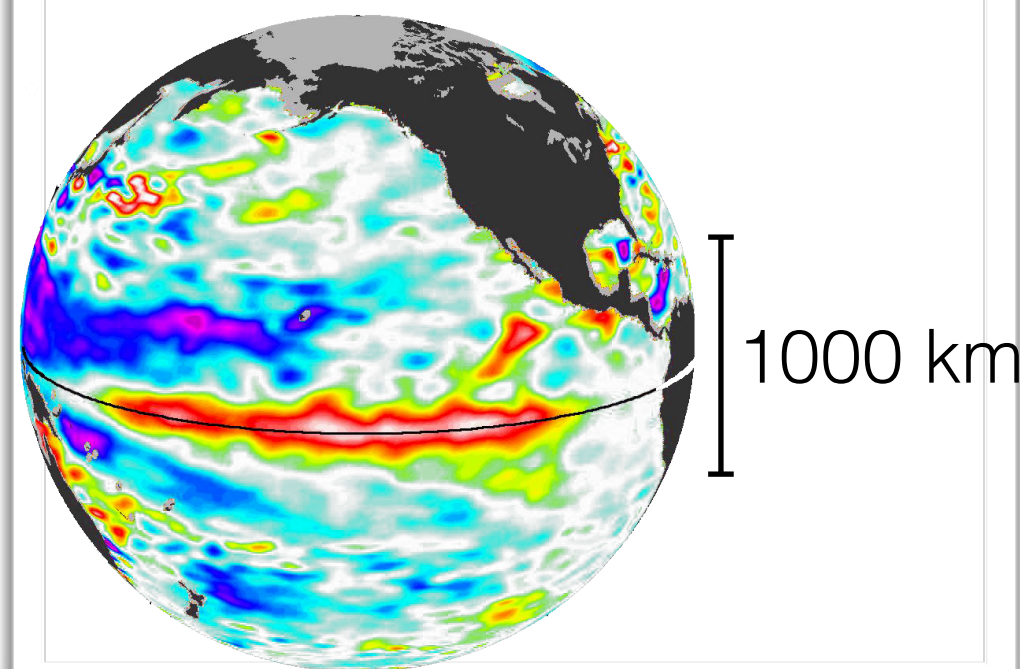
Electronic Properties of Quantum Matter

Topological Insulators

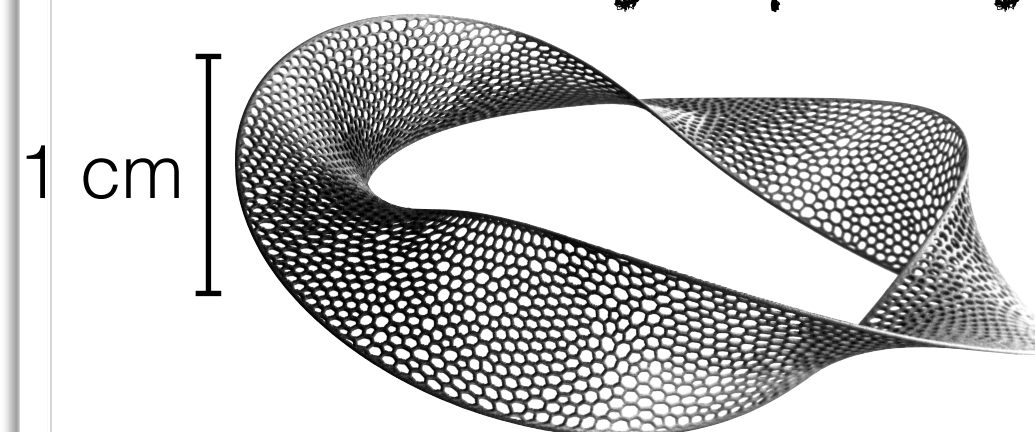


Geofluids

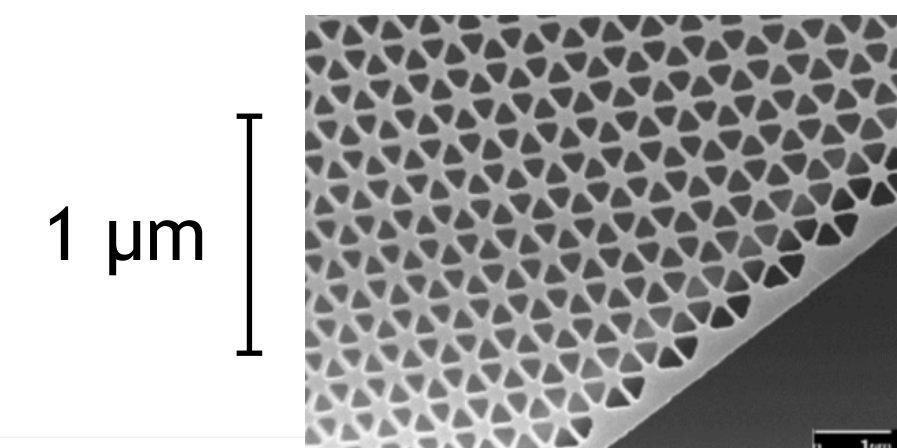
equatorial waves of topological origin



Mechanics / Metamaterials deformations constrained by topology



Optics / Metamaterials optical modes constrained by topology



Topology
global shape

Topological Matter in Lyon

Quantum Technologies



D. Carpentier,
P. Delplace

B. Huard

Electronic and Magnetic Properties of Matter



A. Fedorenko



P. Holdsworth



T. Rolscilde



L. Savary



E. Orignac

Geofluids



P. Delplace



A. Venaille

Topology
global shape



K. Gawedzki



J. Kellendonk



J.-M. Stephan

(Institute Camille Jourdan)

Mechanics / Metamaterials



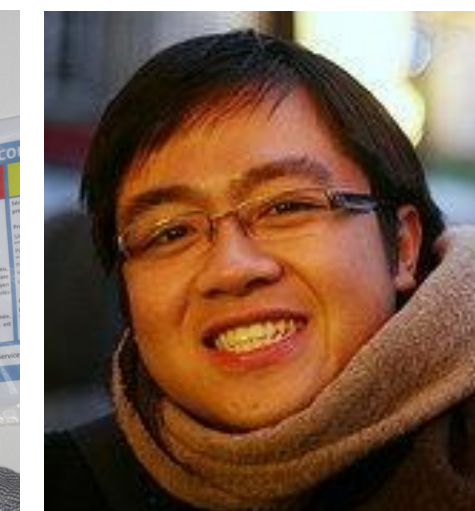
D. Bartolo

D. Carpentier

Optics / Metamaterials



L. Ferrier



H.-S. Nguyen

D. Carpentier, P.
Delplace

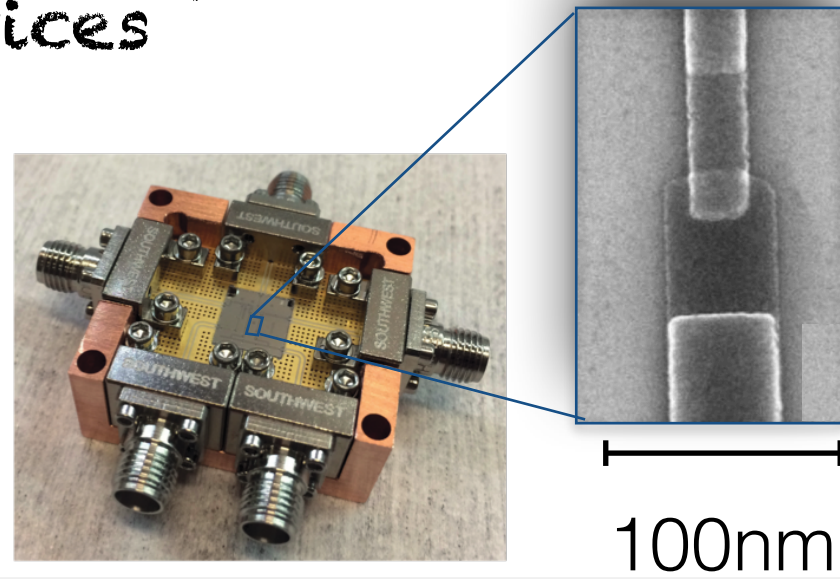
(Lyon Institute of Nanotechnology)

ToRe Breakthrough Project



Quantum Technologies

Robust quantum devices



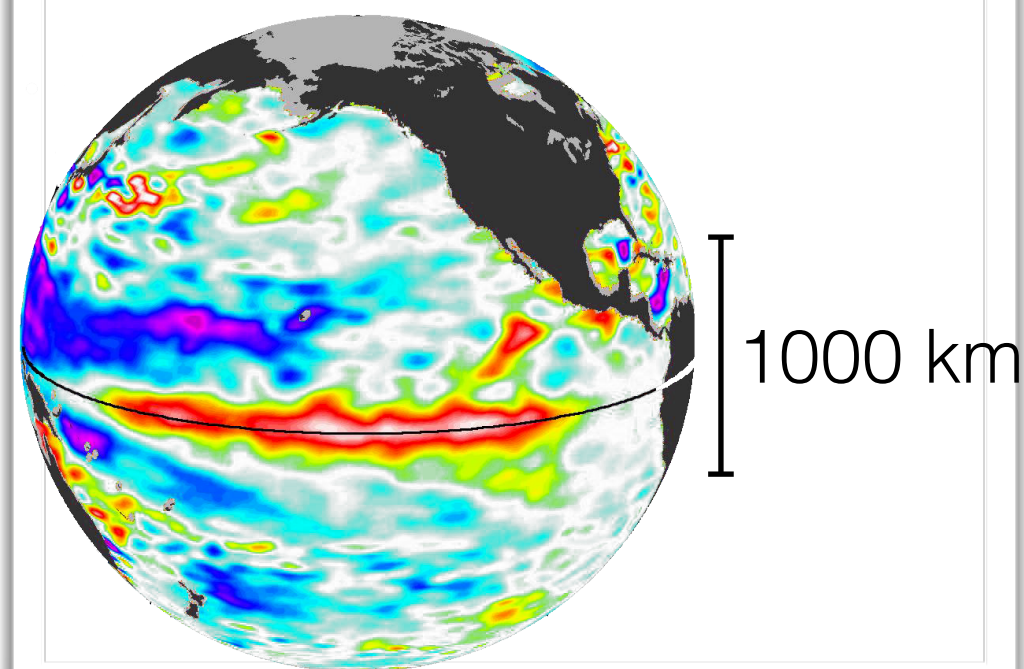
Electronic Properties of Quantum Matter

Topological Insulators

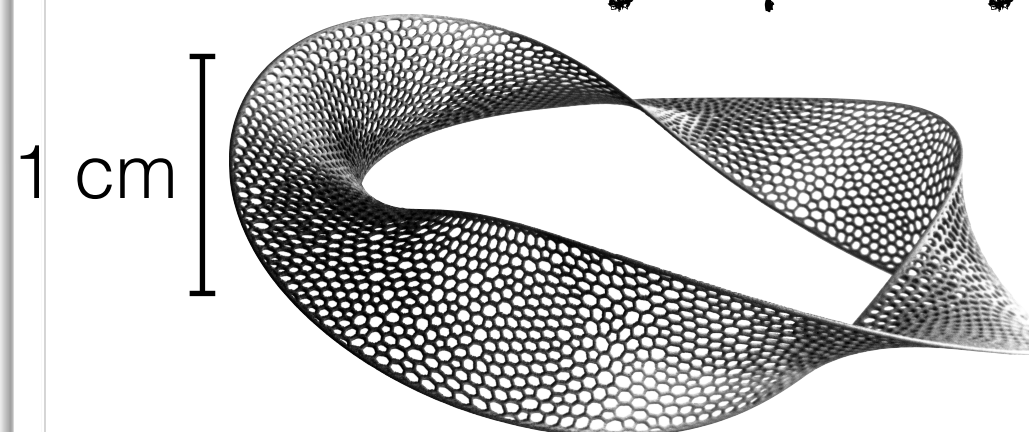


Geofluids

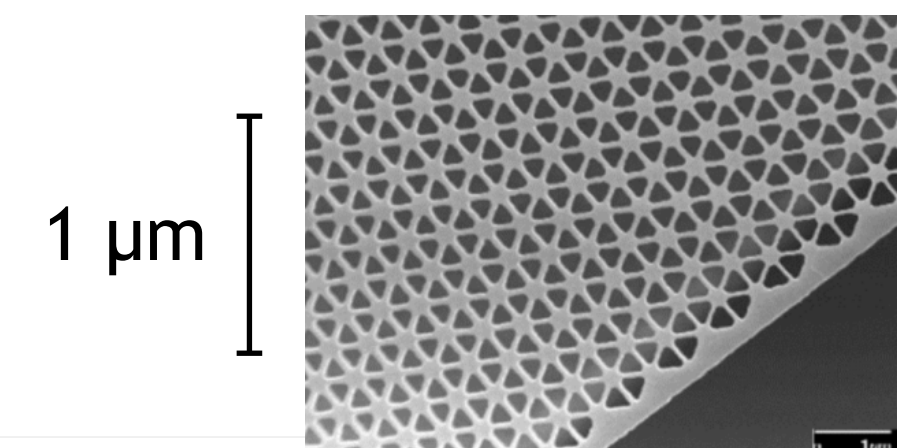
equatorial waves of topological origin



Mechanics / Metamaterials
deformations
constrained by topology



Optics / Metamaterials
optical modes
constrained by topology



Topology
global shape

Outline

1. Electronic Properties of Quantum Matter

Topological Insulators



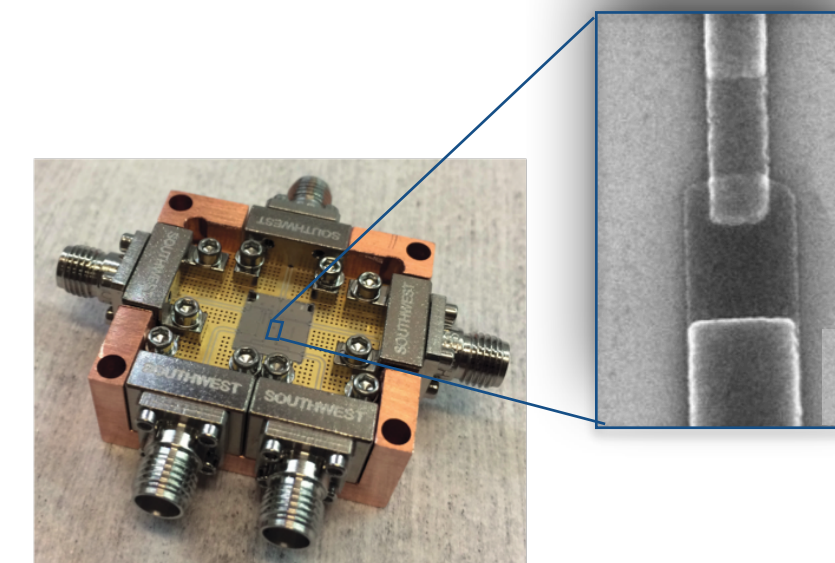
Metals



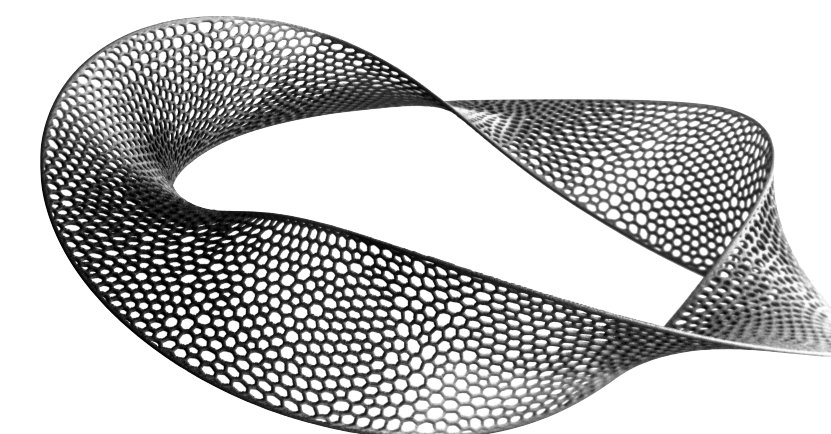
Insulators

2. Quantum Technologies

Robust quantum devices



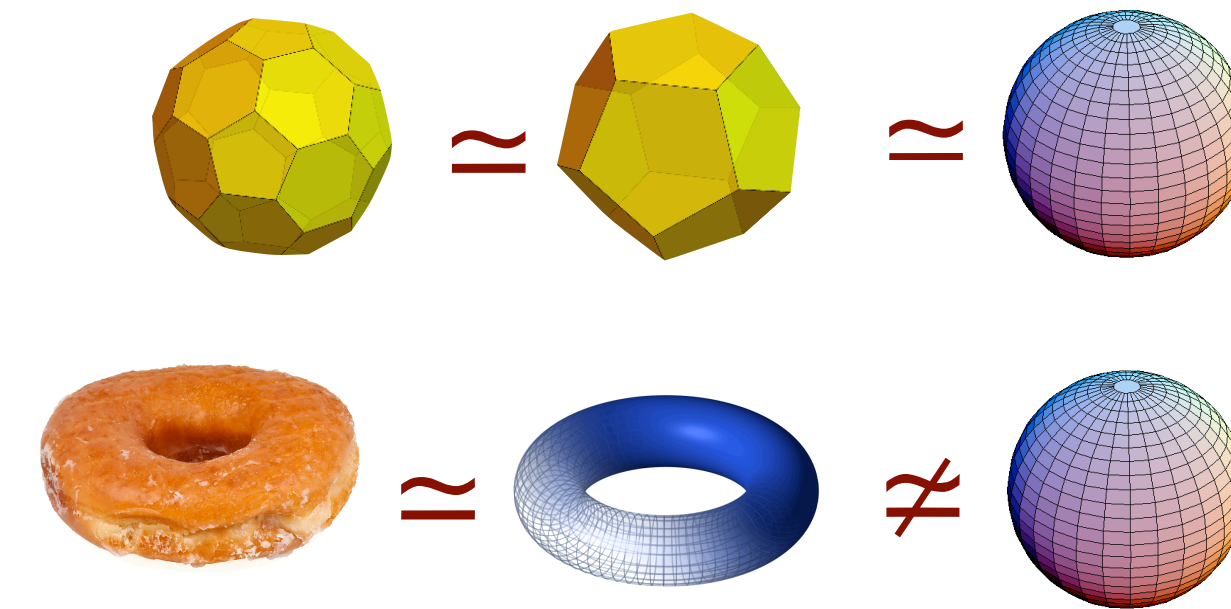
3. Mechanics / Metamaterials deformations constrained by topology



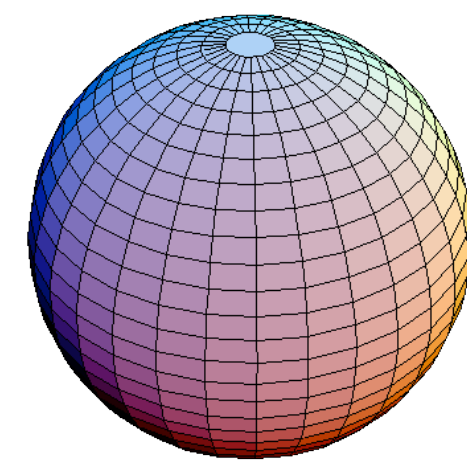
Topology

Topology: aims at classifying objects

- ▶ identifies properties of objects that are preserved under continuous deformations
- ▶ uses **integer number** to distinguish classes of objects



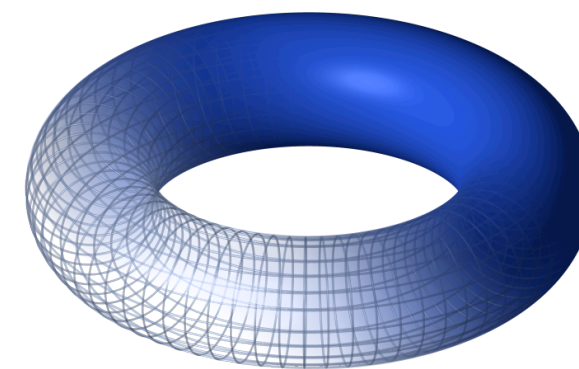
Example of 2d surfaces :



$$\chi = 2$$

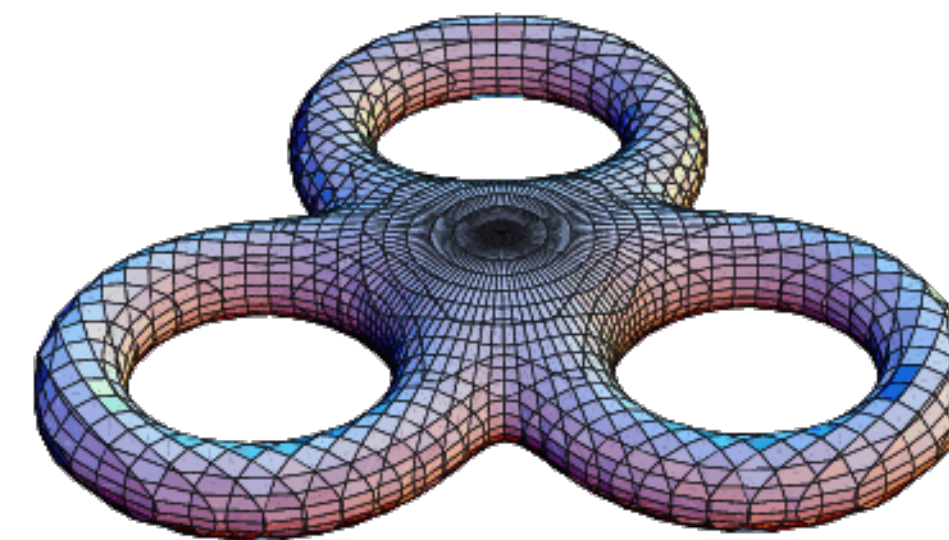
(Euler characteristic)

≠



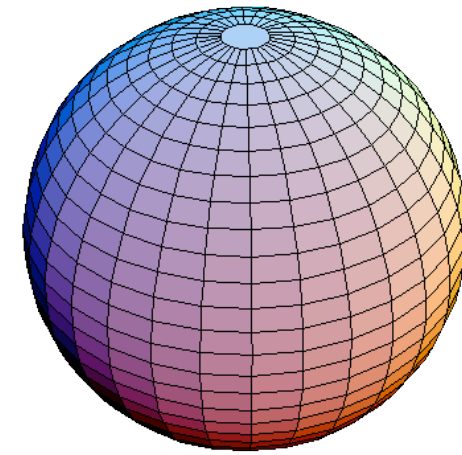
$$\chi = 0$$

≠



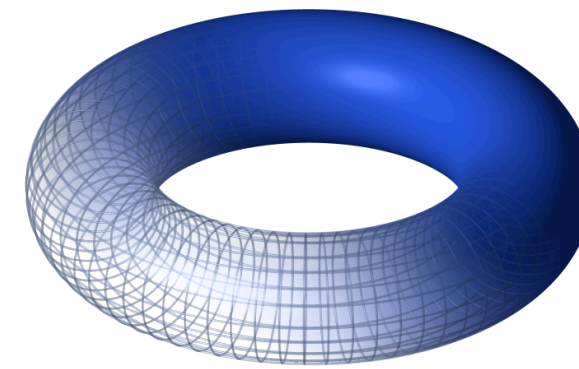
$$\chi = -4$$

Topology



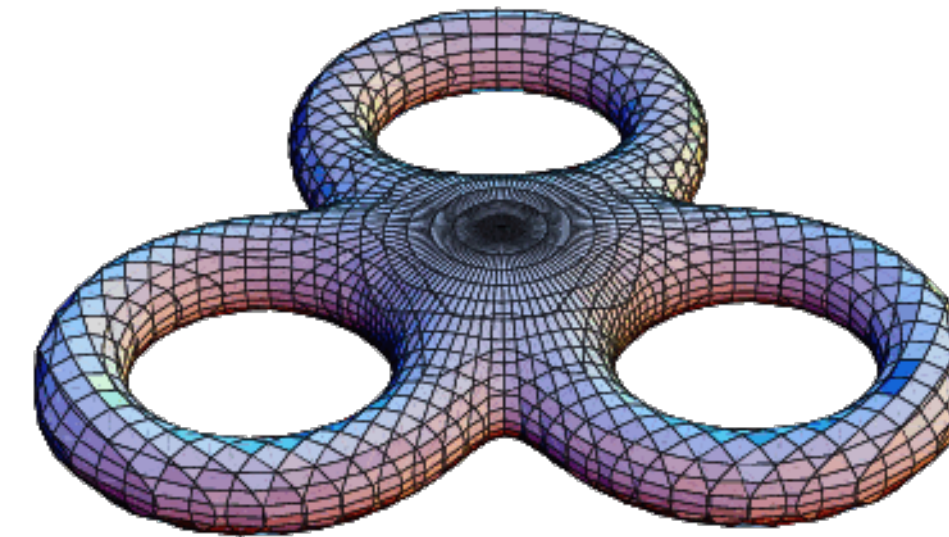
$$\chi = 2$$

≠



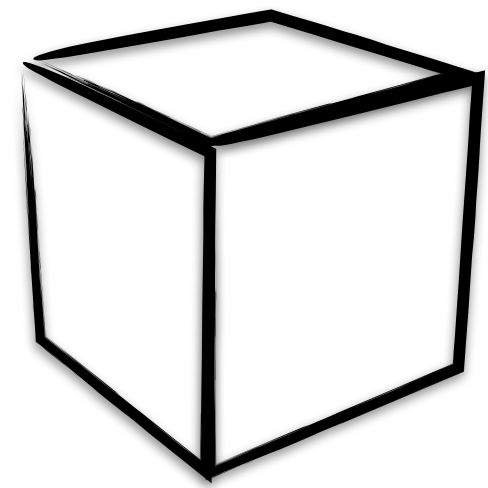
$$\chi = 0$$

≠

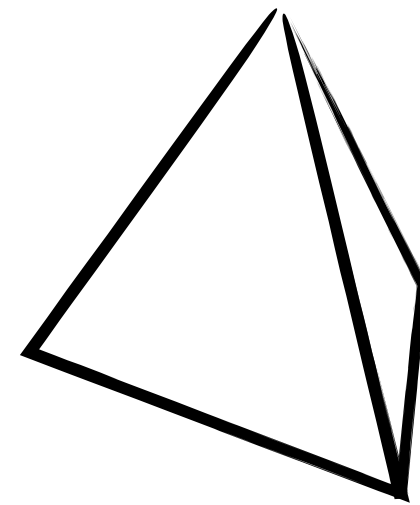


$$\chi = -4$$

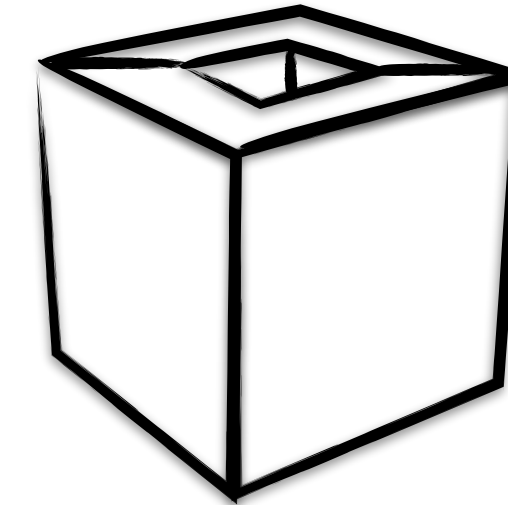
► For polygons: **Euler characteristic** $\chi = \# \text{vertices} - \# \text{edges} + \# \text{faces}$



$$\chi = 8 - 12 + 6 = +2$$

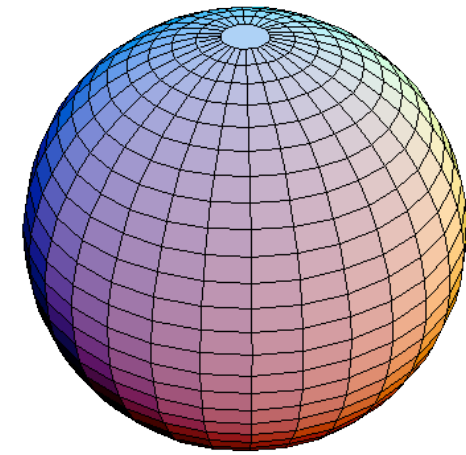


$$\chi = 4 - 6 + 4 = +2$$



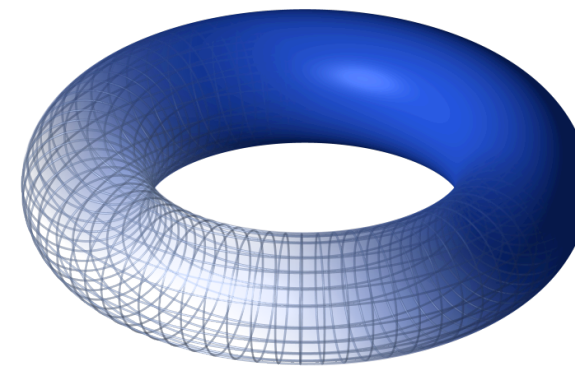
$$\chi = 16 - 28 + 12 = 0$$

Topology



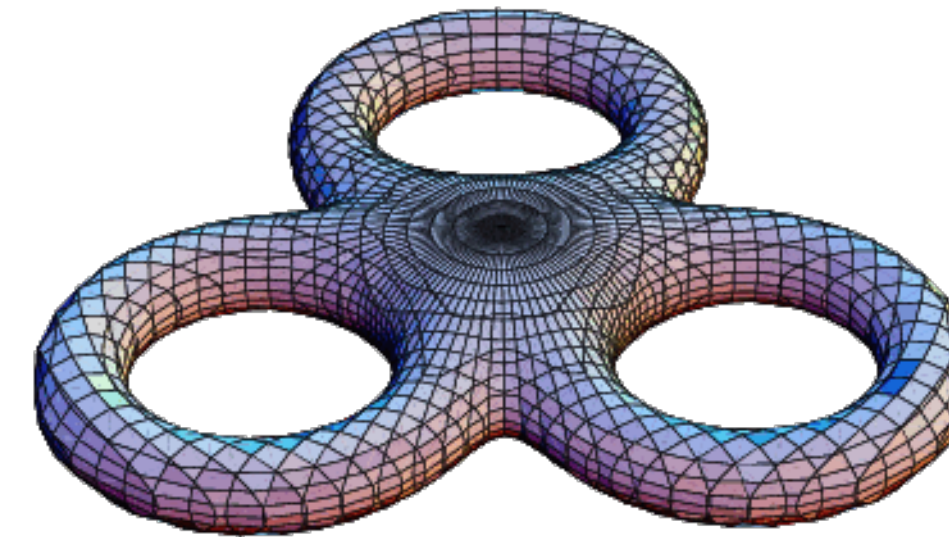
$$\chi = 2, g = 0$$

≠



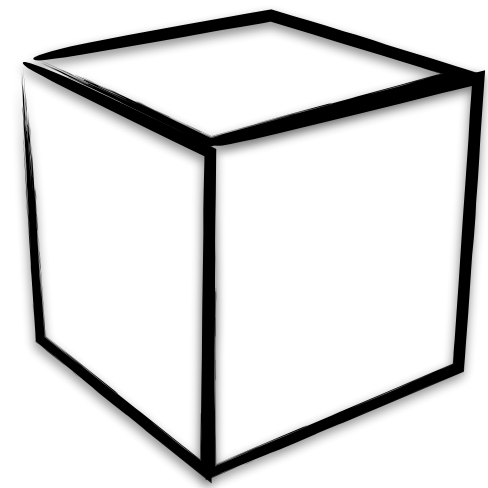
$$\chi = 0, g = 1$$

≠

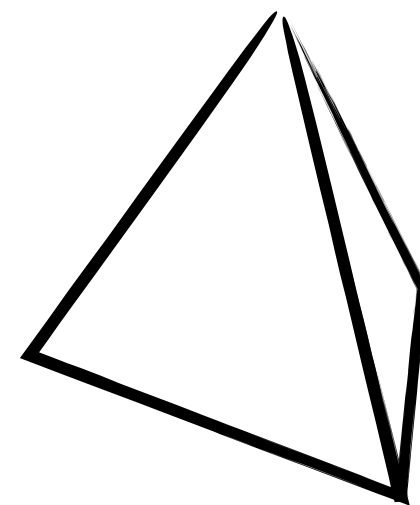


$$\chi = -4, g = 3$$

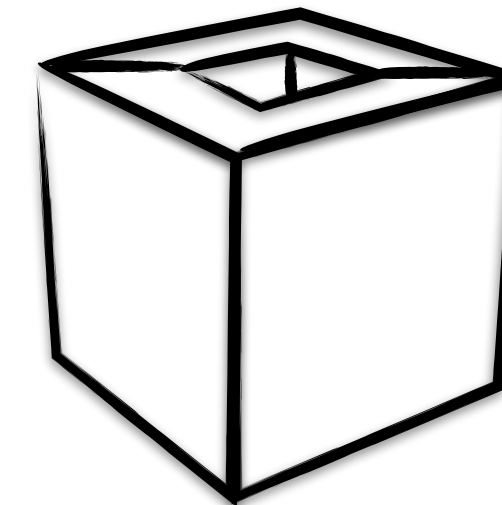
► For polygons: **Euler characteristic** $\chi = \# \text{vertices} - \# \text{edges} + \# \text{faces}$



$$\chi = 8 - 12 + 6 = +2$$



$$\chi = 4 - 6 + 4 = +2$$



$$\chi = 16 - 28 + 12 = 0$$

► **Euler characteristic** \leftrightarrow genus g : $\chi = 2 - 2g$

► Gauss-Bonnet theorem $\chi = \int dS \kappa$

Gaussian curvature : $\kappa = 1/(R_1 R_2)$

- curvature : depends on «local properties»
- Integral of curvature : «global property» (topology)

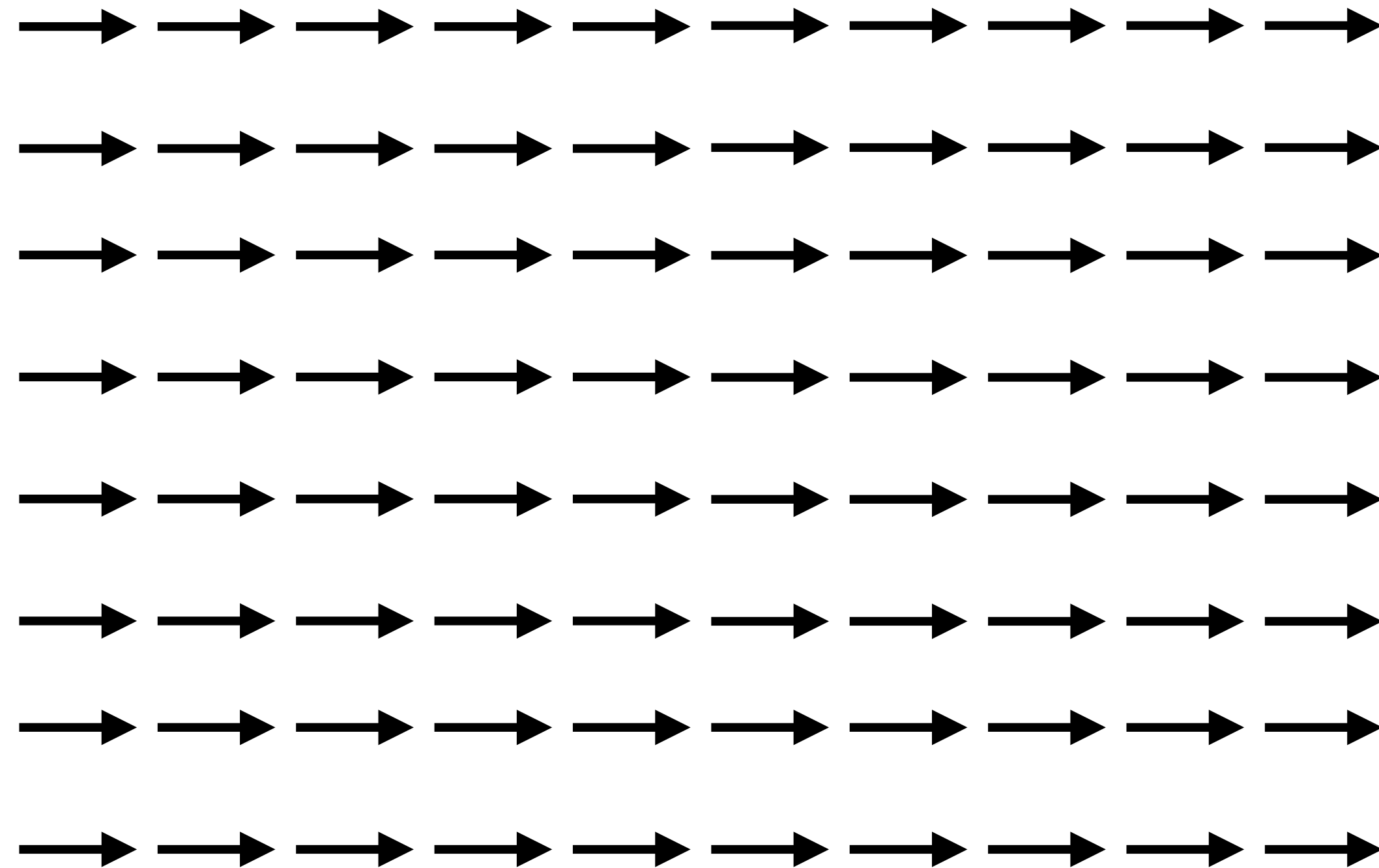
Topology

In condensed matter : **topological defects** of ordered phase / topological **textures**

- ▶ **vortices** (superfluid, superconductor, XY spins),
- ▶ **dislocations and disclinations** (solids, liquid crystals),
- ▶ **hedgehog / skyrmions** (SU(2) spins), etc.

Ordered phase :

- ▶ order parameter $\psi(x) \in \mathbb{C}$
- ▶ spatial order



d=2, complex order parameter

Topology

In condensed matter : **topological defects** of ordered phase / topological **textures**

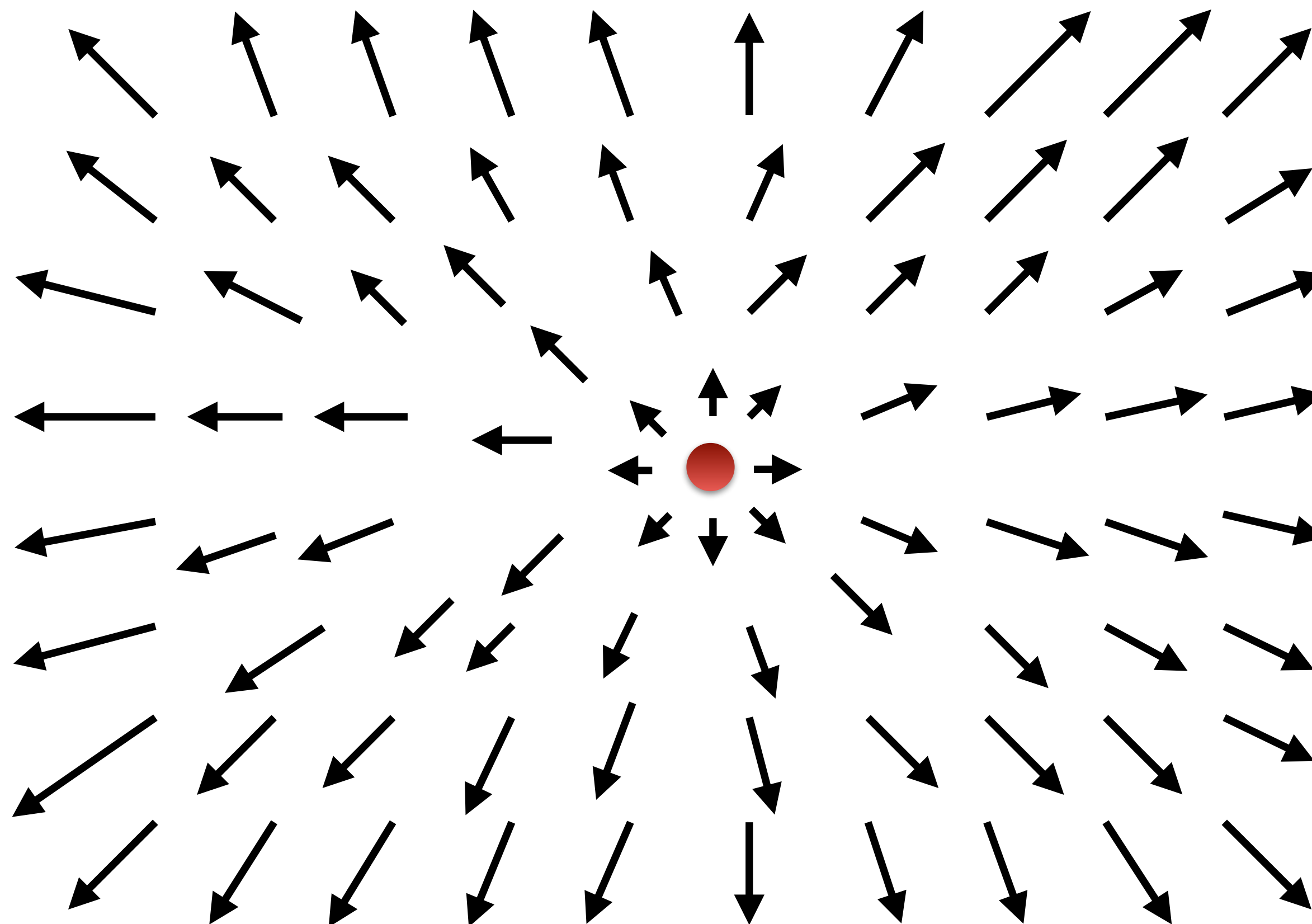
- ▶ **vortices** (superfluid, superconductor, XY spins),
- ▶ **dislocations and disclinations** (solids, liquid crystals),
- ▶ **hedgehog / skyrmions** (SU(2) spins), etc.

Ordered phase :

- ▶ **order parameter**
- ▶ **spatial order**

Associated Defect

- ▶ **singularity** of order field



$d=2$, complex order parameter

Topology

In condensed matter : **topological defects** of ordered phase / topological **textures**

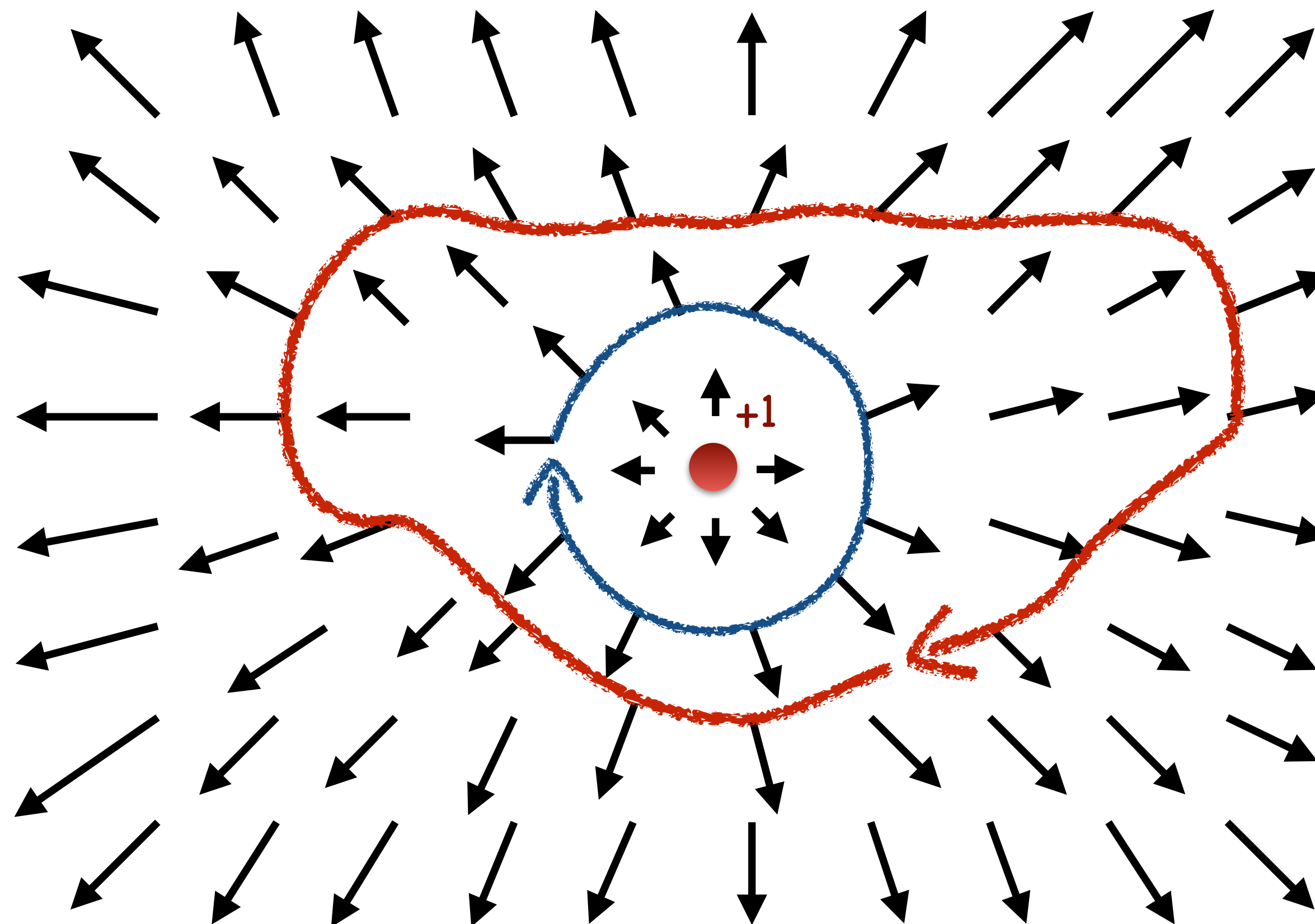
- ▶ **vortices** (superfluid, superconductor, XY spins),
- ▶ **dislocations and disclinations** (solids, liquid crystals),
- ▶ **hedgehog / skyrmions** (SU(2) spins), etc.

Ordered phase :

- ▶ **order parameter**
- ▶ **spatial order**

Associated Defect

- ▶ **singularity** of order field
- ▶ winding of order parameter : **topological number**

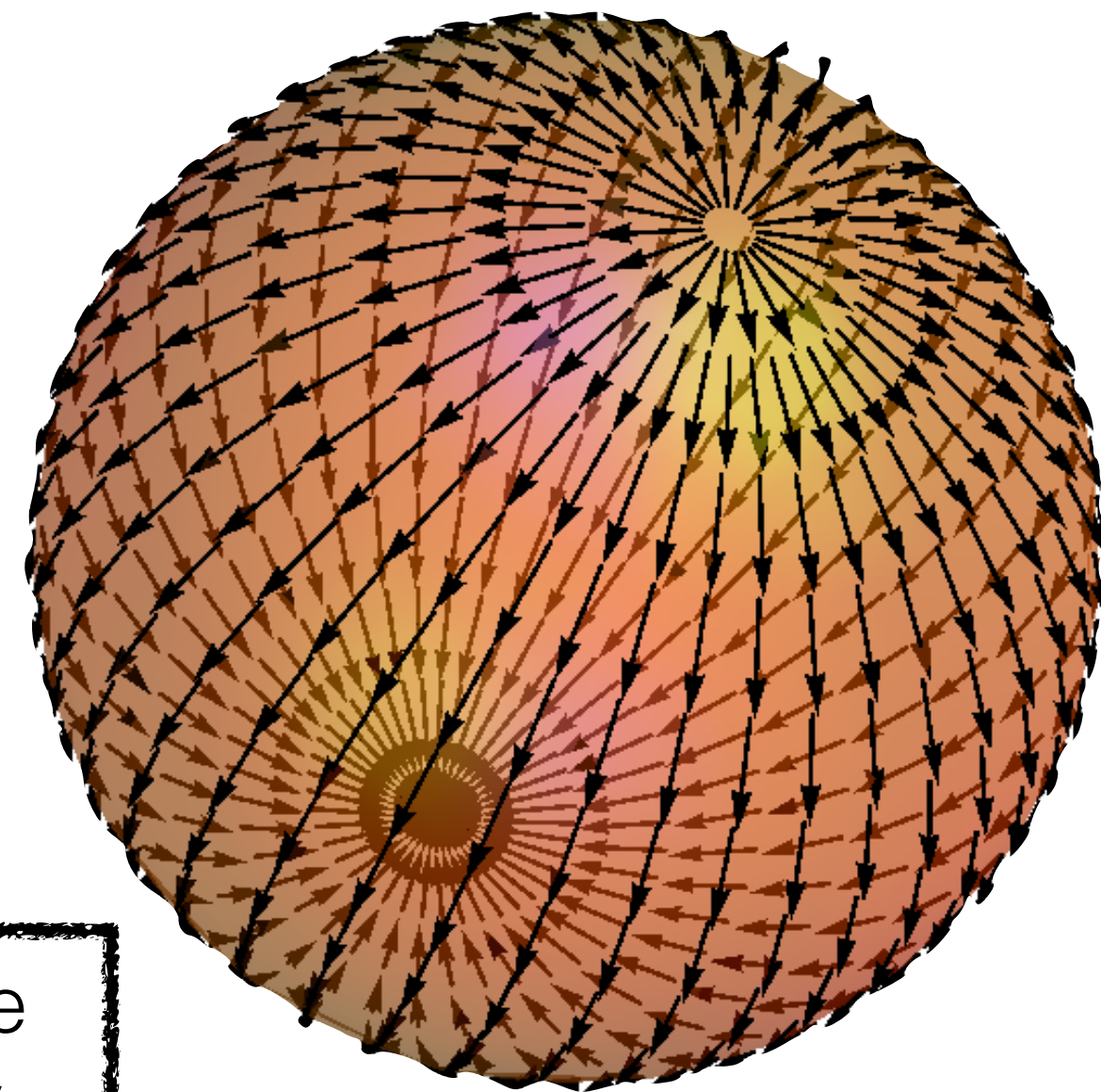


$d=2$, complex order parameter

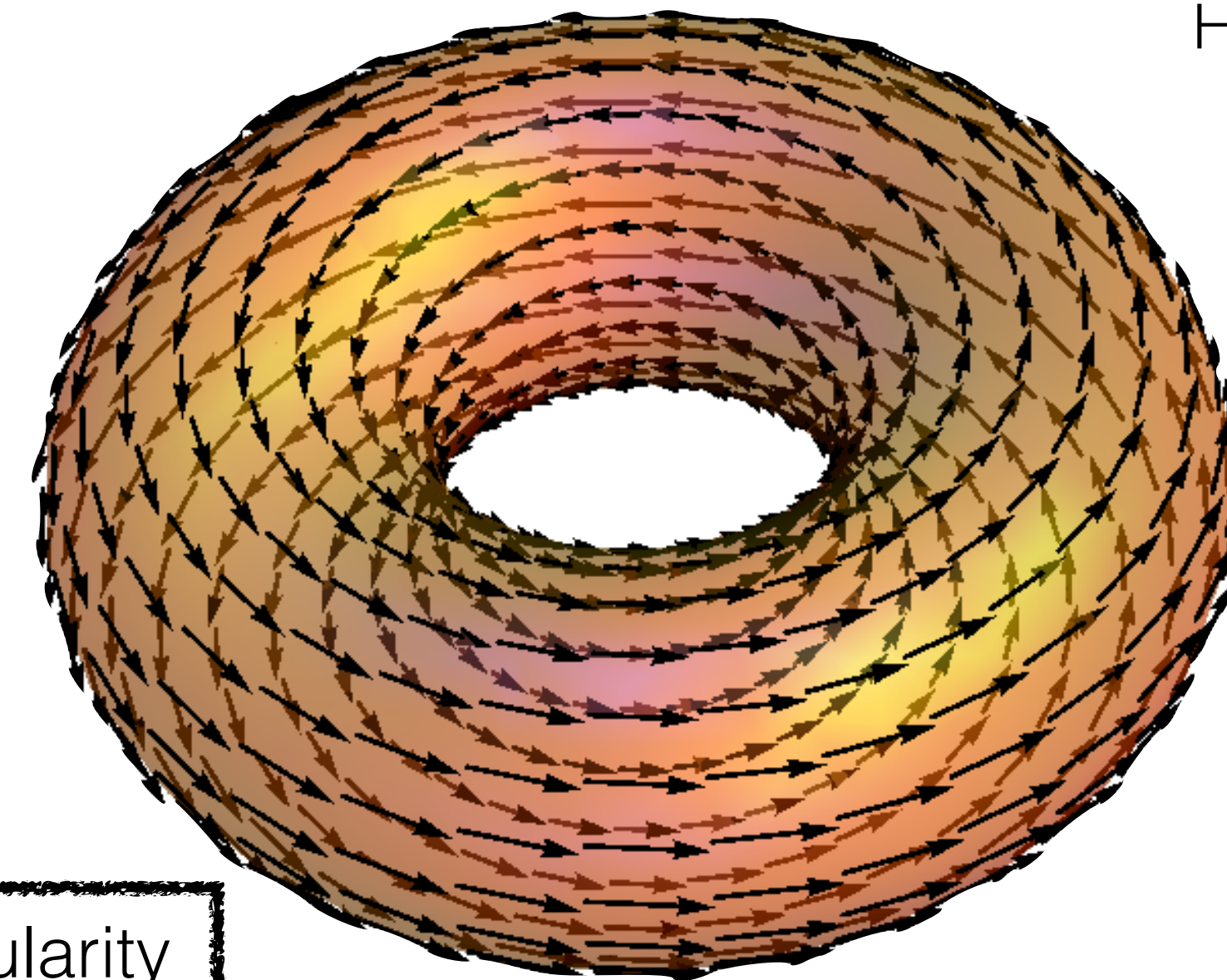
Topology

If vector field defined on a manifold : **non singular vector fields** do not necessarily exist

Hairy ball theorem



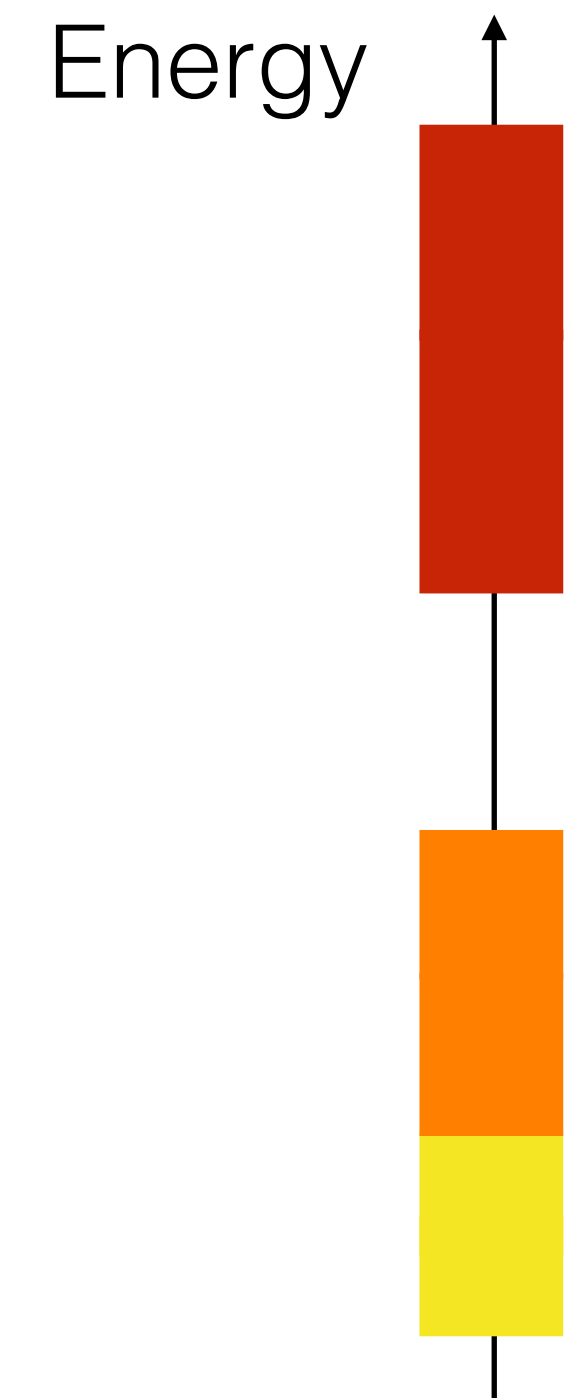
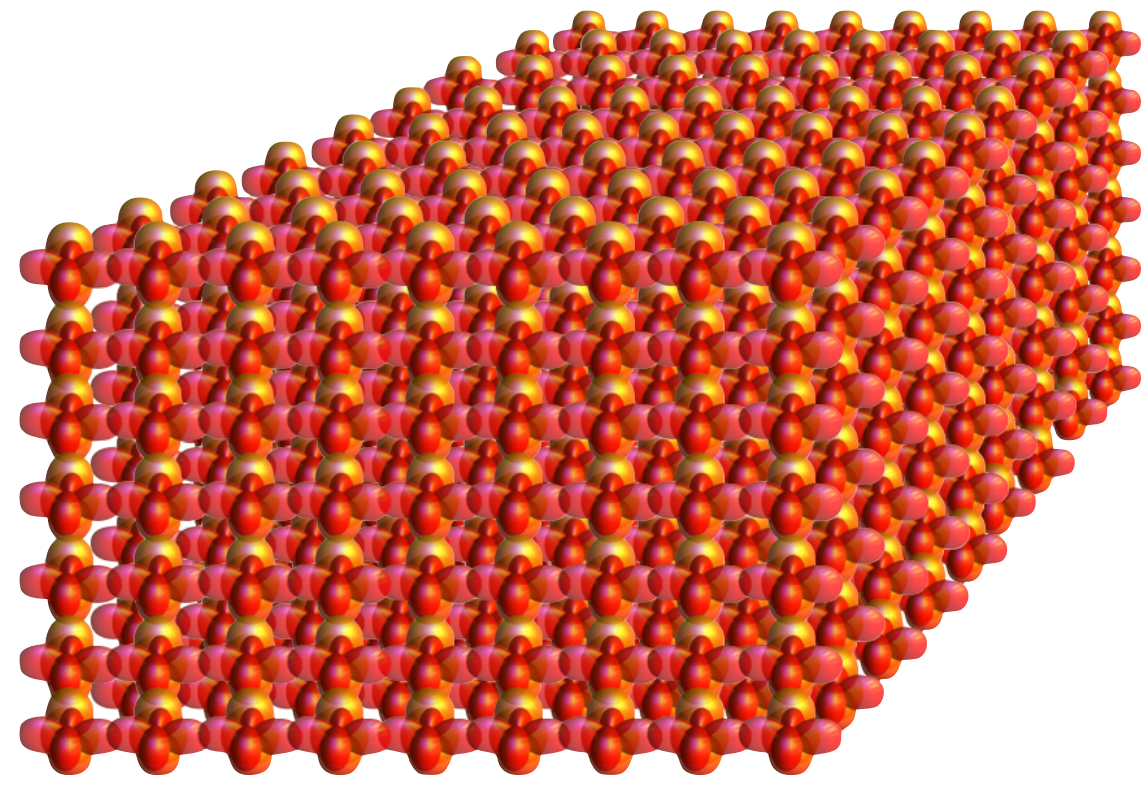
at least one
singularity



no singularity

- ▶ defines a **vector bundle** (manifold + vector space above each point)
- ▶ all vector fields singular \leftrightarrow **non trivial vector bundle**
- ▶ **topological property**, associated with a « topological Chern number »

Band Theory of Electrons in Solids



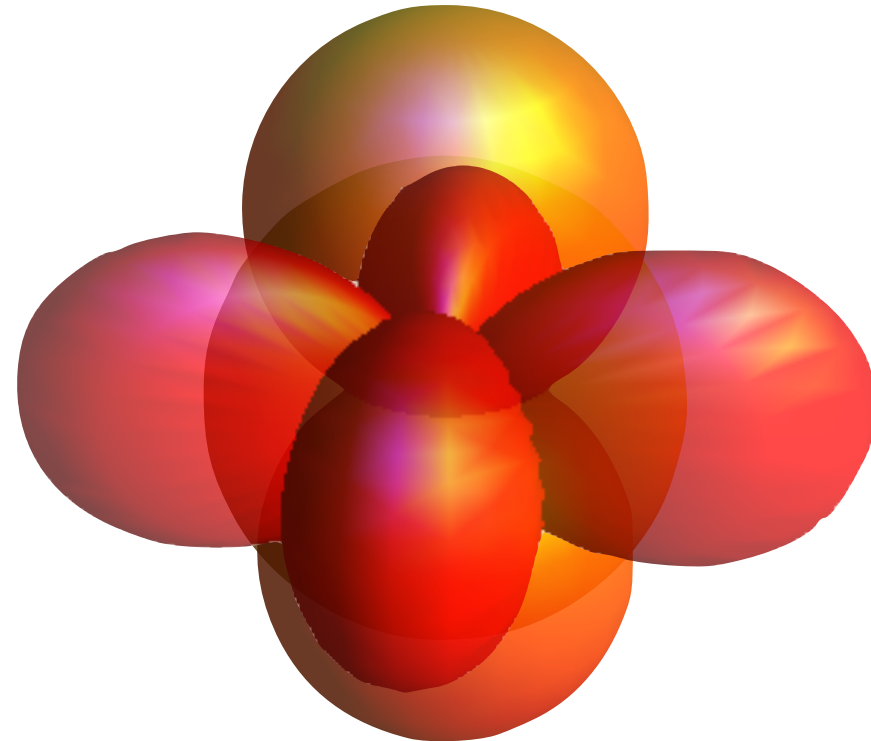
Band Theory of Electrons in Solids

Electronic Orbitals

1s

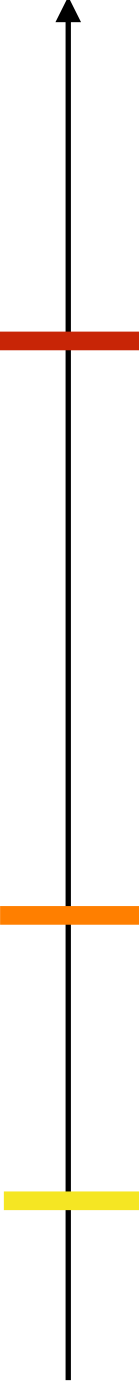
2p

3d



Single atom

Energy



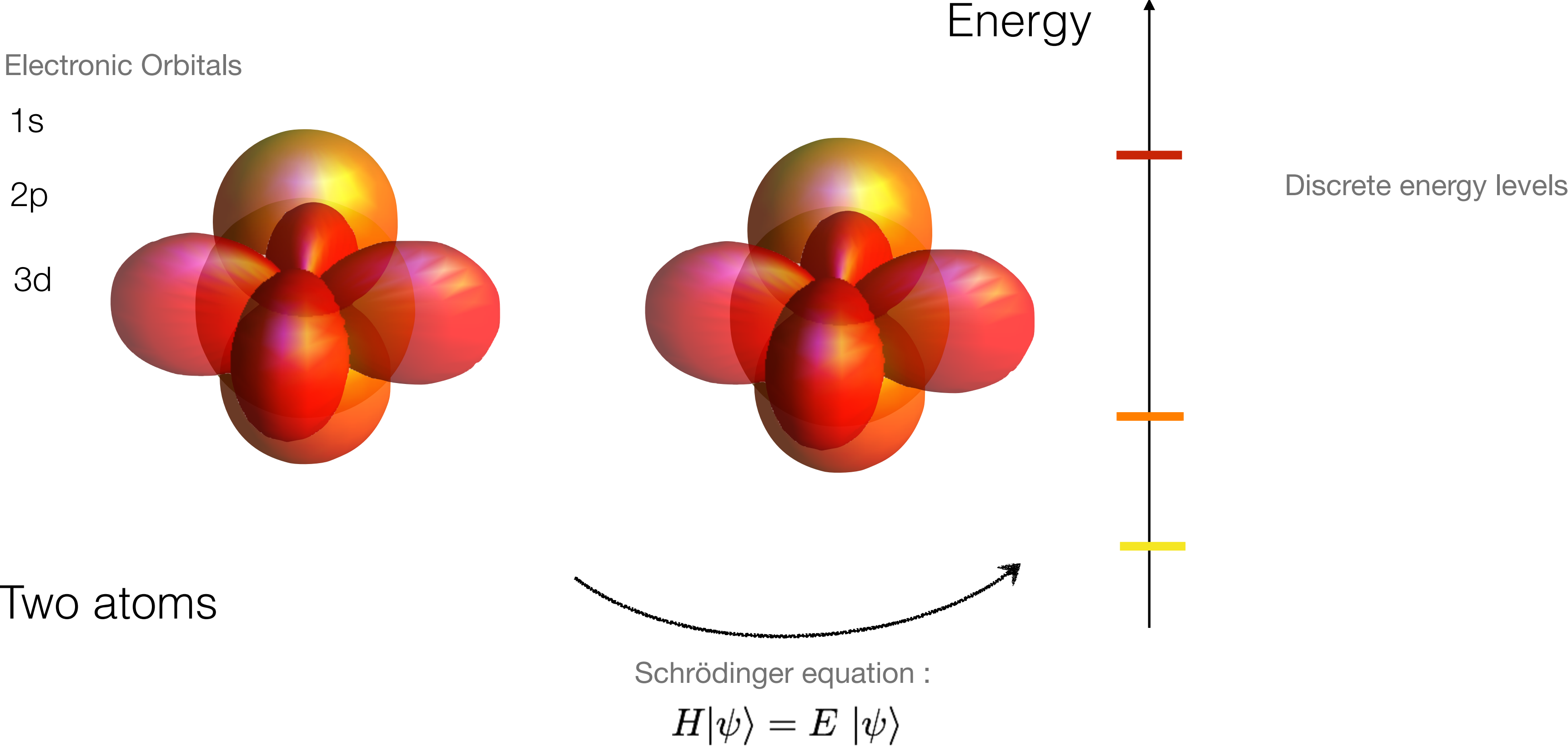
Discrete energy levels



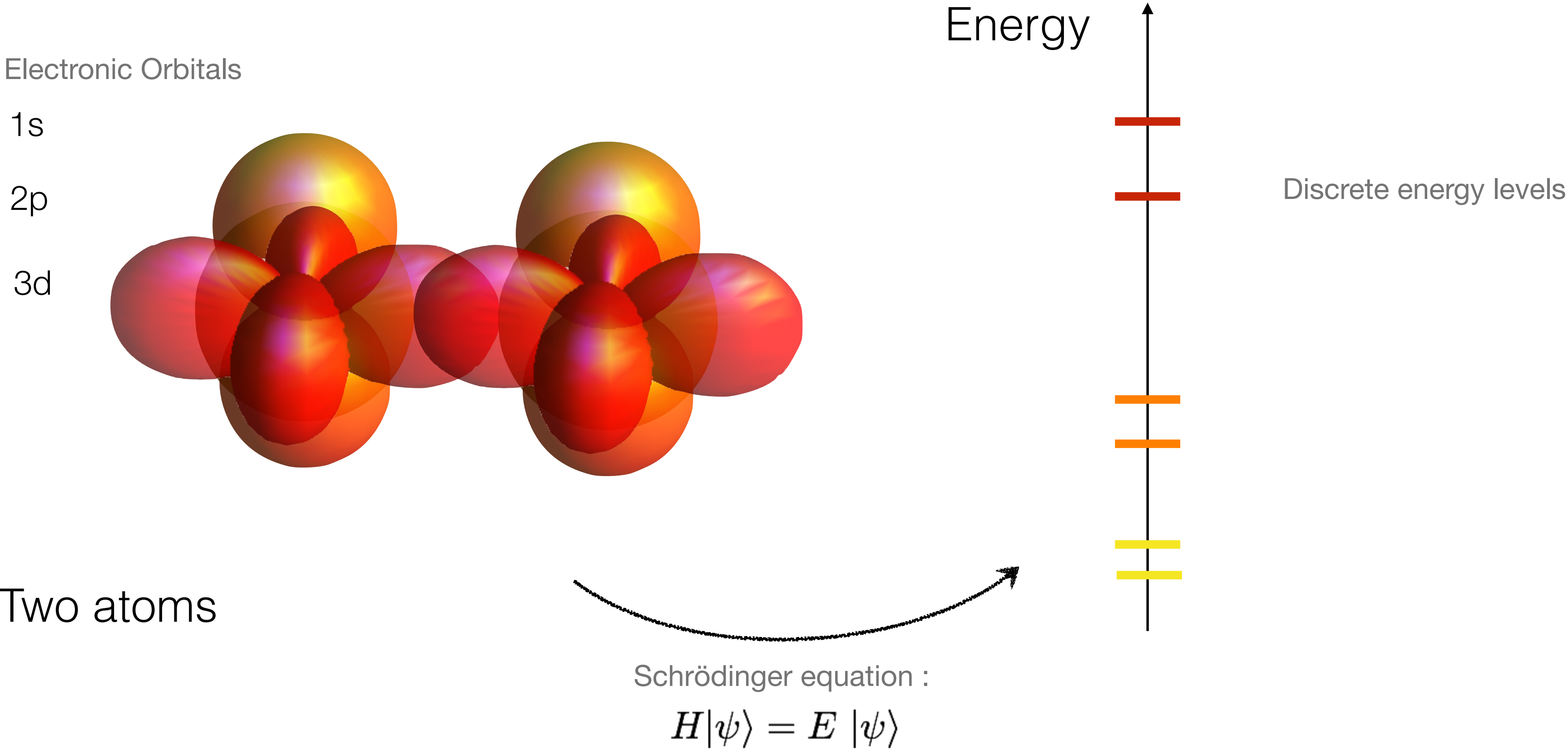
Schrödinger equation :

$$H|\psi\rangle = E |\psi\rangle$$

Band Theory of Electrons in Solids

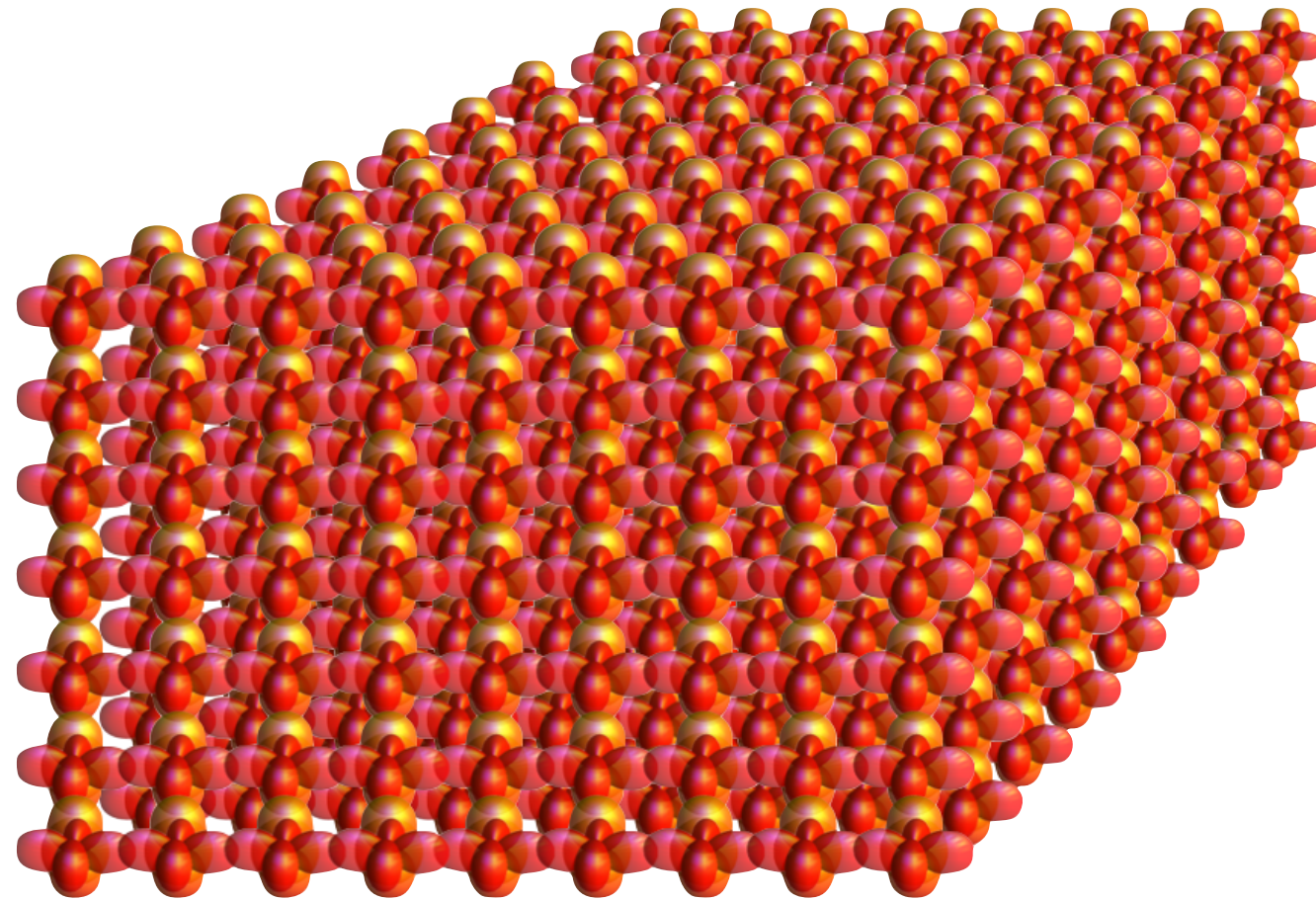


Band Theory of Electrons in Solids



Band Theory of Electrons in Solids

Electronic Orbitals

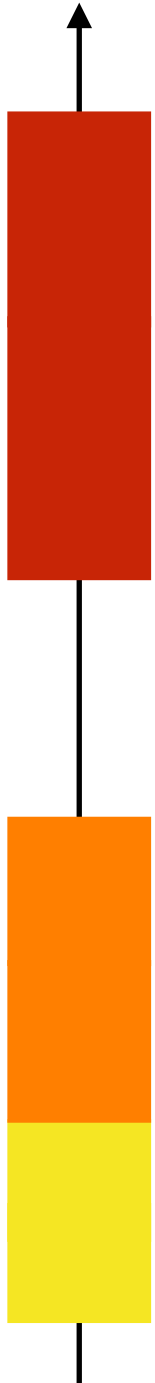


Solids : $\sim 10^{23}$ atoms



Schrödinger equation :
$$H|\psi\rangle = E |\psi\rangle$$

Energy

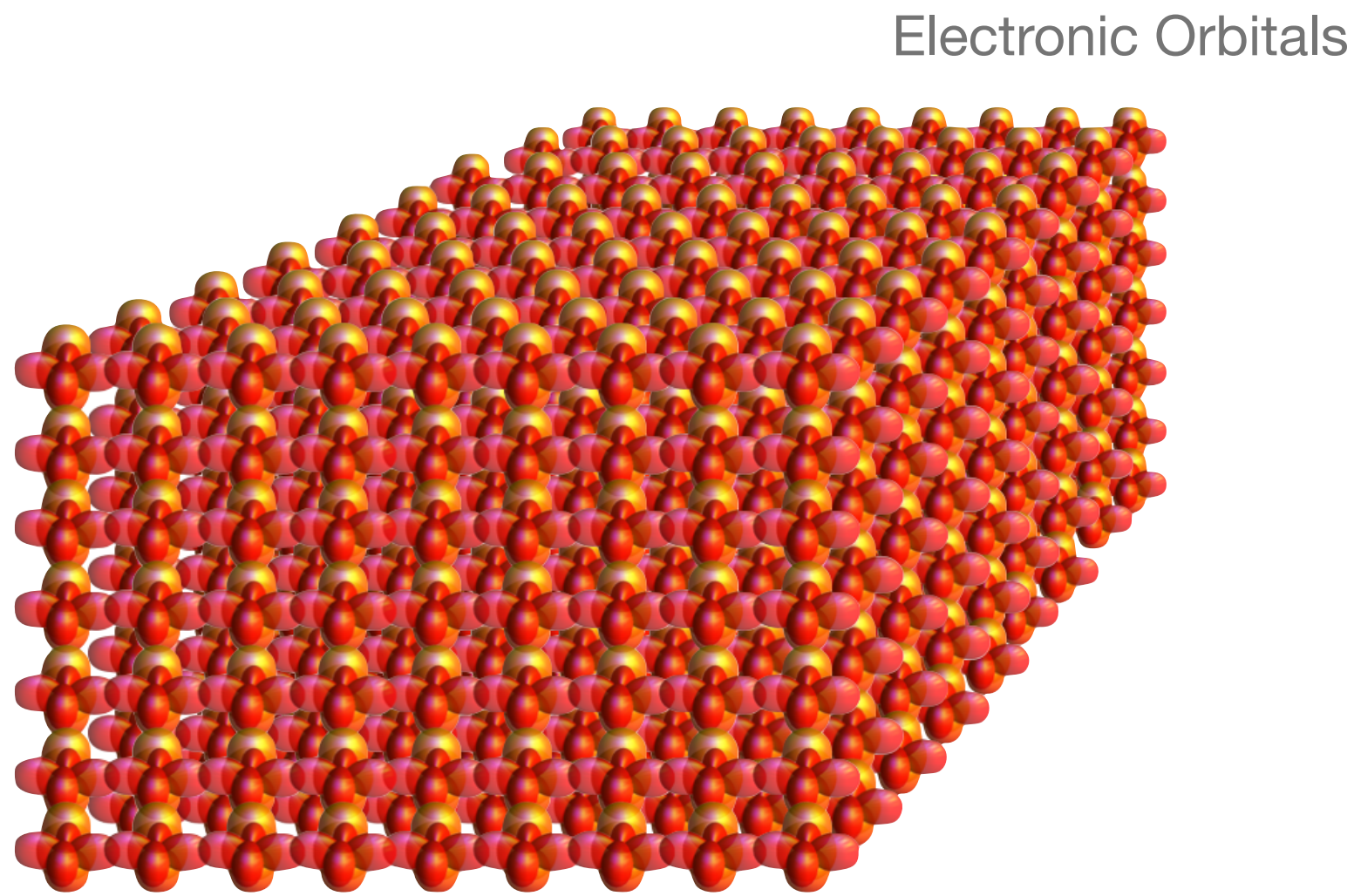


Discrete energy levels

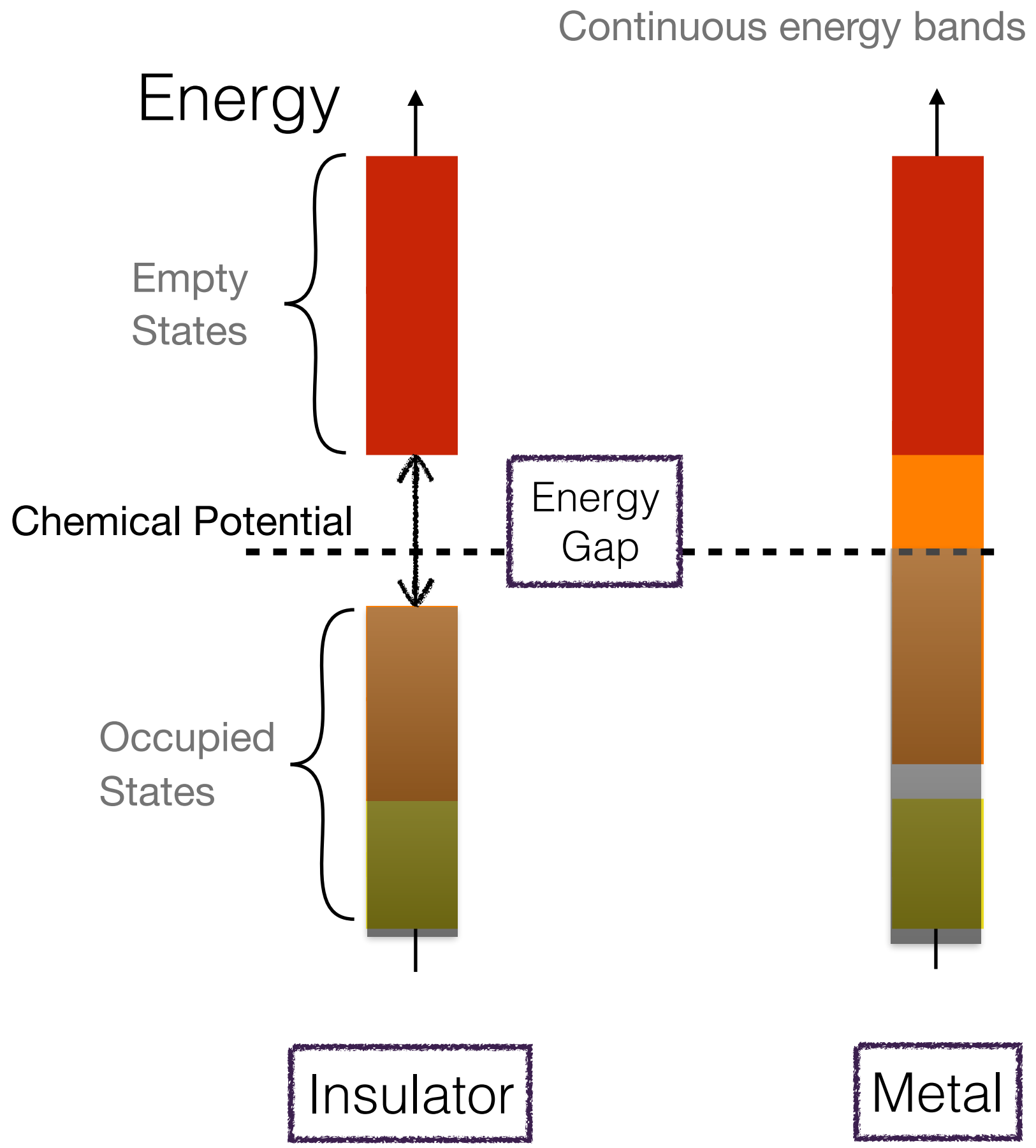


Continuous energy bands

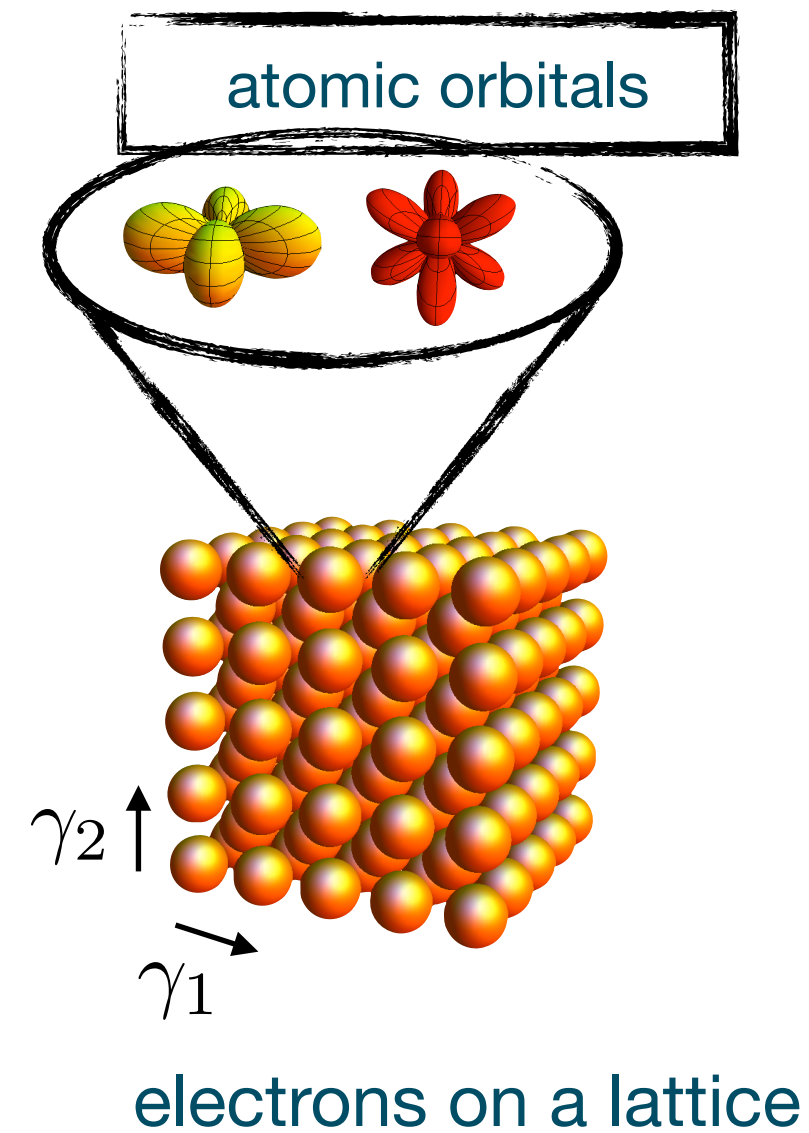
Band Theory of Electrons in Solids



Solids : $\sim 10^{23}$ atoms



Band Theory of Solids



- Band theory : single particle description of electronic states
 - ▶ Diagonalisation of a lattice Hamiltonian : $H_0|\psi\rangle = E|\psi\rangle$
- Periodicity of lattice (**symmetry !**) :
 - ▶ Bravais lattice : translations T_γ that leave physical lattice invariant
 - ▶ if $|\psi\rangle$ eigenstate, then $T_\gamma|\psi\rangle$ also with same energy $T_\gamma\psi(x) = \psi(x - \gamma)$
 - ▶ diagonalize simultaneously H_0 and T_γ
- Bloch wavefunctions :
 - ▶ Eigenstates of translations : $T_\gamma\psi(x) = \psi(x - \gamma) = e^{ik \cdot \gamma}\psi(x)$
 - ▶ labelled by quasi-momentum k
 - ▶ k and $k + G$ label the same eigenvector $|\psi_k\rangle$ if
$$G \cdot \gamma = n 2\pi, \quad n \in \mathbb{Z} \text{ for all } \gamma$$

Band Theory of Solids

- Bloch wavefunctions :

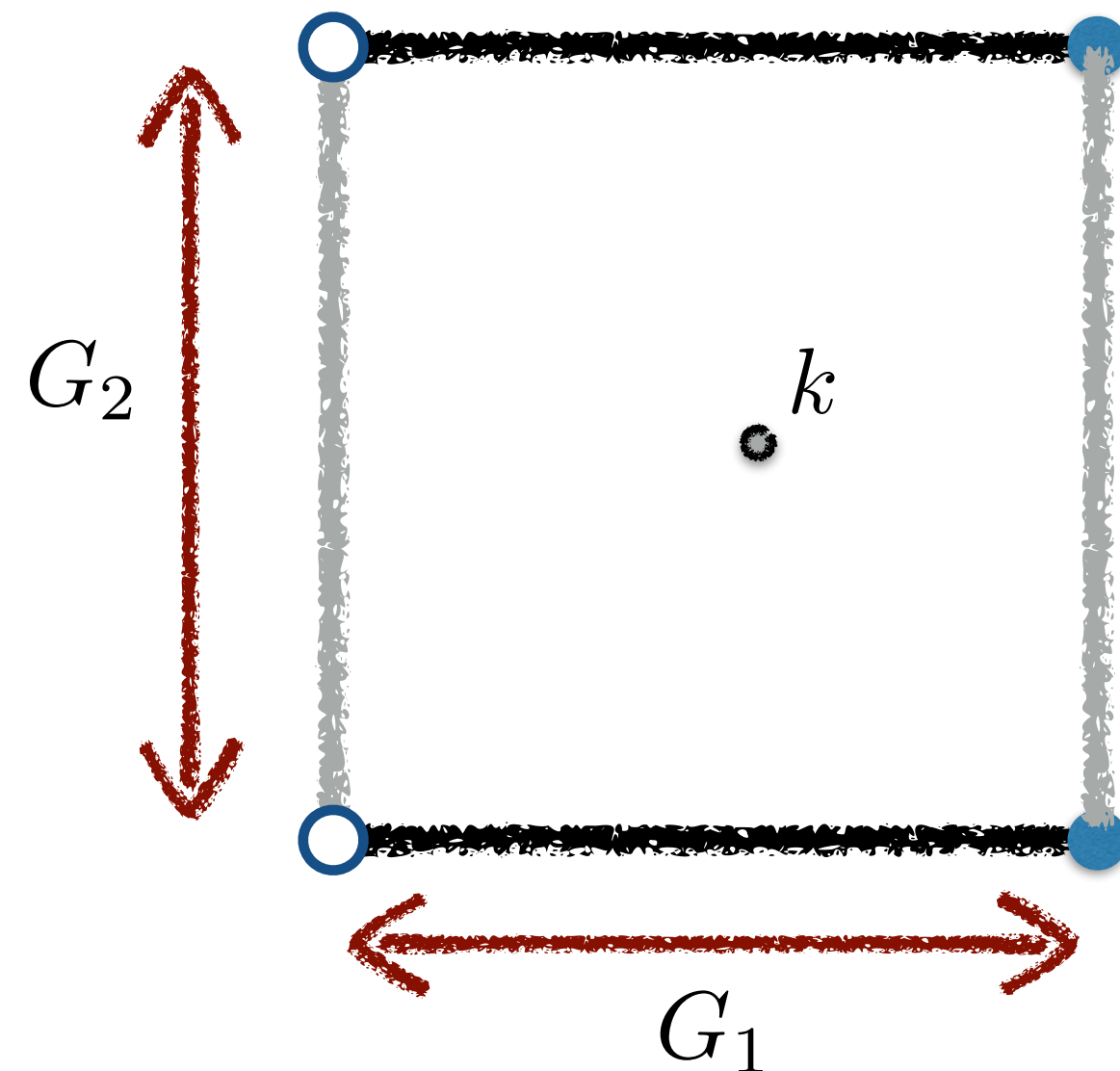
- ▶ Eigenstates of translations : $T_\gamma \psi(x) = \psi(x - \gamma) = e^{ik \cdot x} \psi(x)$

- ▶ labelled by quasi-momentum k

- ▶ k and $k + G$ label the same eigenvector $|\psi_k\rangle$ if

$$G \cdot \gamma = n 2\pi, \quad n \in \mathbb{Z} \text{ for all } \gamma$$

- ▶ k lies in **Brillouin Zone**



Band Theory of Solids

- Bloch wavefunctions :

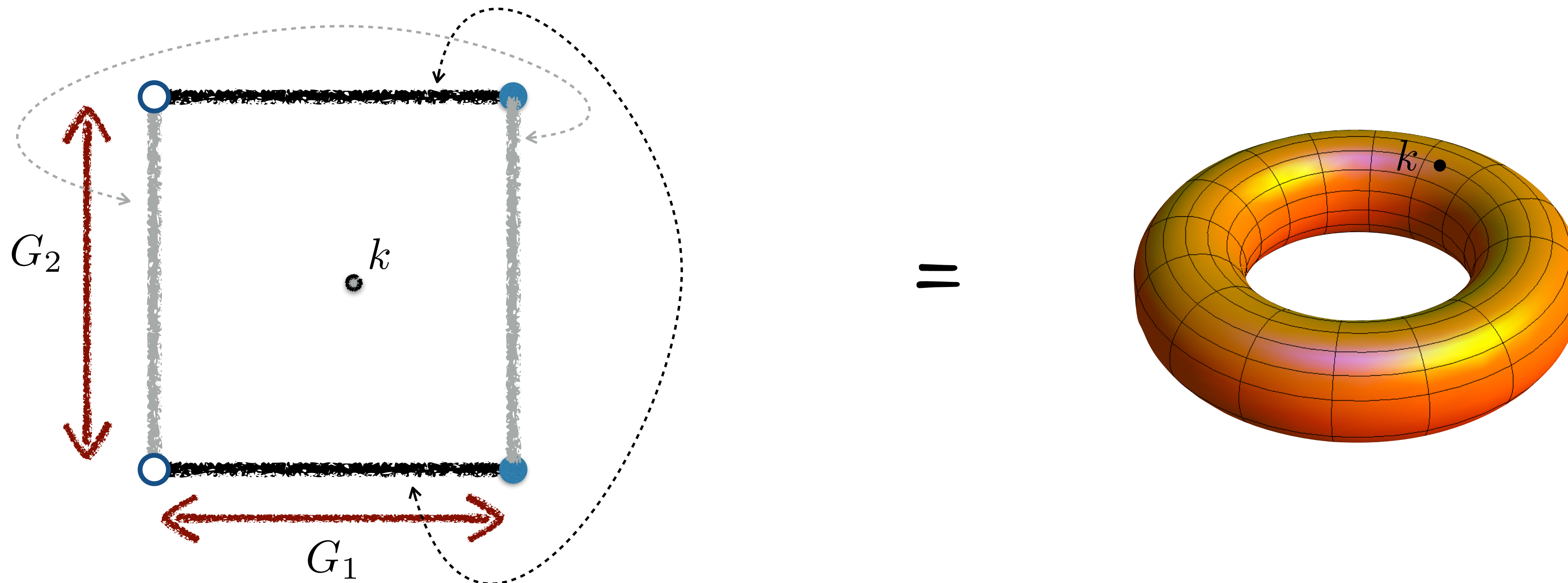
- ▶ Eigenstates of translations : $T_\gamma \psi(x) = \psi(x - \gamma) = e^{ik \cdot x} \psi(x)$

- ▶ labelled by quasi-momentum k

- ▶ k and $k + G$ label the same eigenvector $|\psi_k\rangle$ if

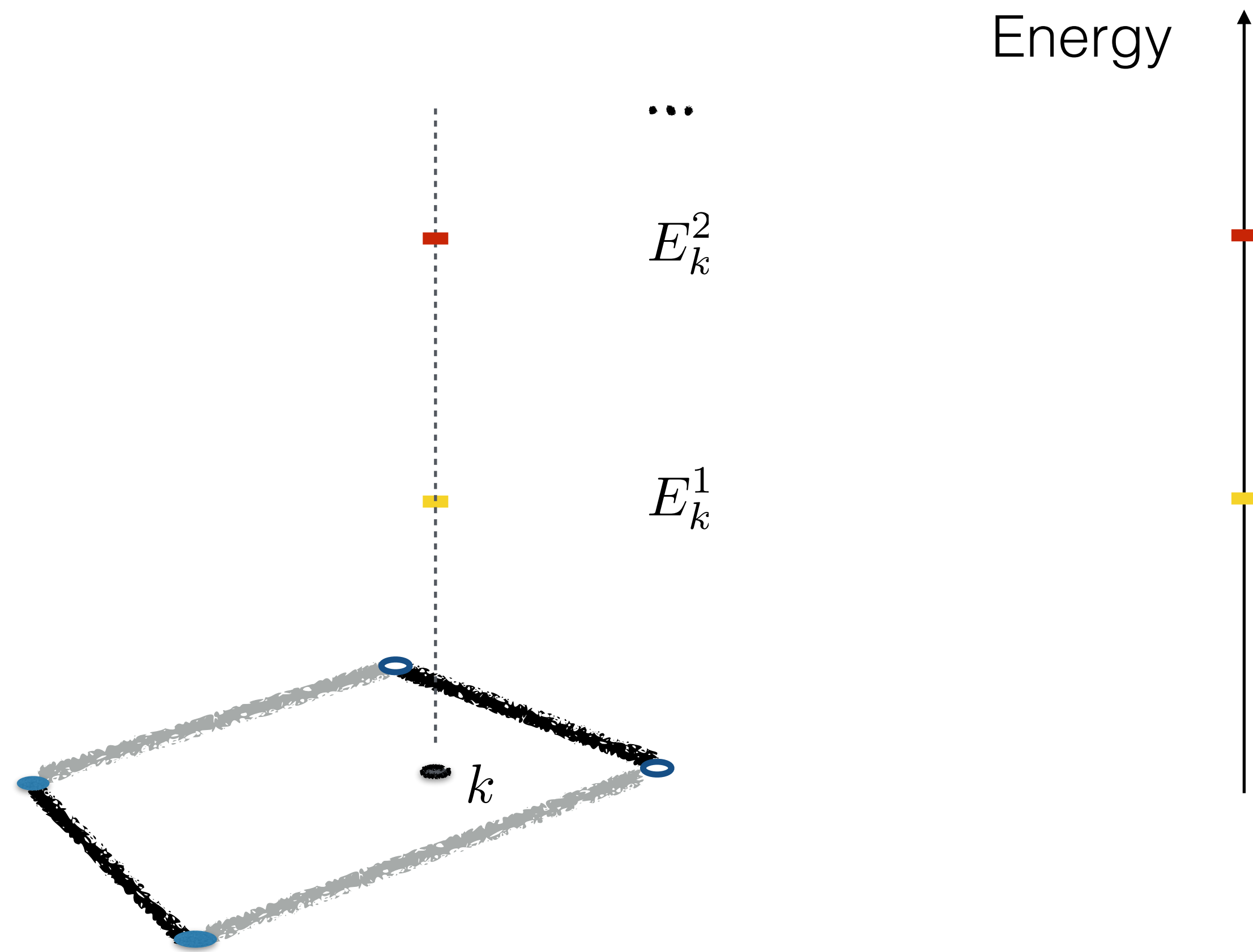
$$G \cdot \gamma = n 2\pi, \quad n \in \mathbb{Z} \text{ for all } \gamma$$

- ▶ k lies in **Brillouin Zone**



Band Theory of Solids

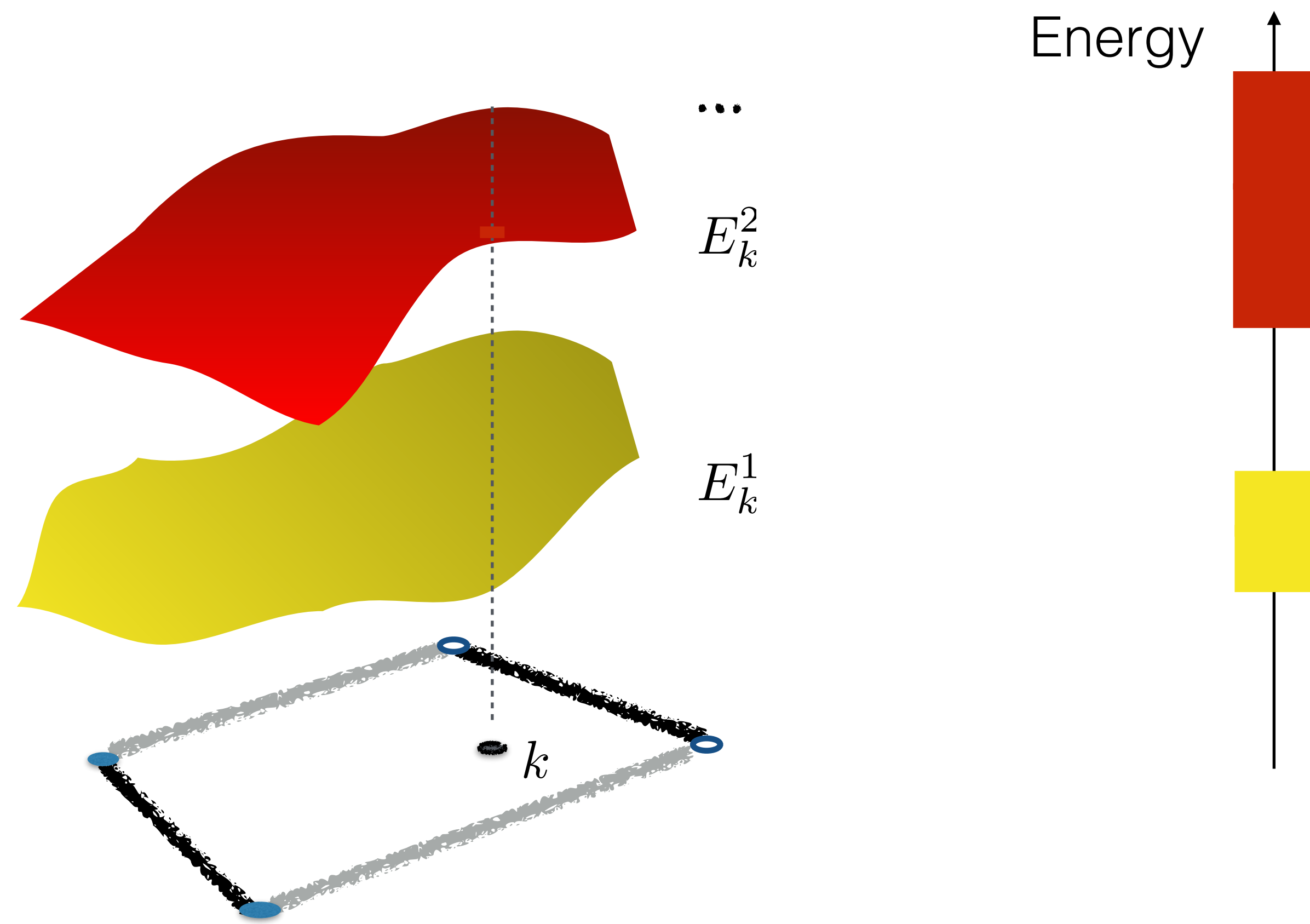
Diagonalization : $H|\psi_k^\alpha\rangle = E_k^\alpha|\psi_k^\alpha\rangle$



Band Theory of Solids

Diagonalization : $H|\psi_k^\alpha\rangle = E_k^\alpha|\psi_k^\alpha\rangle$

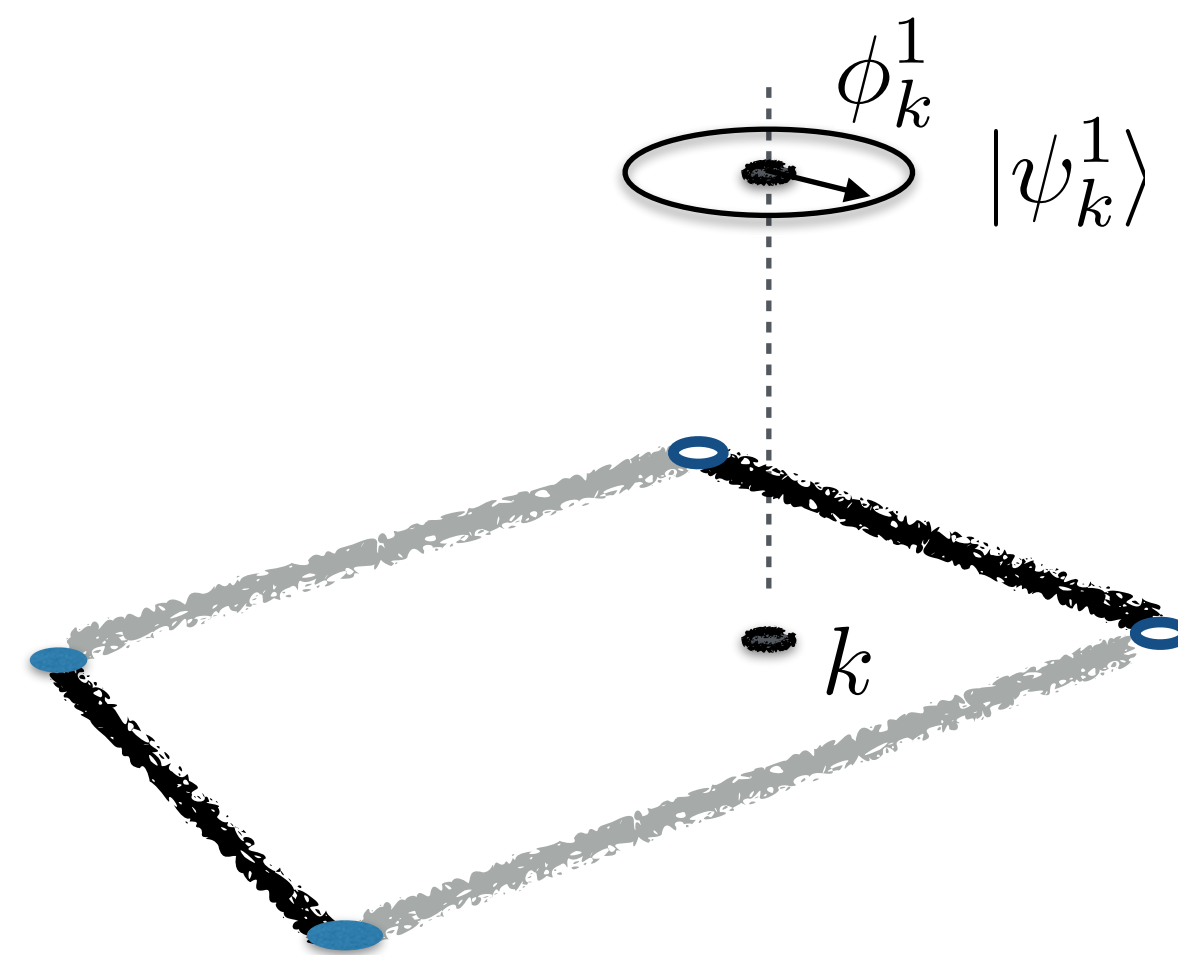
→ Energy bands : surfaces E_k^α over the Brillouin torus



Band Theory of Solids

Diagonalization : $H|\psi_k^\alpha\rangle = E_k^\alpha|\psi_k^\alpha\rangle$

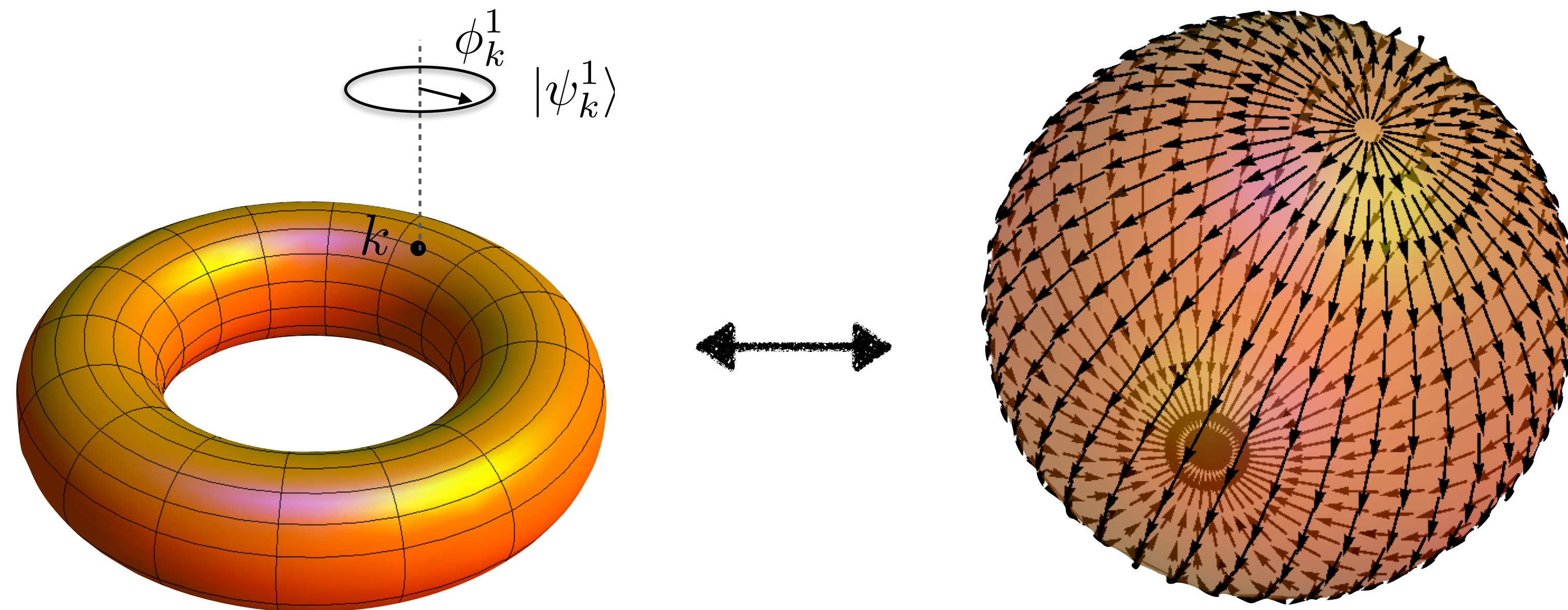
- Each band : **Vector bundle** of vectors $|\psi_k^\alpha\rangle$ over the Brillouin zone
- eigenvectors $|\psi_k^\alpha\rangle$ defined up to a phase $|\psi_k^\alpha\rangle \rightarrow e^{i\phi_k^\alpha}|\psi_k^\alpha\rangle$



Band Theory of Solids

Diagonalization : $H|\psi_k^\alpha\rangle = E_k^\alpha|\psi_k^\alpha\rangle$

- Each band : **Vector bundle** of vectors $|\psi_k^\alpha\rangle$ over the Brillouin zone
- eigenvectors $|\psi_k^\alpha\rangle$ defined up to a phase $|\psi_k^\alpha\rangle \rightarrow e^{i\phi_k^\alpha}|\psi_k^\alpha\rangle$

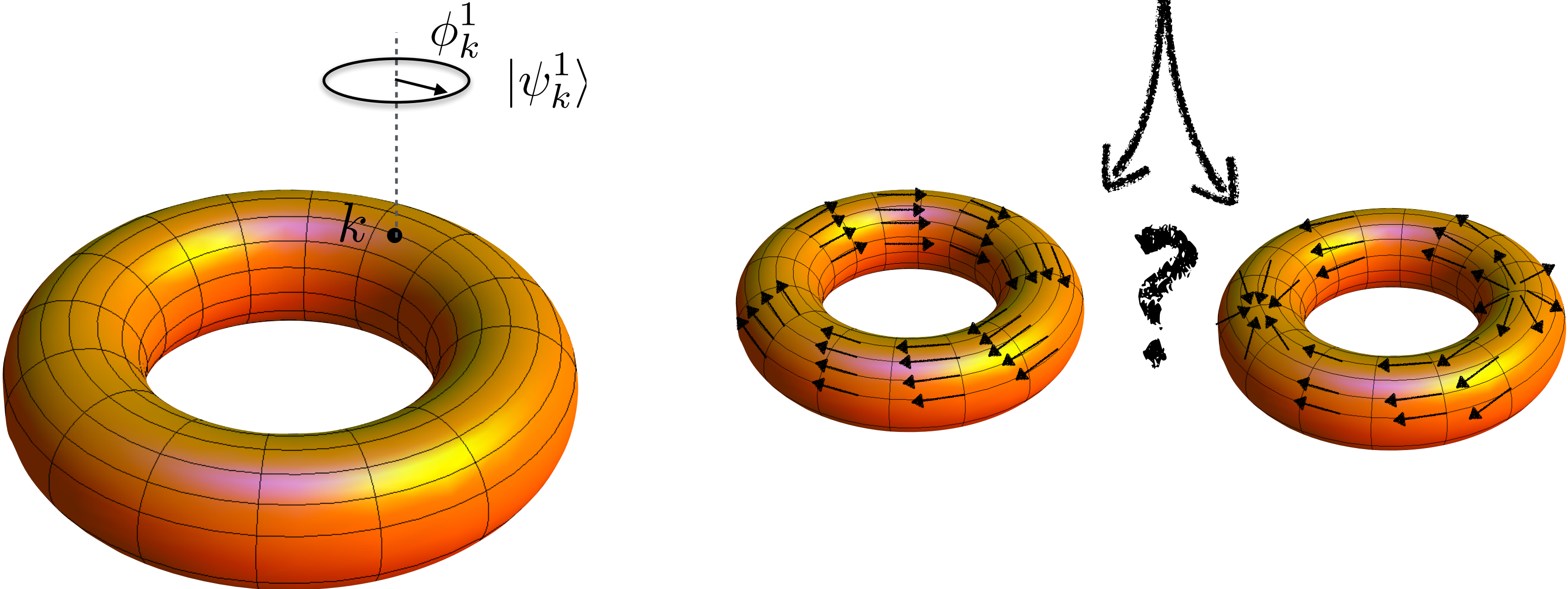


Band Theory of Solids

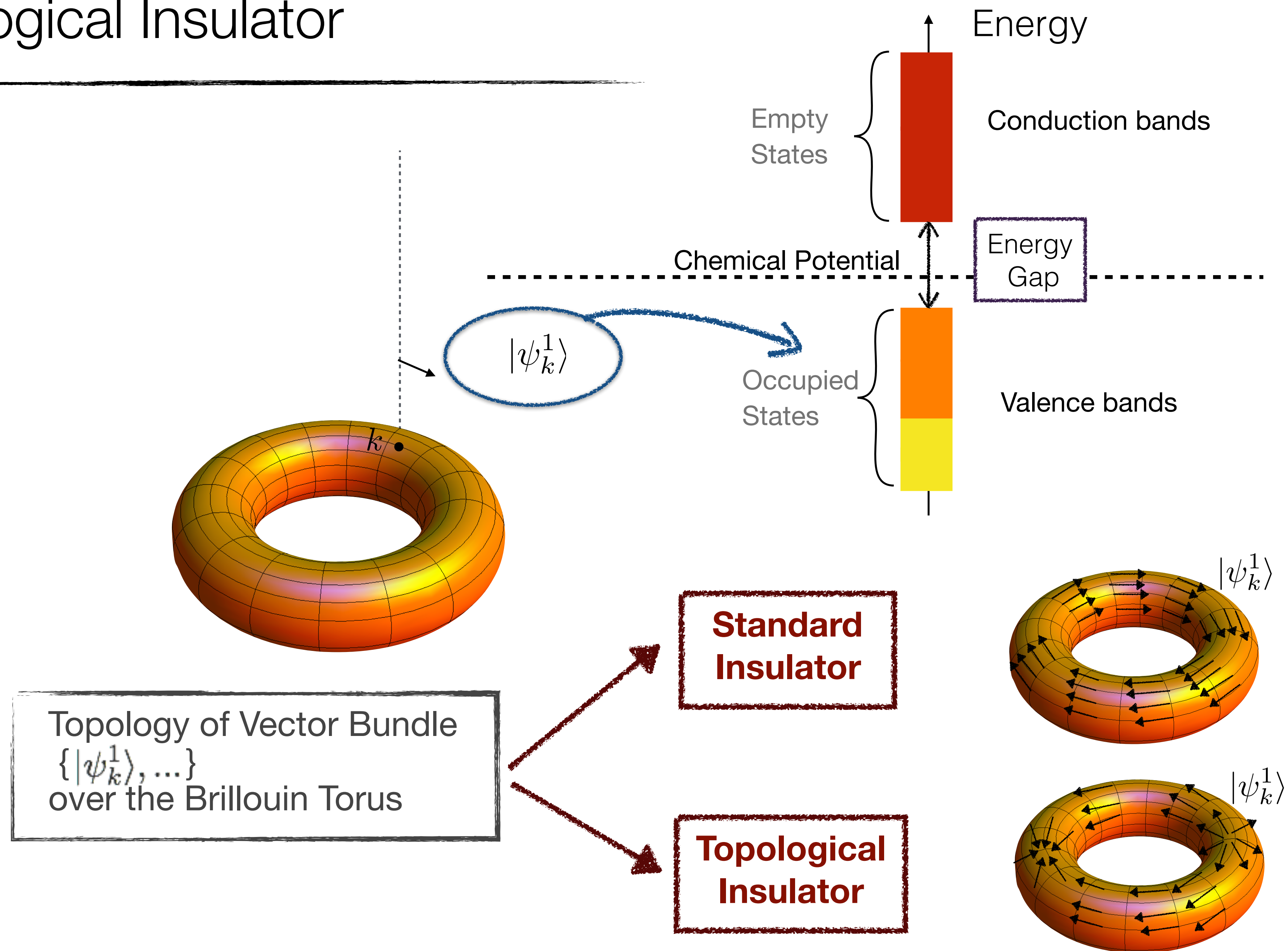
Diagonalization : $H|\psi_k^\alpha\rangle = E_k^\alpha|\psi_k^\alpha\rangle$

- Each band : **Vector bundle** of vectors $|\psi_k^\alpha\rangle$ over the Brillouin zone
- eigenvectors $|\psi_k^\alpha\rangle$ defined up to a phase $|\psi_k^\alpha\rangle \rightarrow e^{i\phi_k^\alpha}|\psi_k^\alpha\rangle$

Is it always topologically trivial ?
(i.e. continuous vector field exists)



Topological Insulator



Topological index for bands

Thouless *et al.*, (1982)
 Berry (1984)
 see also Fruchart *et al.*, (2014)

- **Berry Connexion** form (analogous to electr. potential)

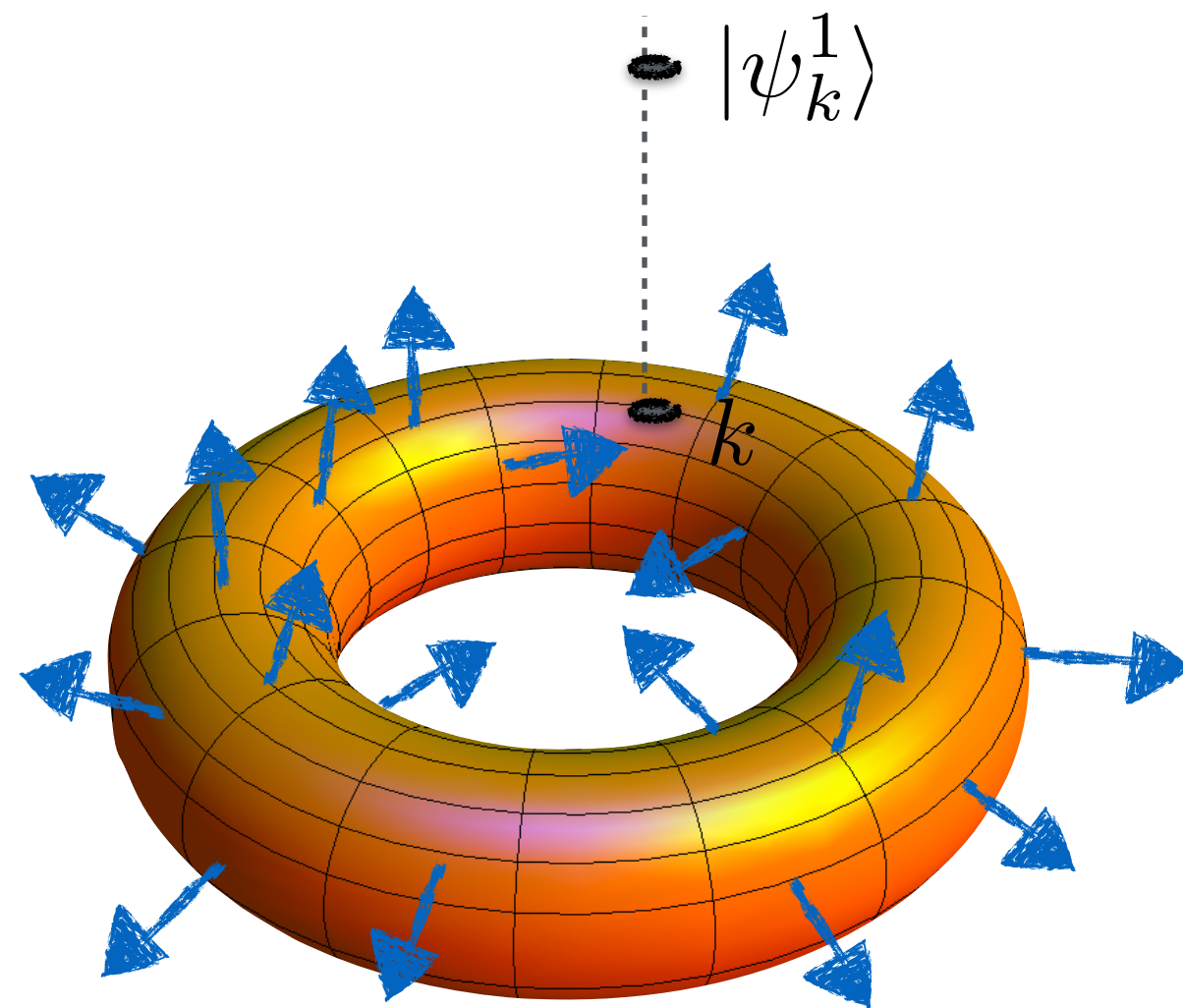
$$A_k = \frac{1}{i} \langle u_k^1 | \nabla_k | u_k^1 \rangle \quad |\psi_k^1\rangle = e^{ik \cdot \hat{r}} |u_k^1\rangle$$

- **Berry curvature** (analogous to Flux) :

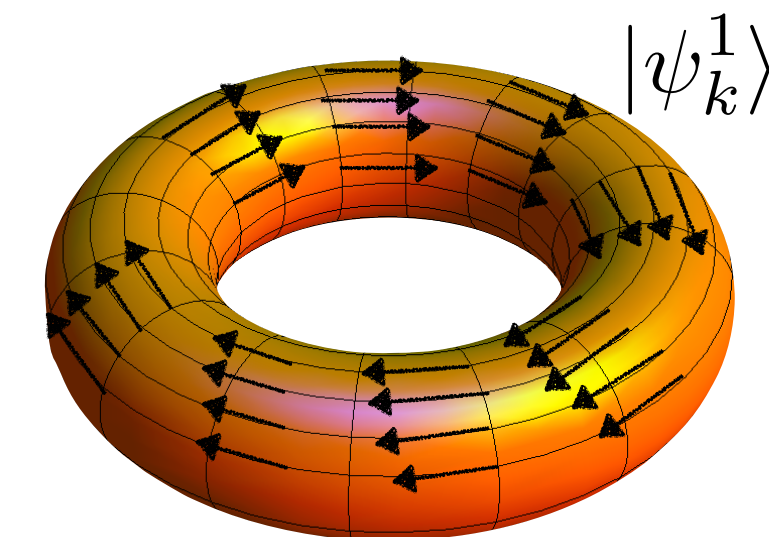
$$F_k = \nabla_k \times A_k$$

- **Topological number** : Chern number

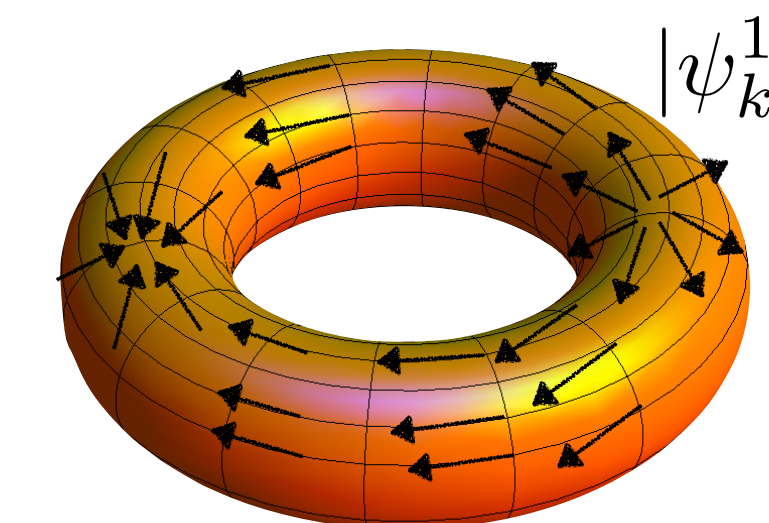
$$C_1 = \frac{1}{2\pi} \int_{\text{BZ}} F$$



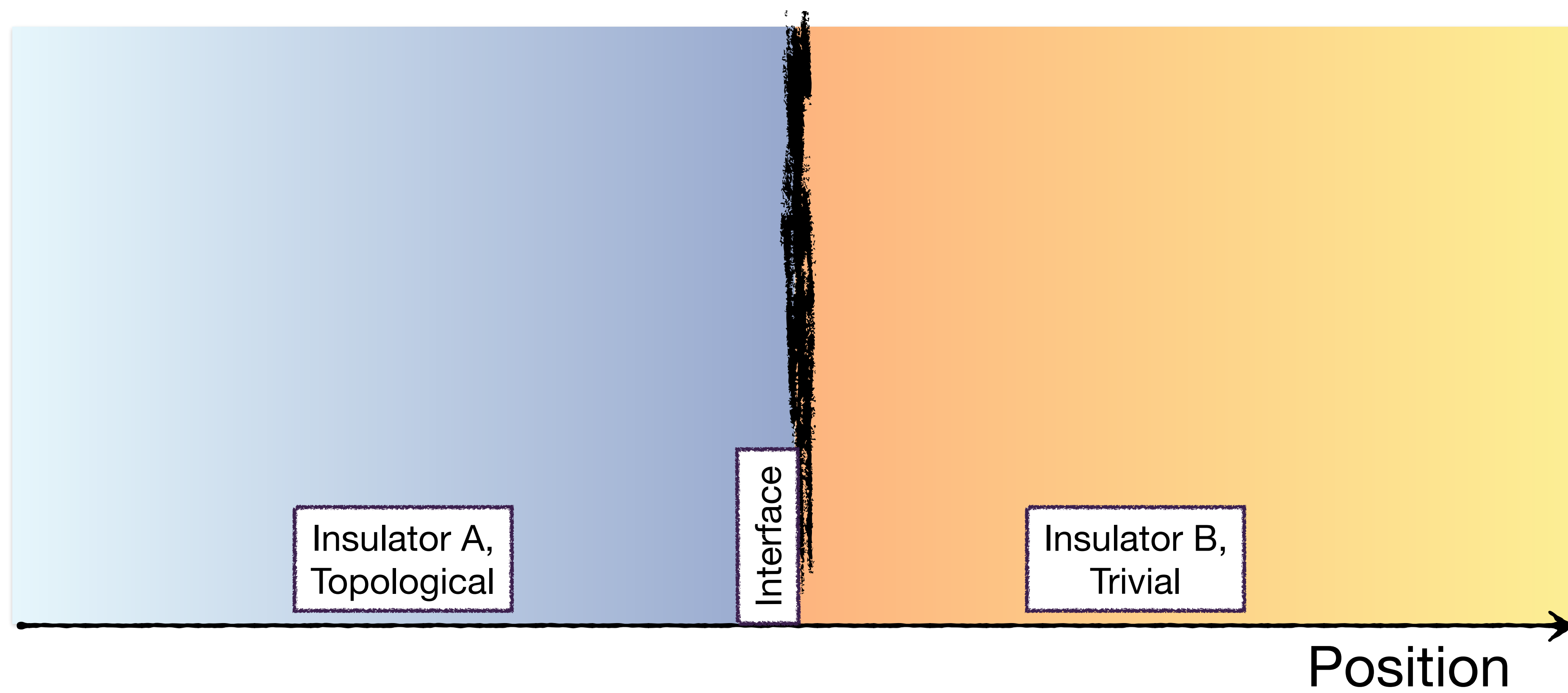
Chern number $C = 0$



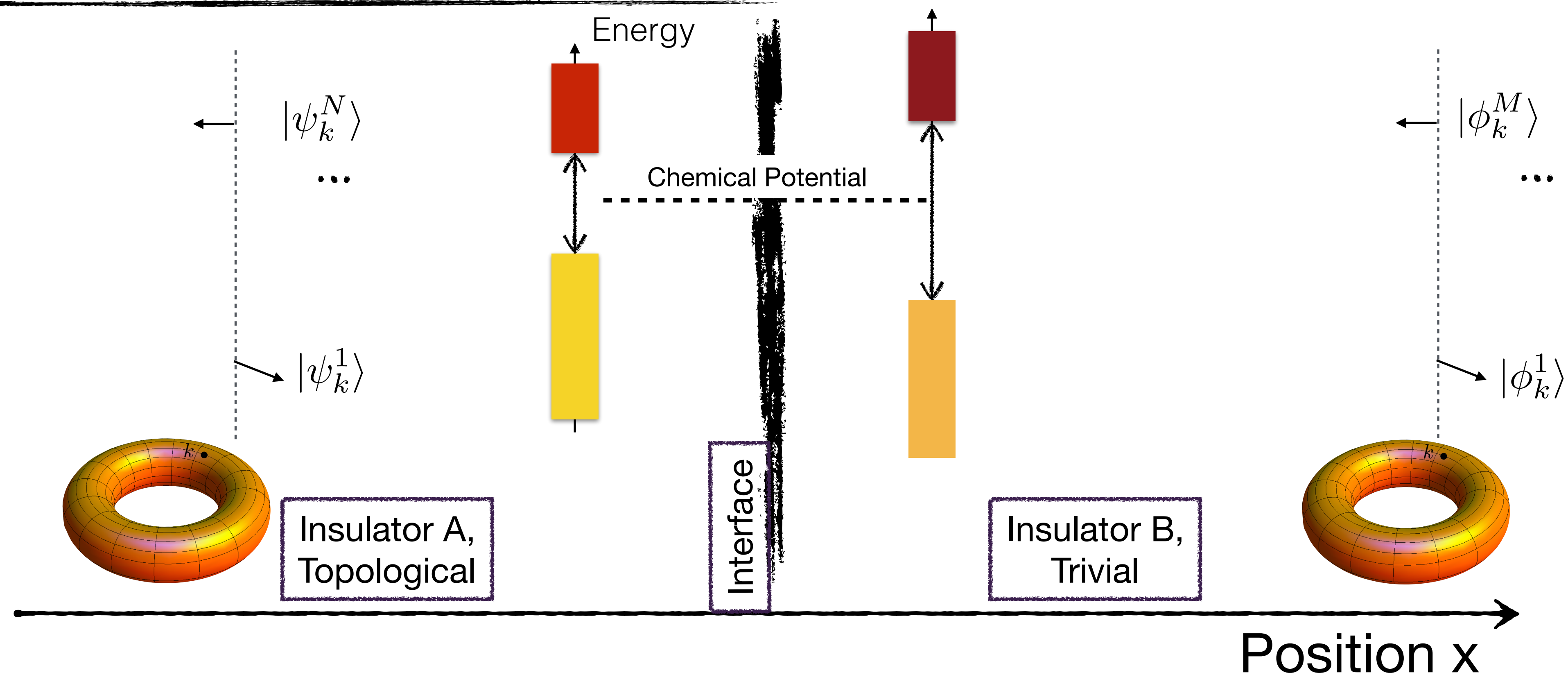
Chern number $C \neq 0$



Interface between Insulators



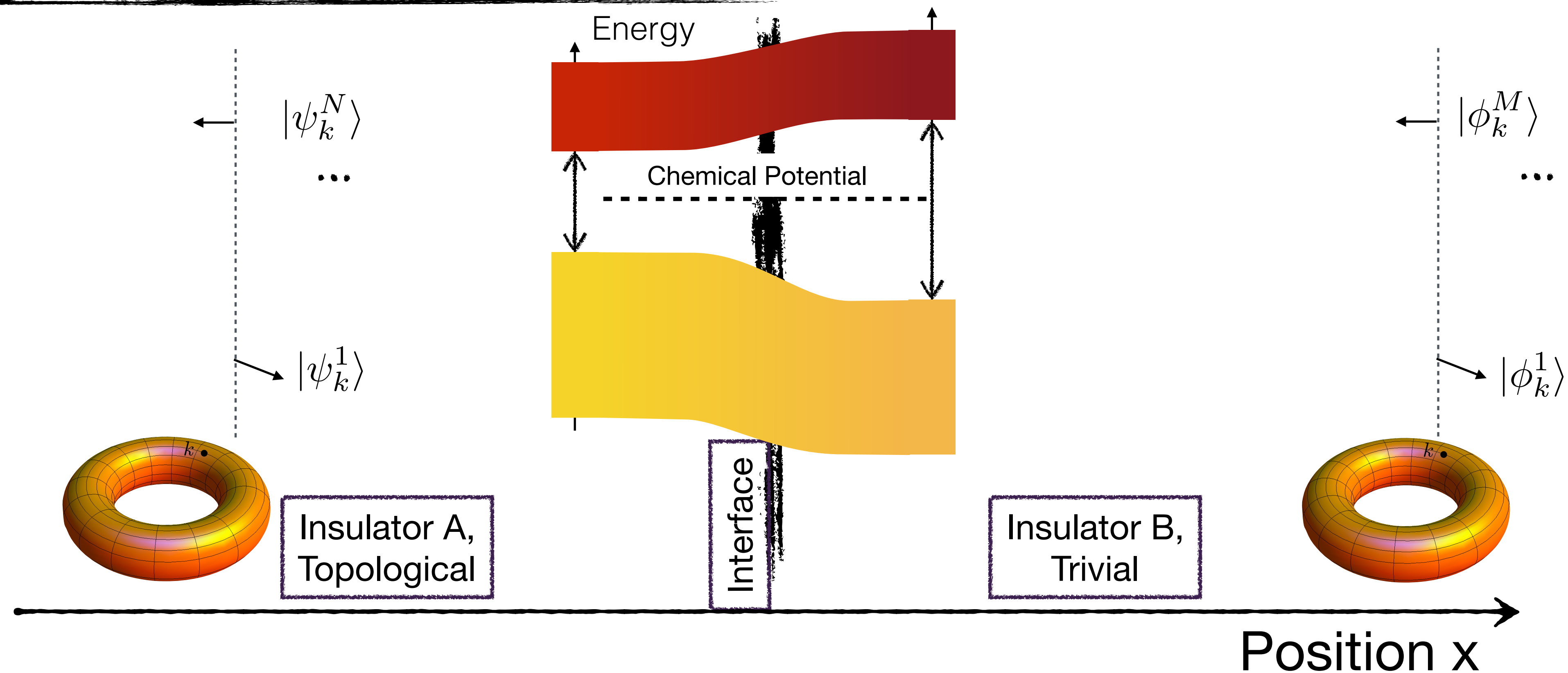
Interface between Insulators



Description of the interface :

- ▶ consider eigenstates $\{|\psi_k^\alpha\rangle\}$ and $\{|\phi_k^\beta\rangle\}$ of materials on both sides
- ▶ extrapolate (position dependent Hamiltonian $H(x)$)

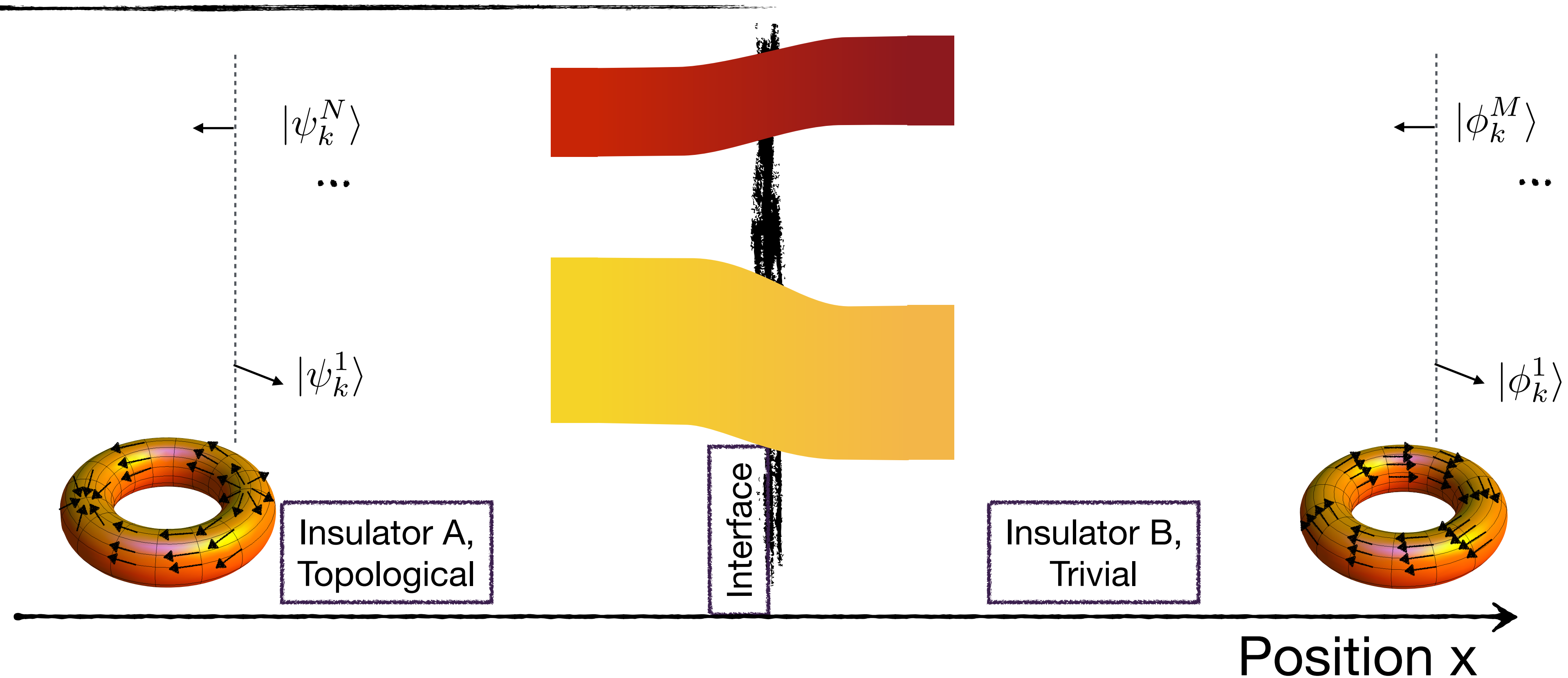
Interface between Insulators



Description of the interface :

- ▶ consider eigenstates $\{|\psi_k^\alpha\rangle\}$ and $\{|\phi_k^\beta\rangle\}$ of materials on both sides
- ▶ extrapolate (position dependent Hamiltonian $H(x)$)

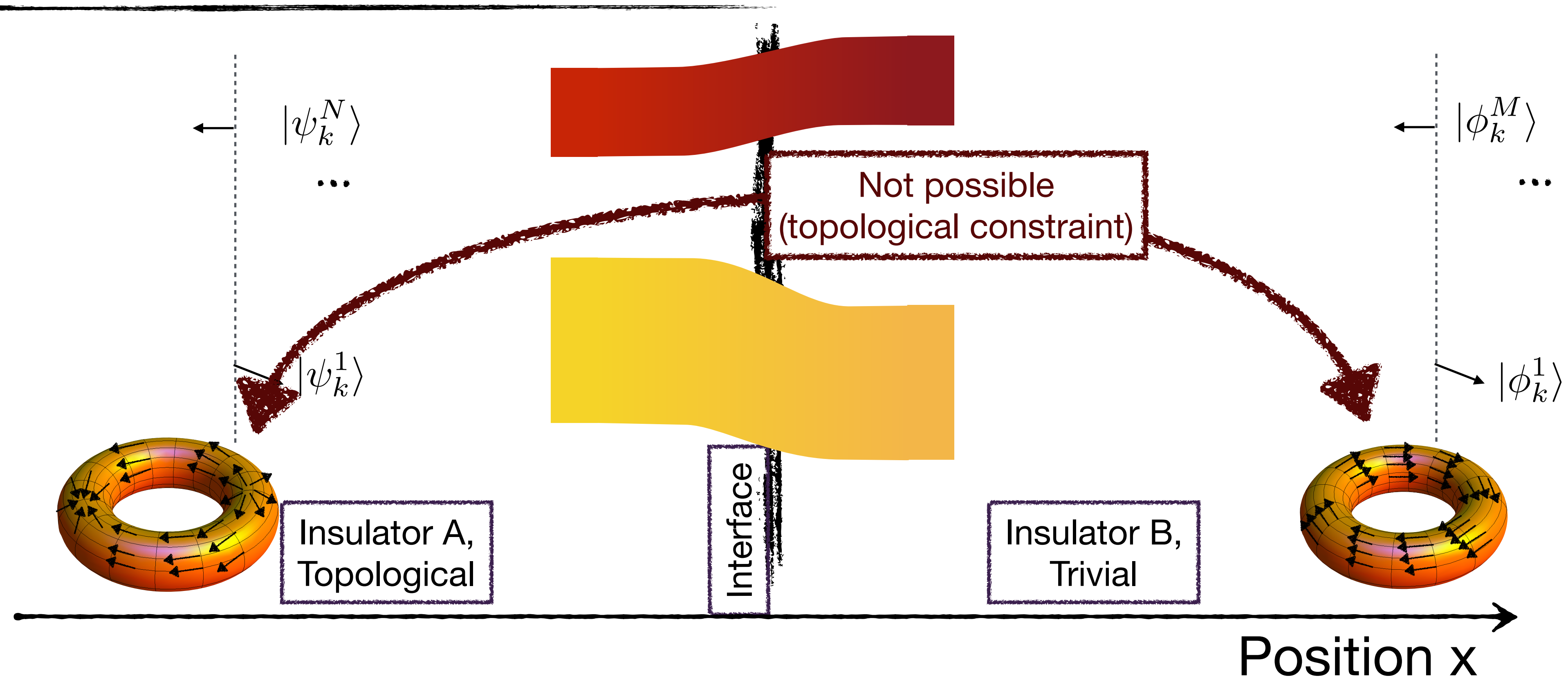
Interface between Insulators



Description of the interface :

- ▶ consider eigenstates $\{|\psi_k^\alpha\rangle\}$ and $\{|\phi_k^\beta\rangle\}$ of materials on both sides
- ▶ extrapolate (position dependent Hamiltonian $H(x)$)

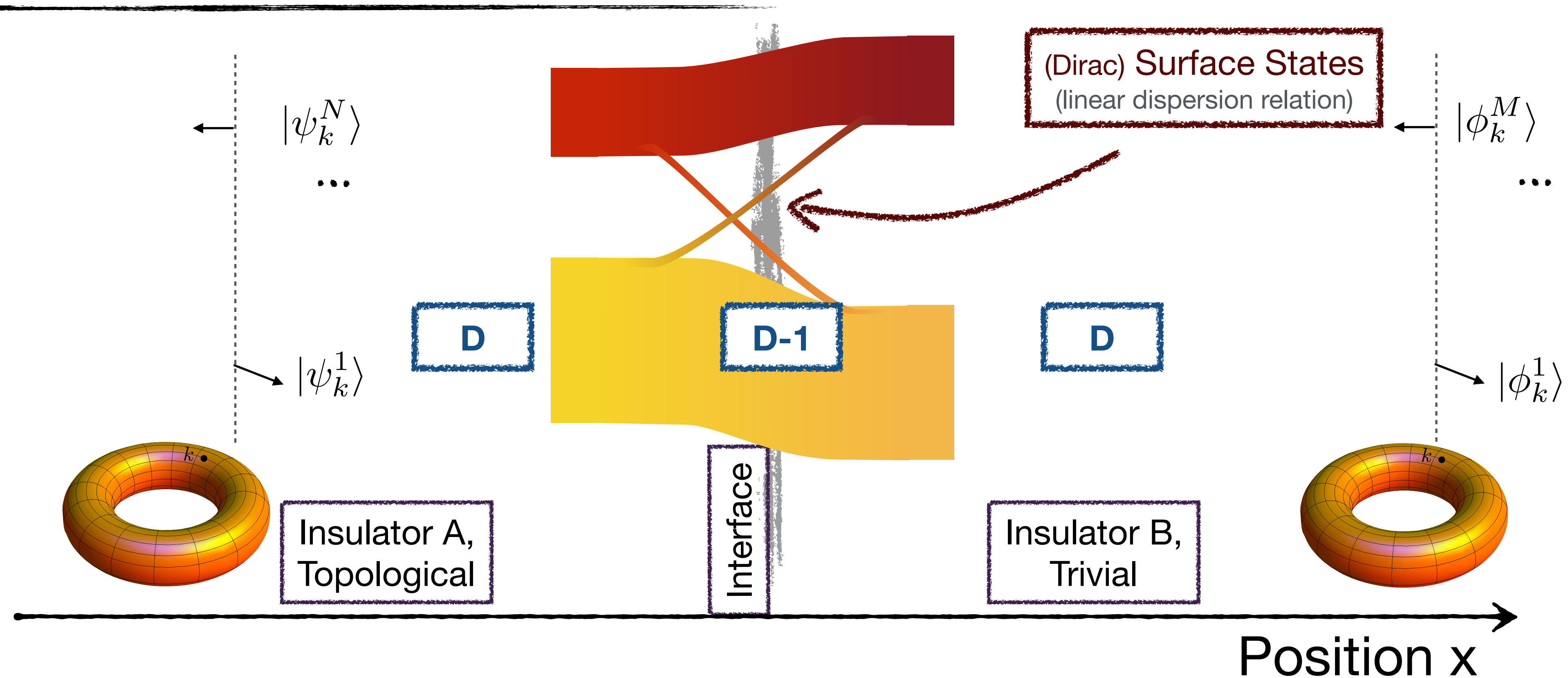
Interface between Insulators



Description of the interface :

- ▶ consider eigenstates $\{|\psi_k^\alpha\rangle\}$ and $\{|\phi_k^\beta\rangle\}$ of materials on both sides
- ▶ extrapolate (position dependent Hamiltonian $H(x)$)

Interface between Insulators

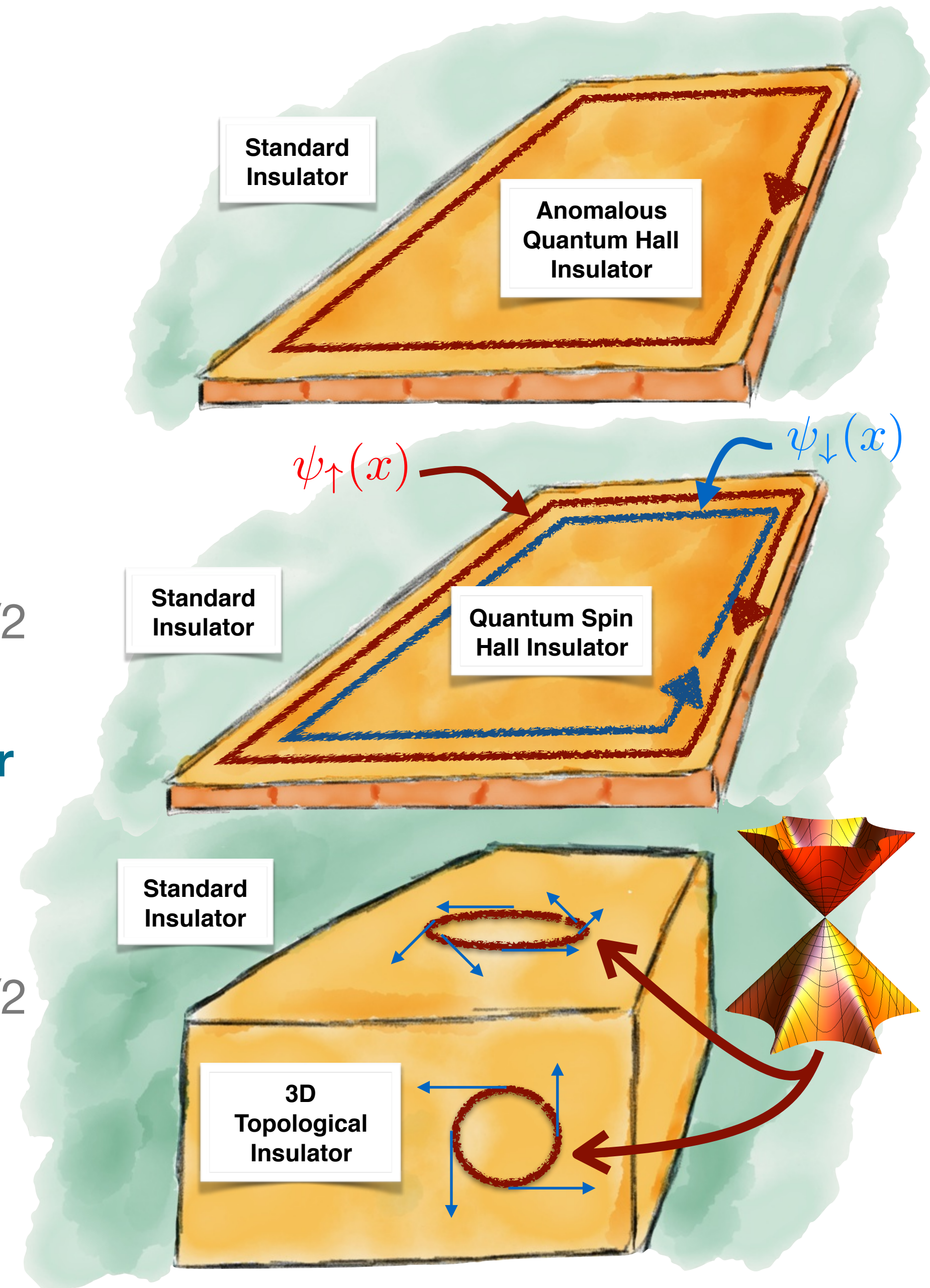


Description of the interface :

- ▶ consider eigenstates $\{|\psi_k^\alpha\rangle\}$ and $\{|\phi_k^\beta\rangle\}$ of materials on both sides
- ▶ extrapolate (position dependent Hamiltonian $H(x)$)

Topological Surface States

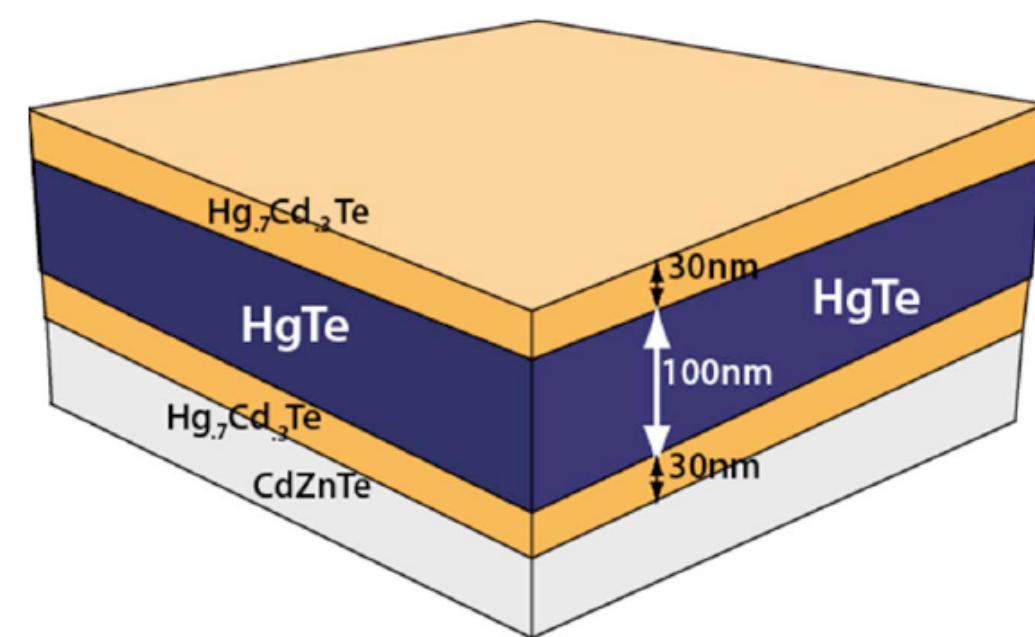
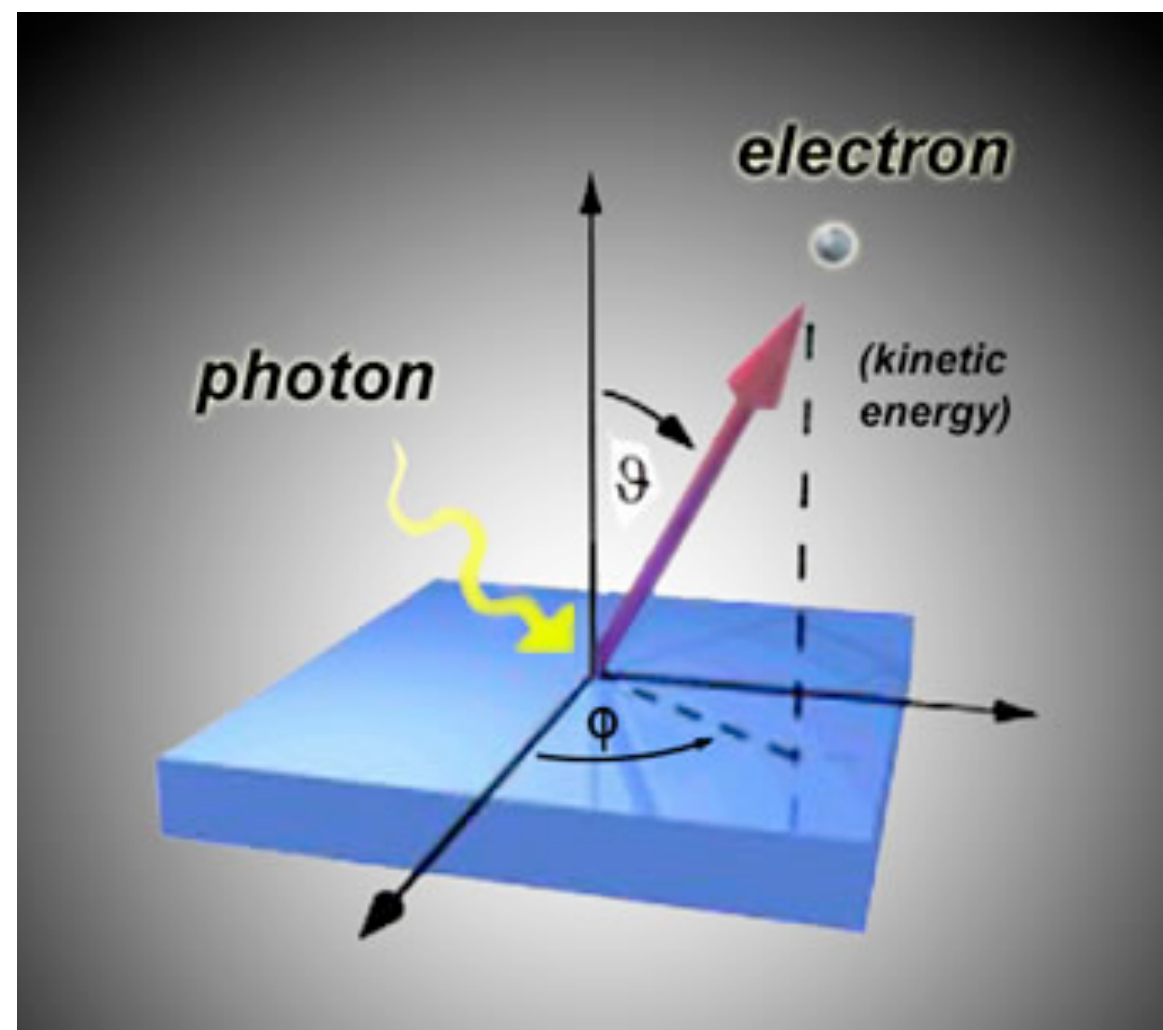
- Quantum Hall Insulator
 - ▶ Two dimensions
 - ▶ breaks Time-Reversal Symmetry
 - ▶ Chern index
 - **Chiral edge states**
- Quantum Spin Hall Insulator
 - ▶ Two dimensions
 - ▶ Time-Reversal Symmetry + spins 1/2
 - ▶ Kane-Mele Z_2 index
 - **Helical edge states : Kramers pair**
- 3D Topological Insulators
 - ▶ Three dimensions
 - ▶ Time-Reversal Symmetry + spins 1/2
 - ▶ Kane-Mele Z_2 index
 - (odd number of) **Dirac cone**



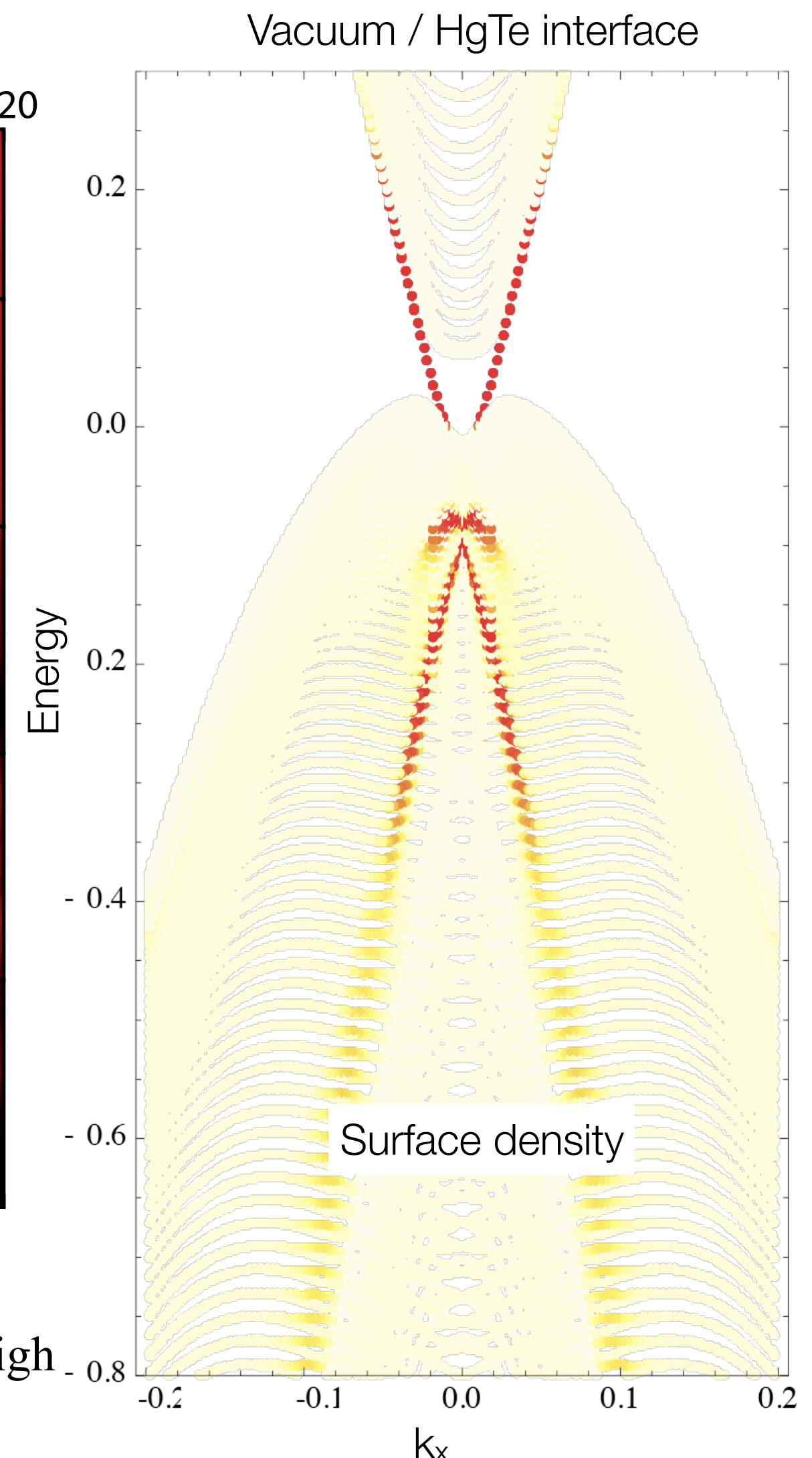
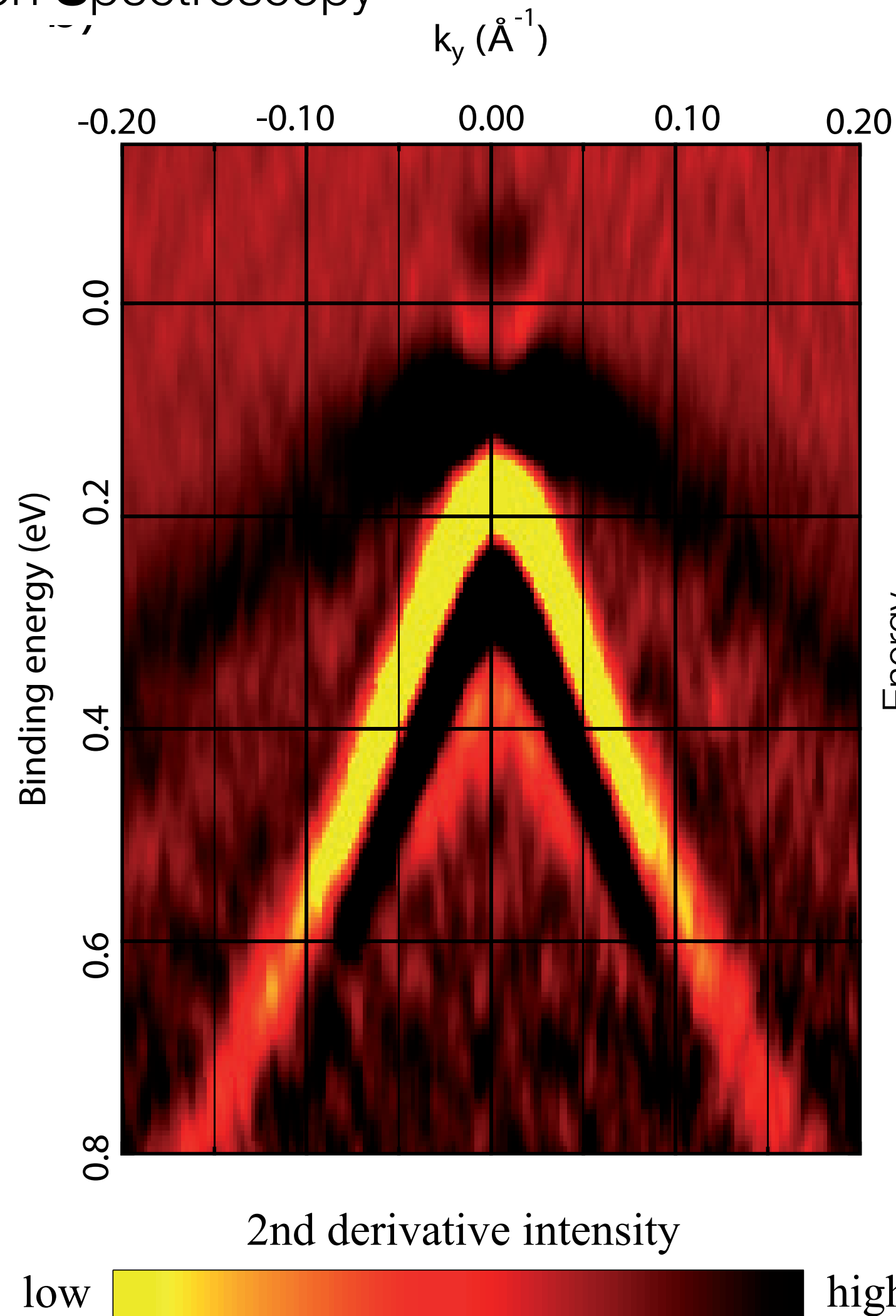
Probing Topological Insulator : Surface States

Coll. Inst. Néel / CEA Leti / ENS-Lyon
O. Crauste, *et al.*, arXiv:1307.2008

Angle Resolved PhotoEmission Spectroscopy



3D strained HgTe

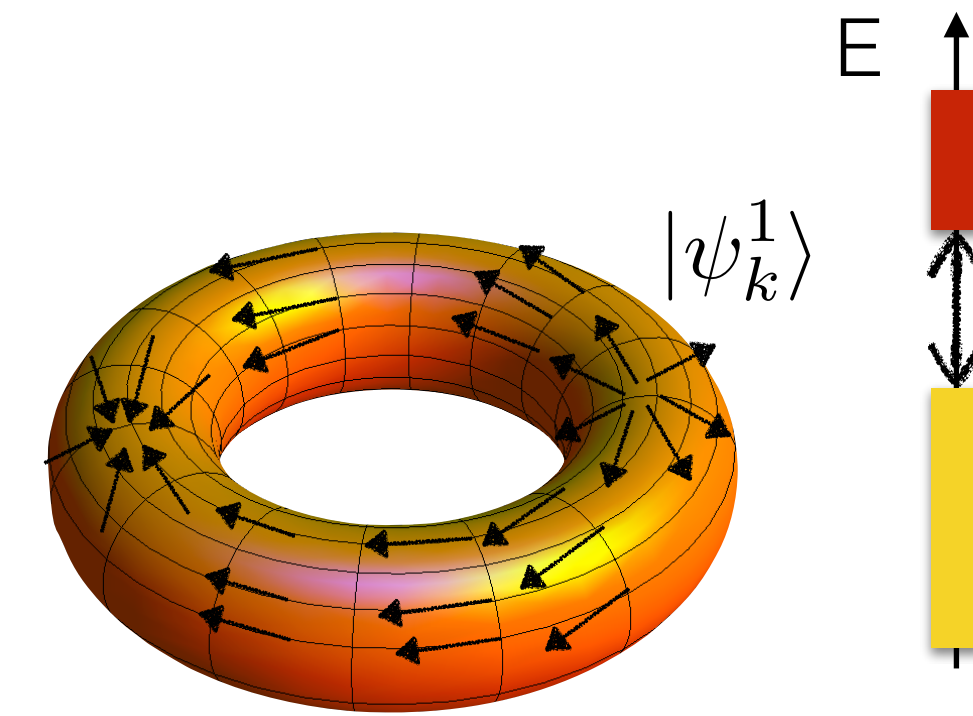


Summary

What is a topological gapped phase ?

► Bulk topological property :

- property of states below the gap
- no continuous Bloch states over Brillouin zone



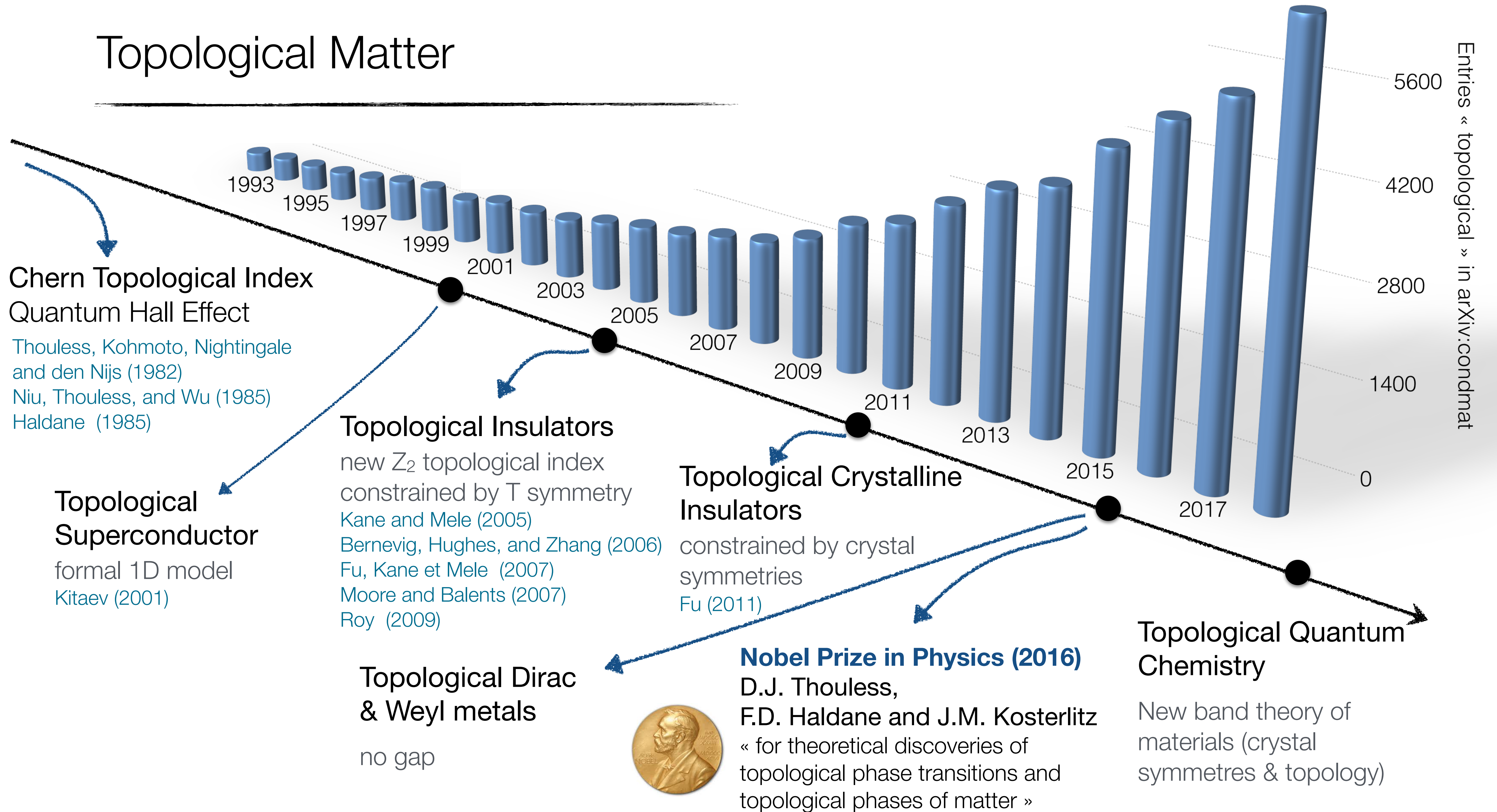
Topological number : Chern number $C_1 = \frac{1}{2\pi} \int_{\text{BZ}} F$

► Surface / edge states (inside gap)

- robust (related to topology)
- unique metals (not conventional)



Topological Matter



Topological Quantum Chemistry

... combining symmetry representations and topology

Catalogue of Topological Electronic Materials

Tiantian Zhang, Yi Jiang, Zhida Song, He Huang, Yuqing He, Zhong Fang, Hongming Weng, Chen Fang
[Nature, 566, 475 \(2019\)](#)

Towards ideal topological materials: Comprehensive database searches using symmetry indicators

Feng Tang, Hoi Chun Po, Ashvin Vishwanath, Xiangang Wan
[Nature 566, 486 \(2019\)](#)

The (High Quality) Topological Materials In The World

M. G. Vergniory, L. Elcoro, C. Felser, B. A. Bernevig, Z. Wang
[Nature 566, 480 \(2019\)](#)

Abstract: « Topological Quantum Chemistry (TQC) links the chemical and symmetry structure of a given material with its topological properties.

Out of **26938** stoichiometric materials in our filtered ICSD database, we find **2861** topological insulators (TI) and **2936** topological semimetals.

Remarkably, our exhaustive results show that a large proportion (~ **24%** !) of all materials in nature are topological

We added an open-source code and end-user button on the Bilbao Crystallographic Server (BCS)

Outline

1. Electronic Properties of Quantum Matter

Topological Insulators



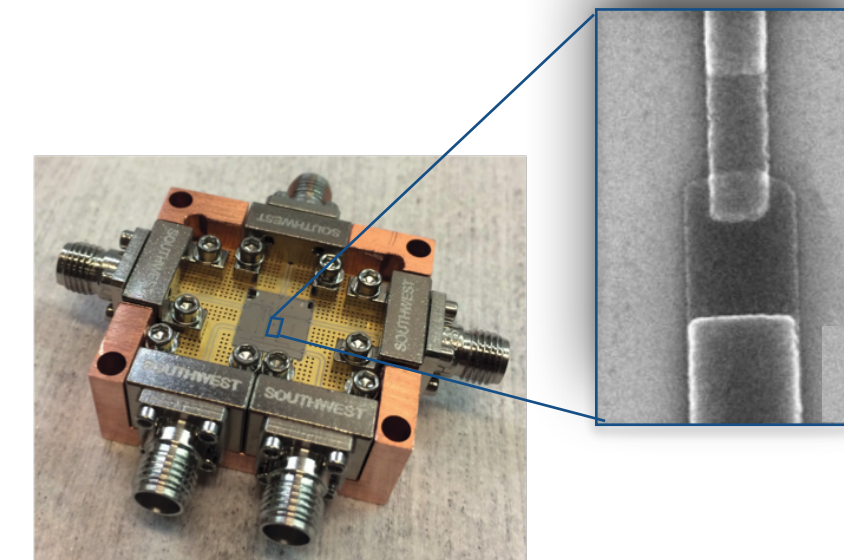
Metals



Insulators

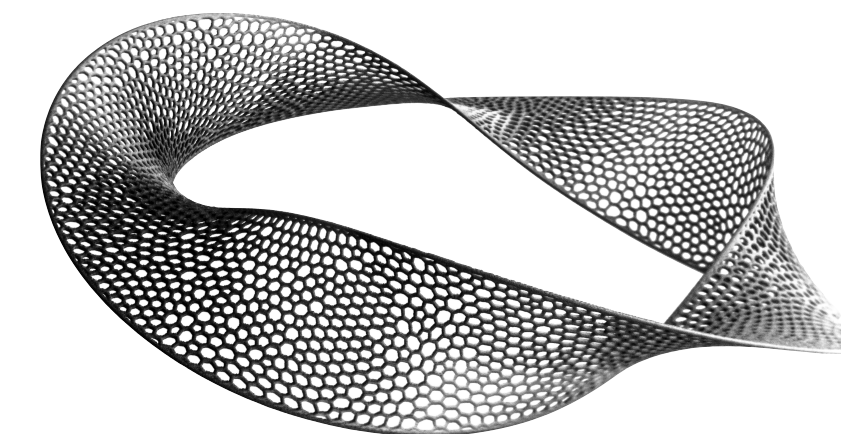
2. Quantum Technologies

Robust quantum devices

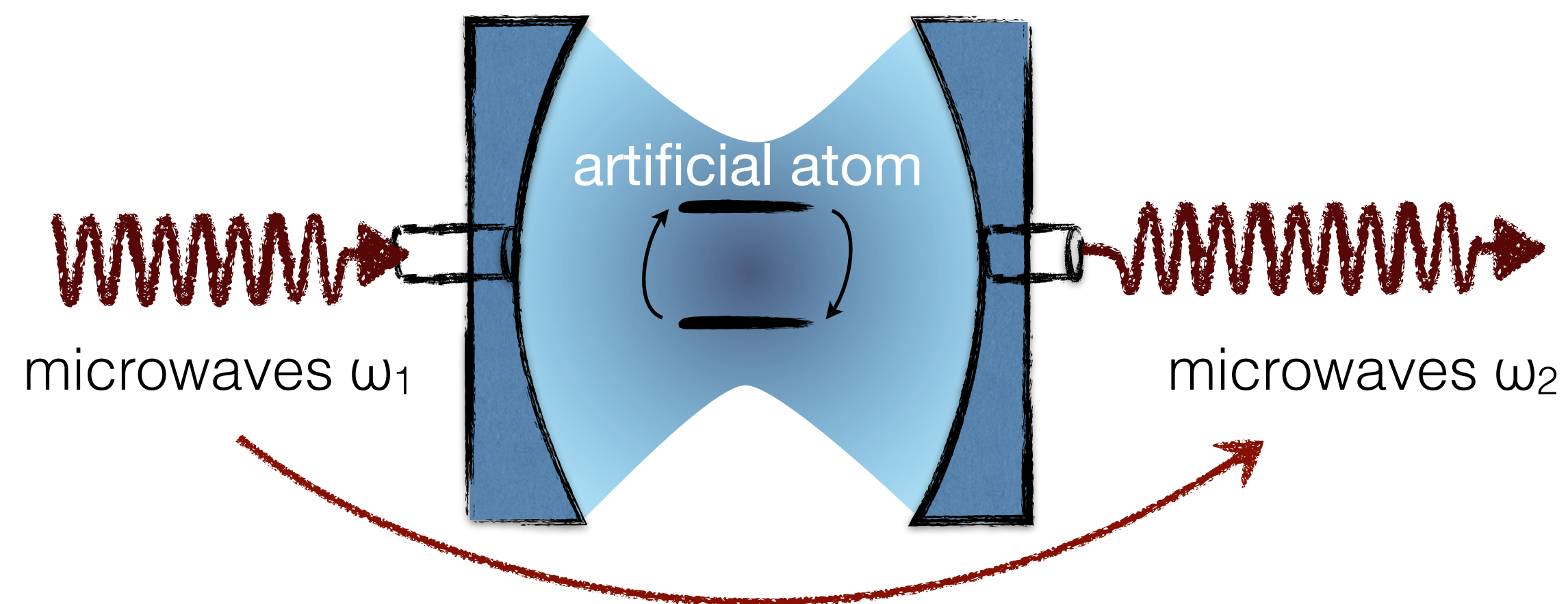


3. Mechanics / Metamaterials

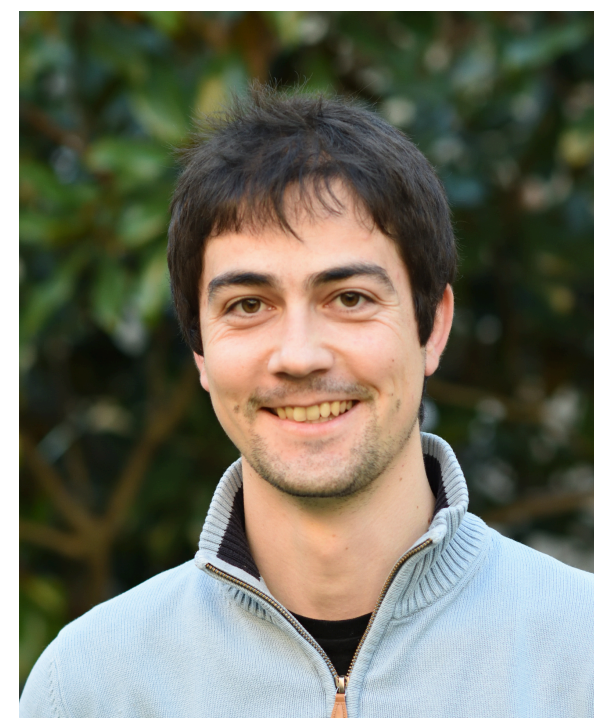
deformations constrained by topology



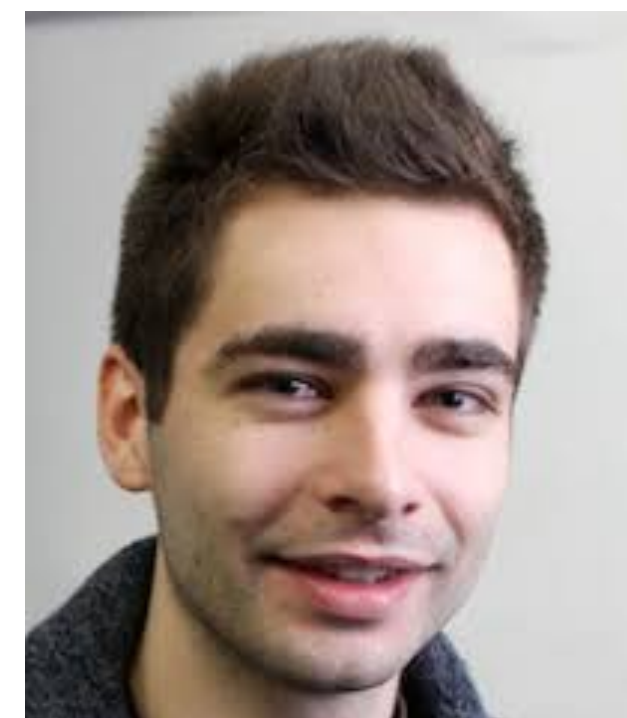
A Topological Energy Pump



Transfer of energy, quantized power
imposed by topology



Clément Dutreix
(Univ. Bordeaux)



Quentin Ficheux
(Univ. Maryland)



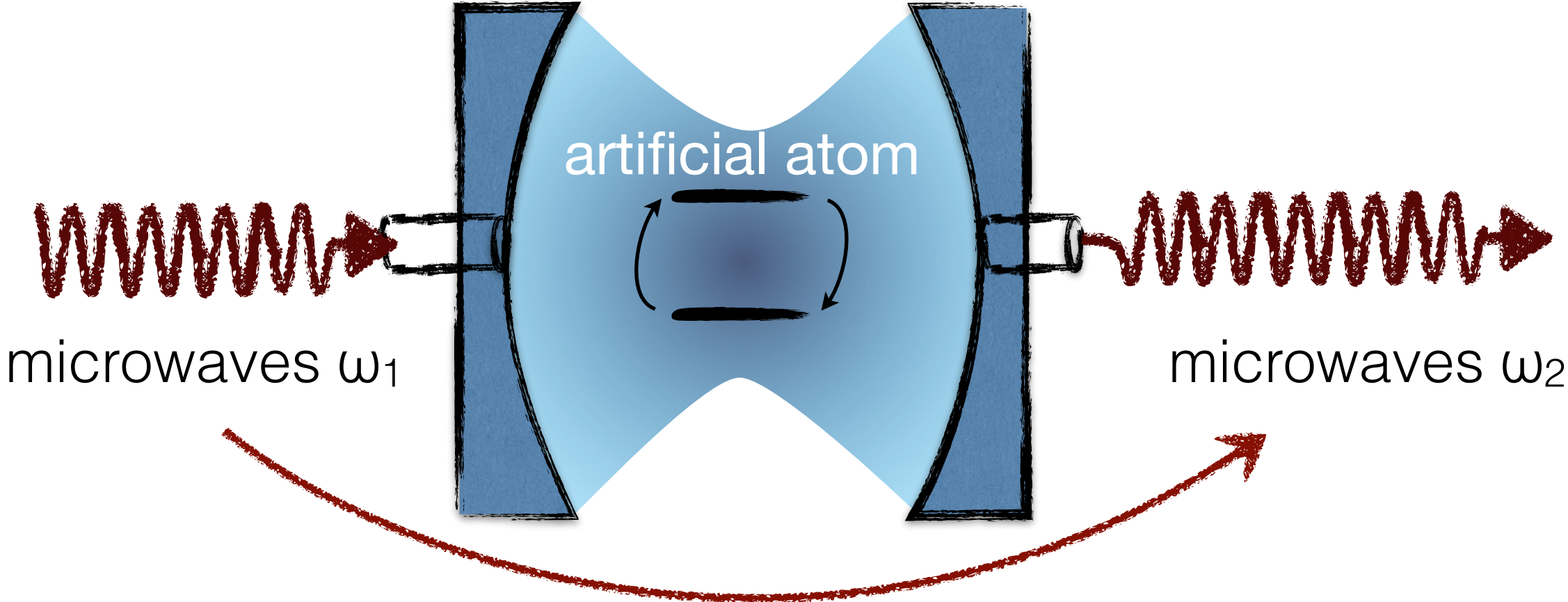
Pierre Delplace
(ENS Lyon)



Benjamin Huard
(ENS Lyon)

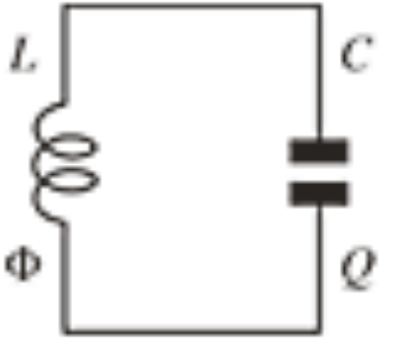
C. Dutreix et al., in preparation (2019)

A Topological Energy Pump

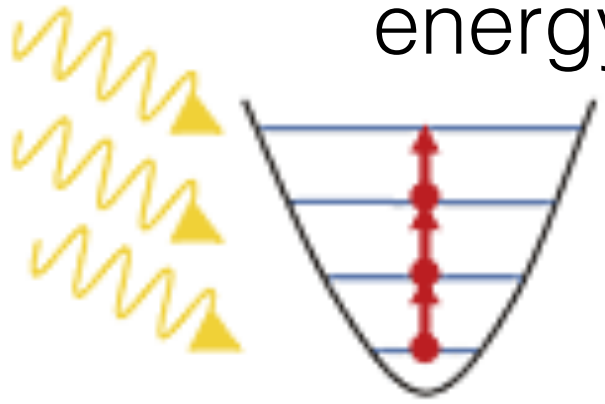


Transfer of energy, quantized power imposed by topology

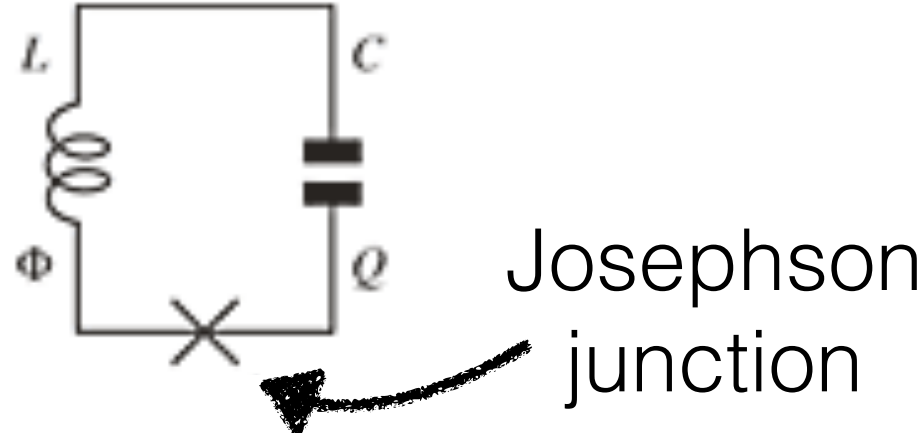
LC circuit



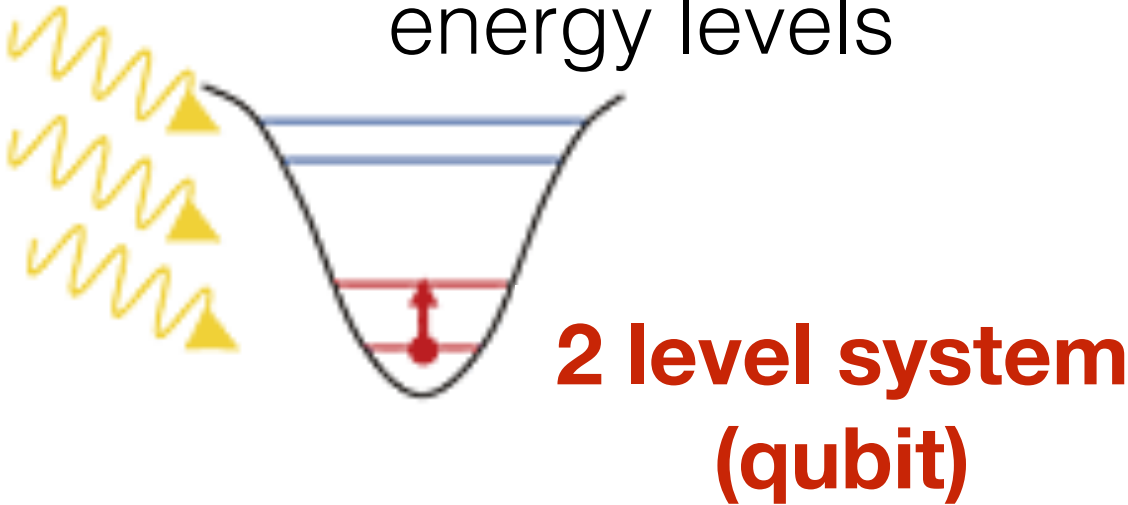
Equally spaced energy levels



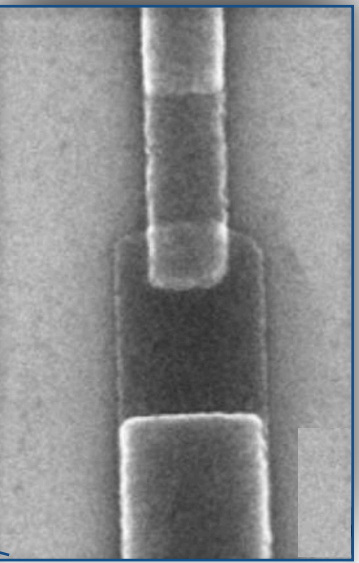
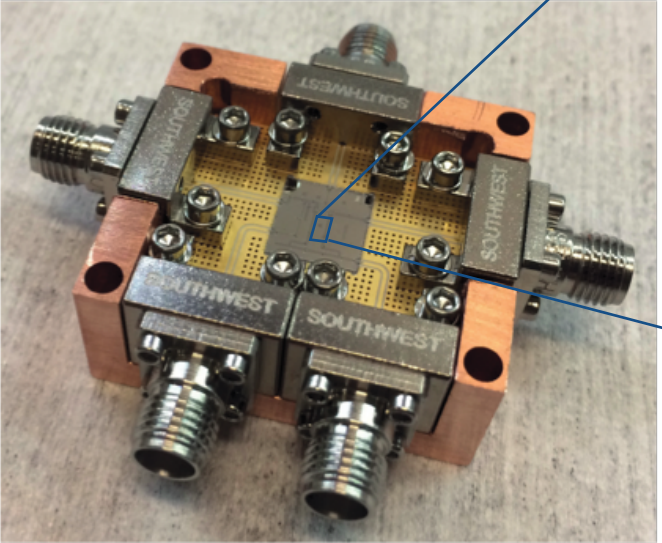
LC circuit with Josephson junction



Unequally spaced energy levels

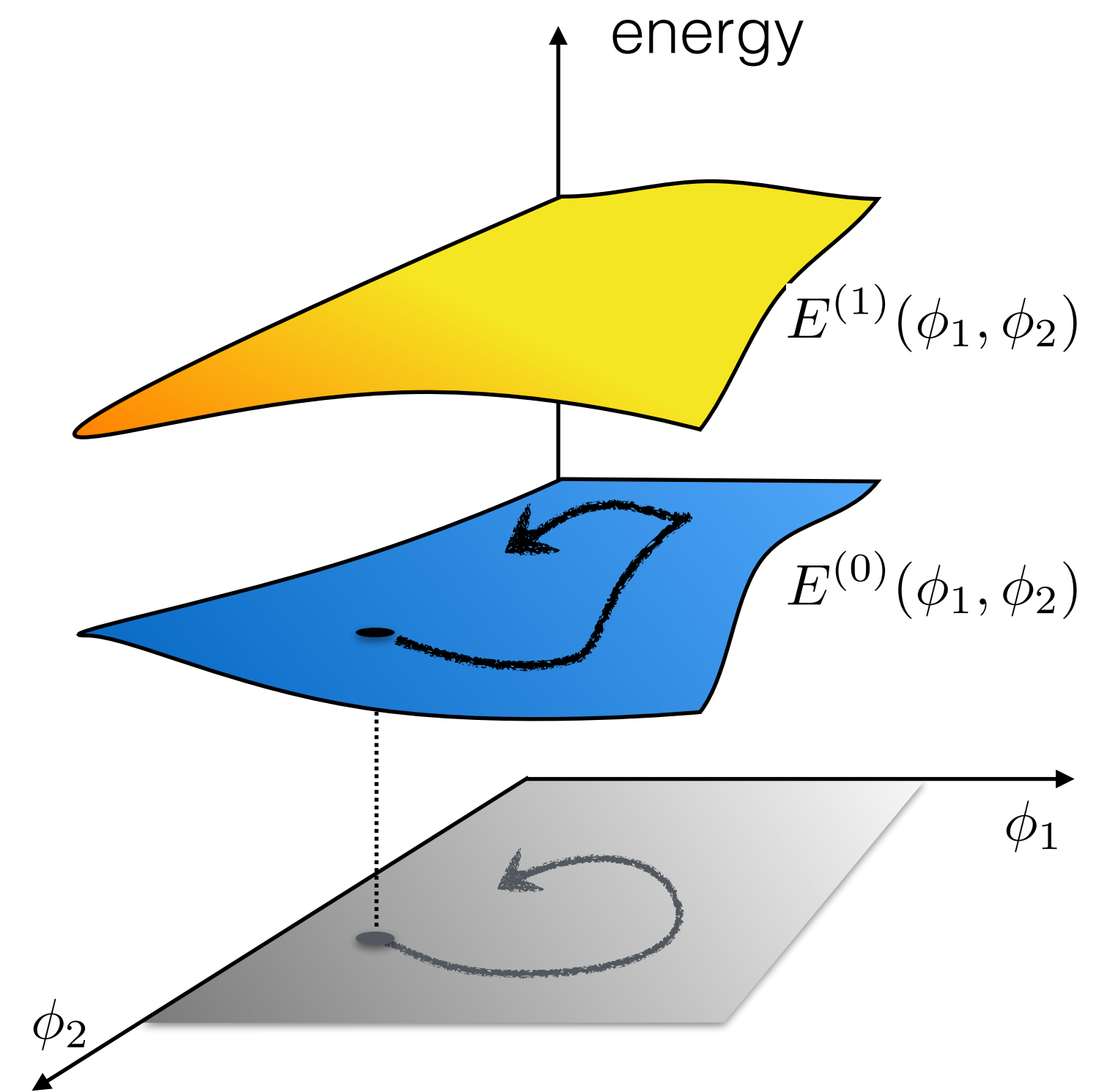
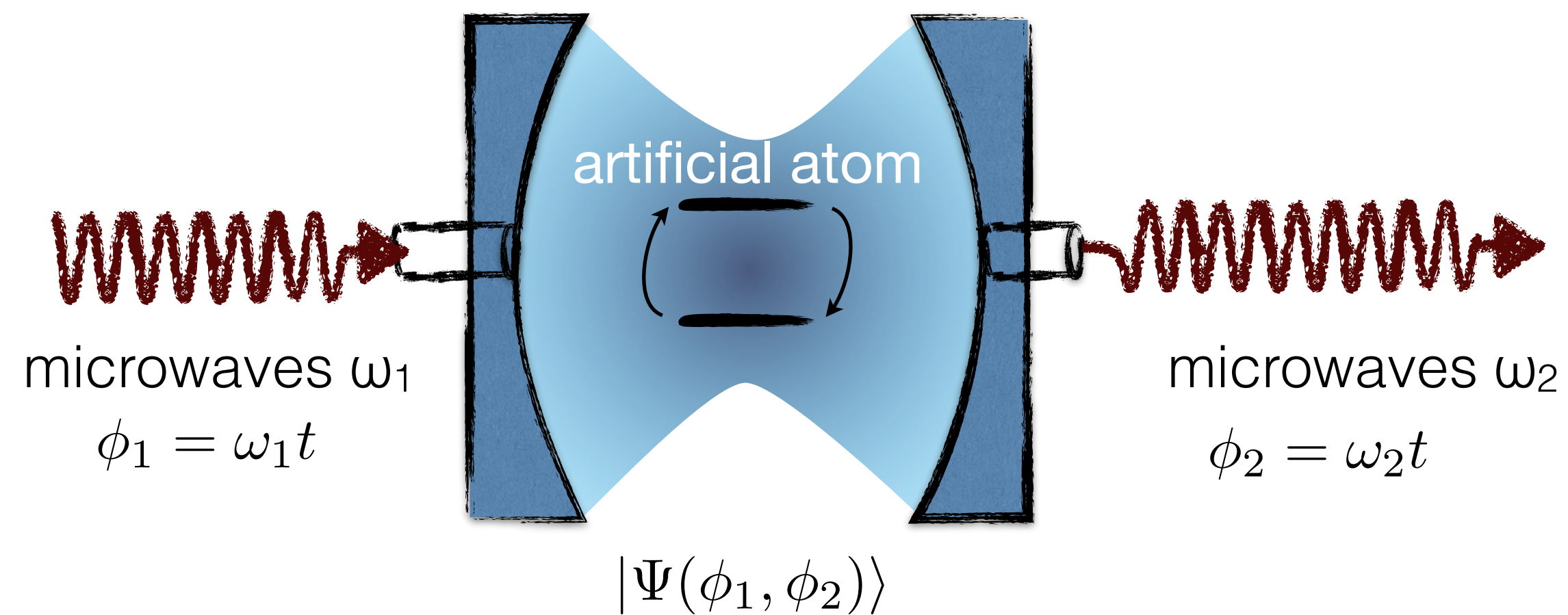


Quantum Circuit



Josephson Junction

A Topological Energy Pump



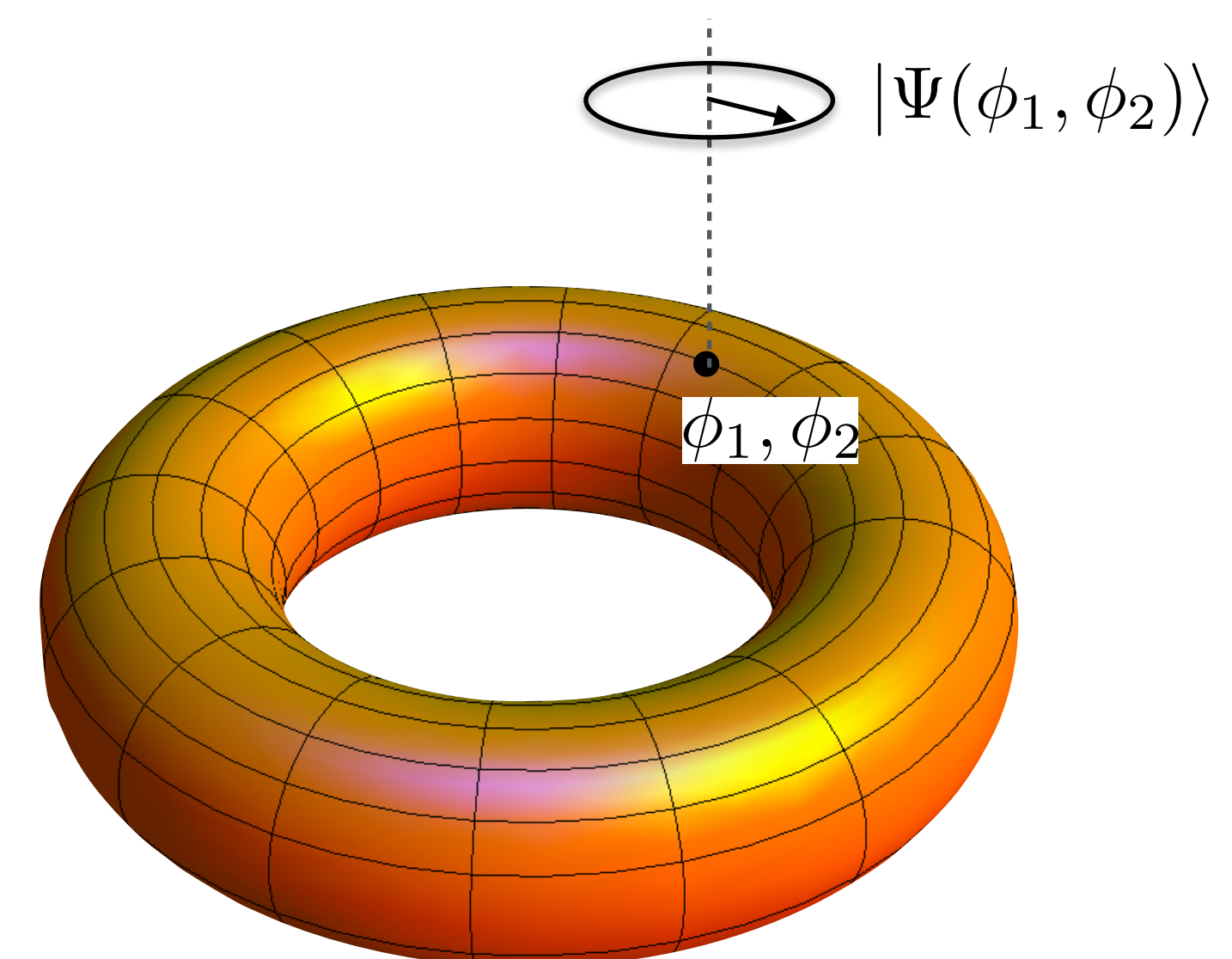
Periodicity

- Hamiltonian $H(\phi_1, \phi_2)$ is 2π periodic in ϕ_1, ϕ_2

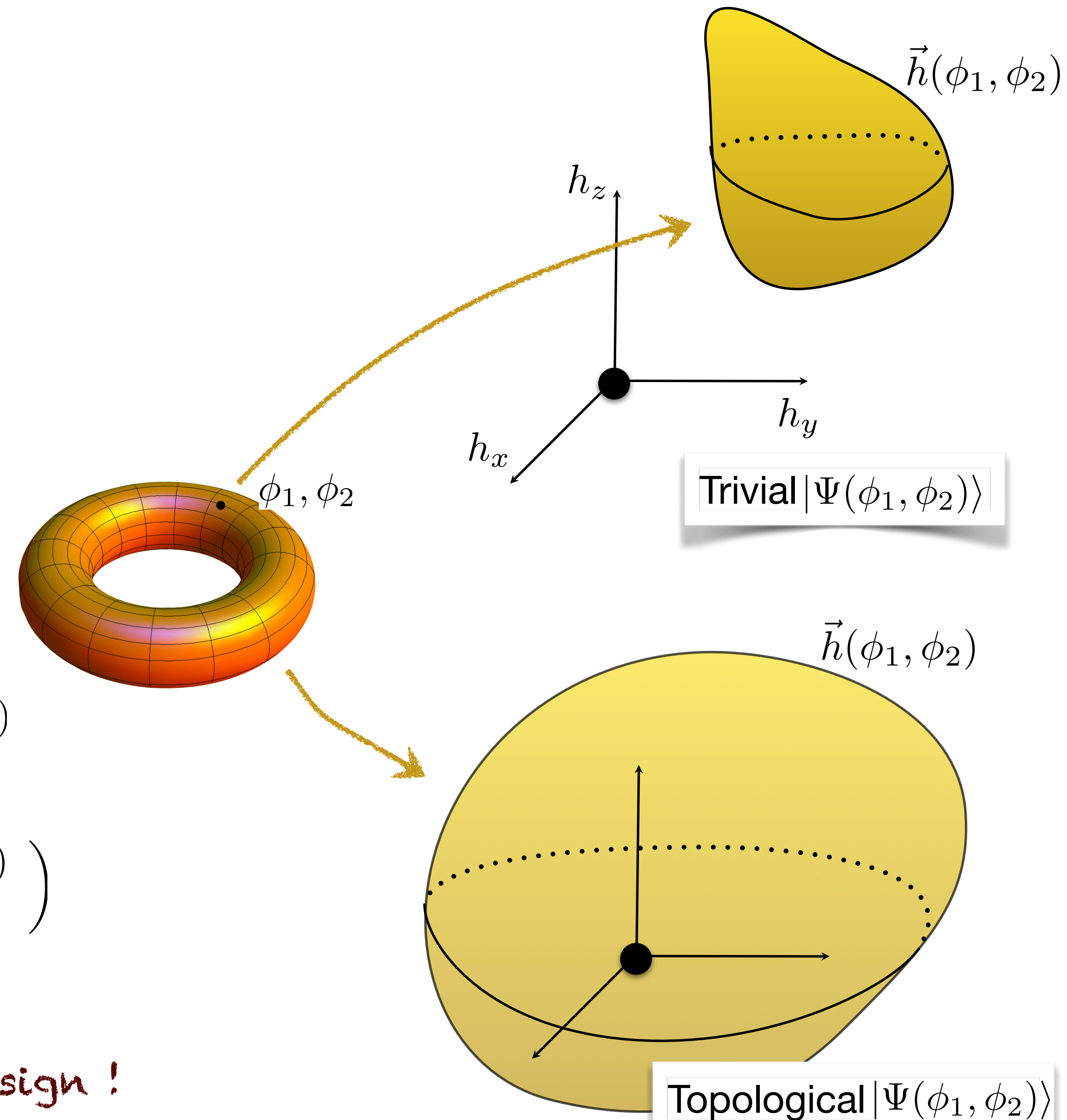
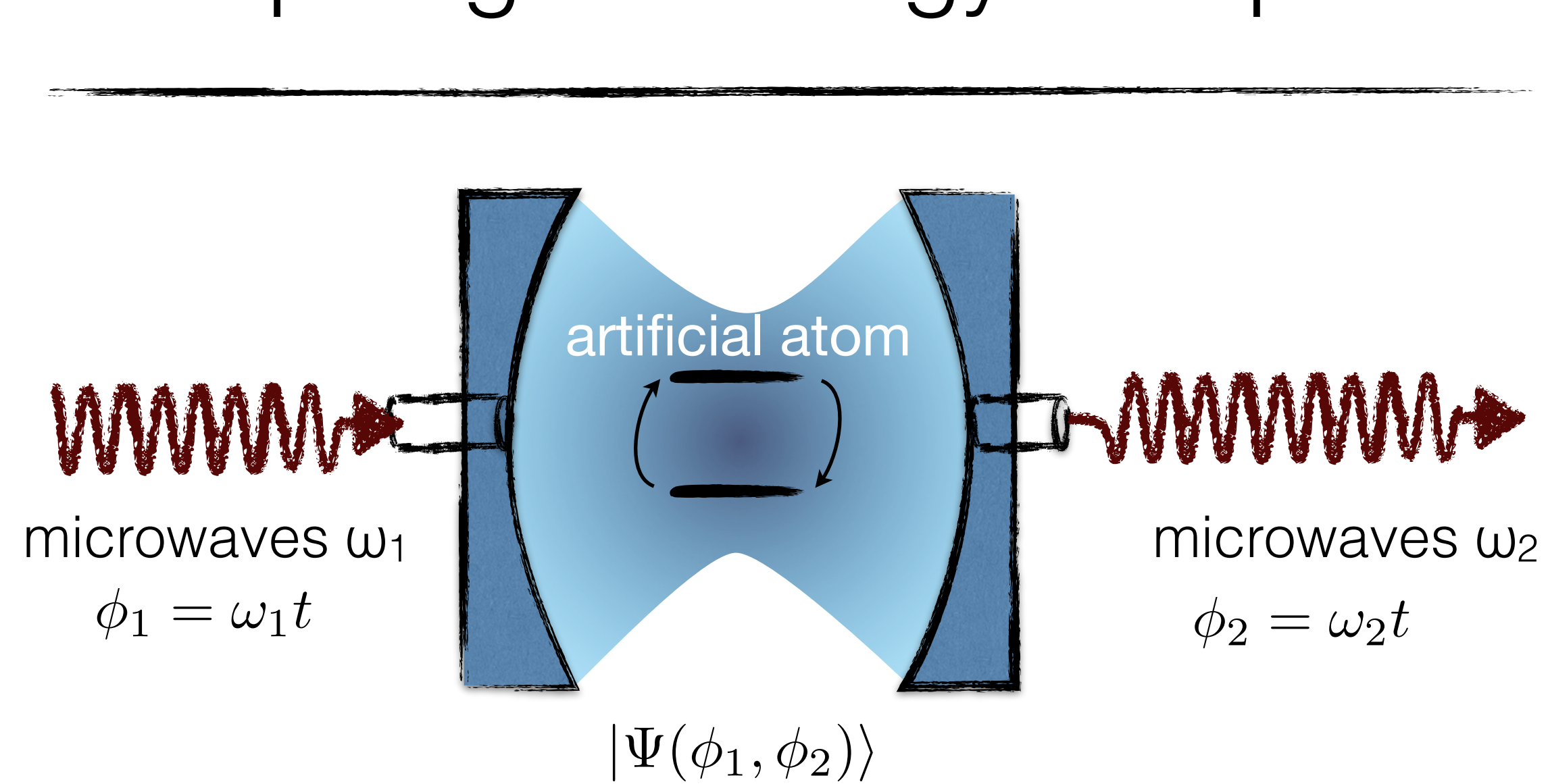
Adiabatic dynamics

- gap in energy for all ϕ_1, ϕ_2
- quantum system remains in 1 (ground) state

→ topological (driven) quantum state ?



A Topological Energy Pump



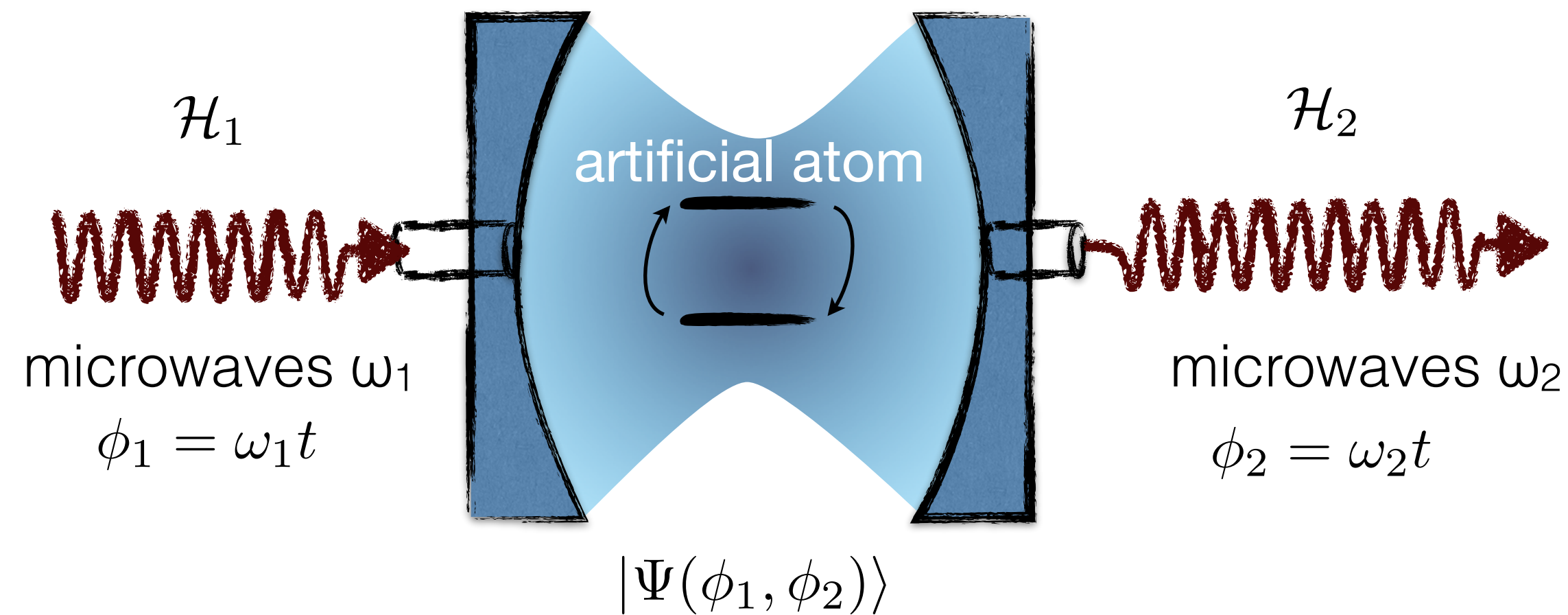
2-level system : recipe for topology

- general parametrization in terms of vector $\vec{h} = (h_x, h_y, h_z)$
- driven 2-level system : map $\vec{h}(\phi_1, \phi_2)$

$$H(\phi_1, \phi_2) = \vec{h} \begin{pmatrix} h_z(\phi_1, \phi_2) & h_x(\phi_1, \phi_2) - ih_y(\phi_1, \phi_2) \\ h_x(\phi_1, \phi_2) + ih_y(\phi_1, \phi_2) & -h_z(\phi_1, \phi_2) \end{pmatrix}$$

necessary condition : all couplings must change sign !

A Topological Energy Pump



Chern number (topological index)

$$c_{12}^{(\Psi)} = \frac{1}{2\pi} \int d\phi_1 d\phi_2 F_{1,2}^{(\Psi)}$$

→ Average Berry curvature

$$\overline{F_{1,2}^{(\Psi)}} = \frac{1}{2\pi} c_{12}^{(\Psi)}$$

Berry curvature of the eigenstate $|\Psi(\phi_1, \phi_2)\rangle$

General Formalism for topological pumping

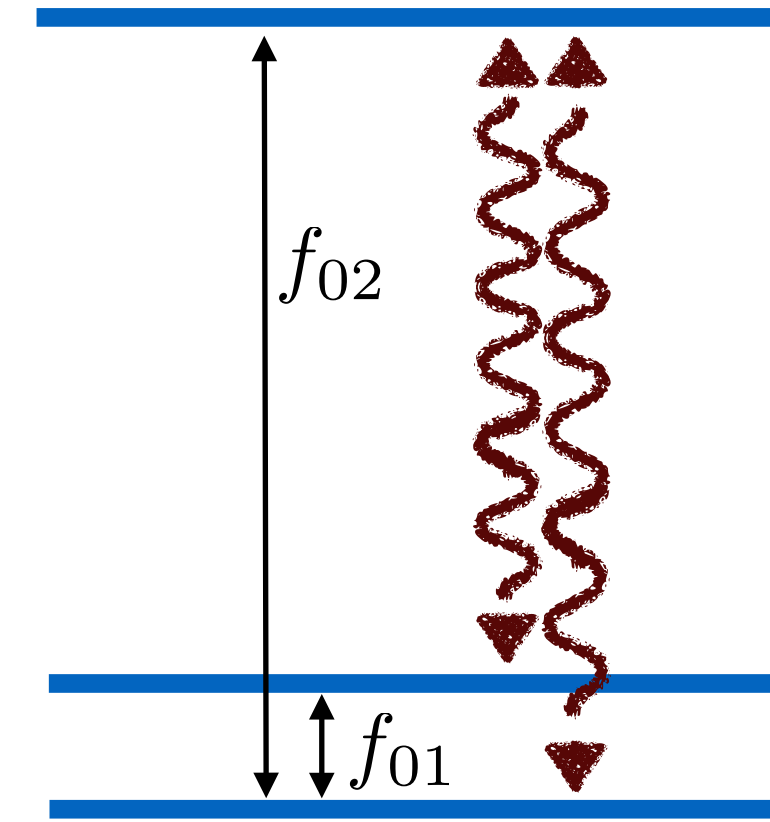
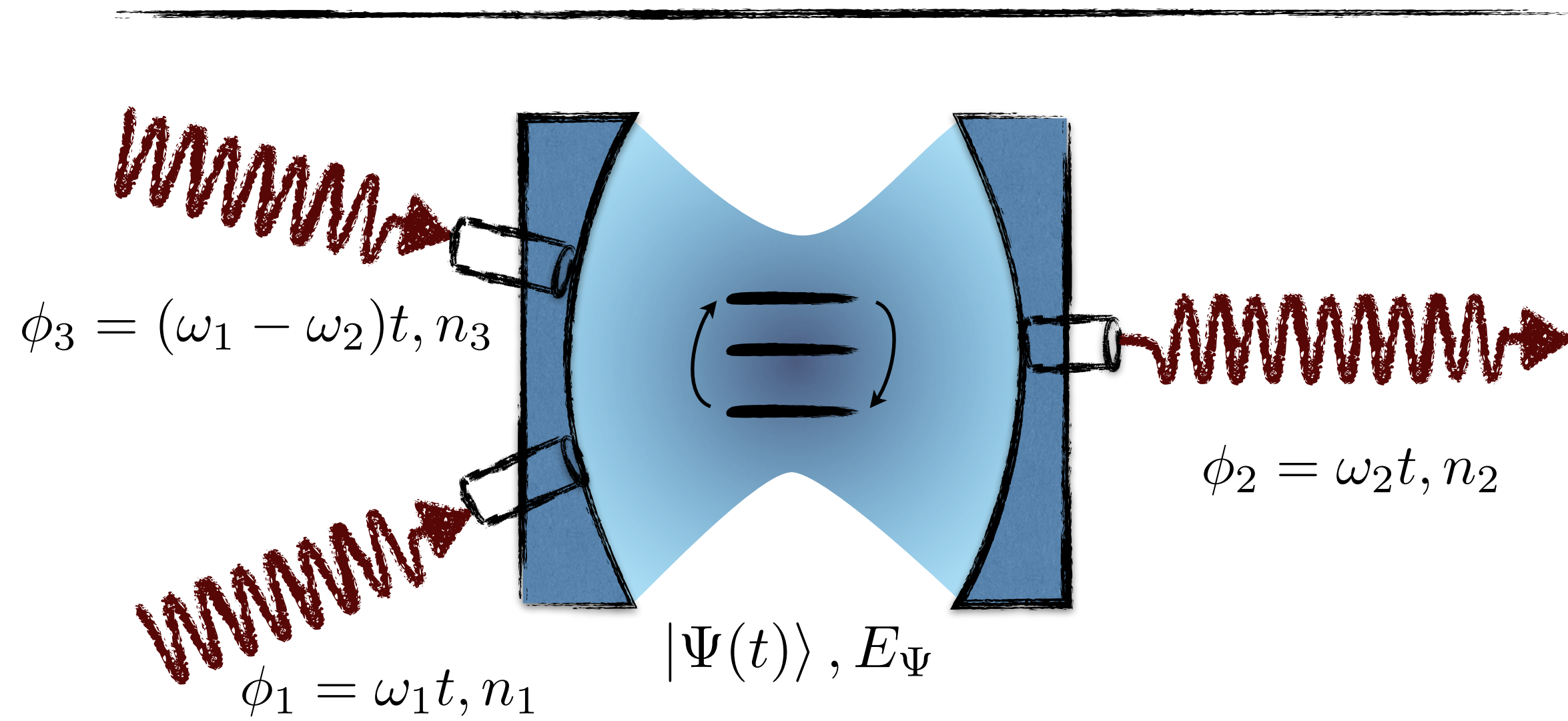
- each bath: conjugate classical variables n_α, ϕ_α , Hamiltonian \mathcal{H}_α
- equation of motion:

$$\dot{n}_\alpha = -\frac{\partial \mathcal{H}_\alpha}{\partial \phi_\alpha} - \langle \Psi(t) | \frac{\partial H}{\partial \phi_\alpha} | \Psi(t) \rangle = -\frac{\partial \mathcal{H}_\alpha}{\partial \phi_\alpha} - \frac{\partial E_\Psi}{\partial \phi_\alpha} - \hbar \sum_{\beta} F_{\alpha,\beta}^{(\Psi)} \dot{\phi}_\beta$$

$$\dot{\phi}_\alpha = \frac{\partial \mathcal{H}_\alpha}{\partial n_\alpha} = \omega_\alpha$$

- power transfer $\Delta \mathcal{E}_1 = \dot{n}_1 \omega_1 = \frac{\hbar}{2\pi} c_{12} \omega_1 \omega_2$

A Topological Energy Pump : effective 2-level system (qutrit)



ω_{01}	$2\pi \times 200$ MHz
ω_{02}	$2\pi \times 5$ GHz
Ω_{01}	$2\pi \times 100$ MHz
Ω_{02}	$2\pi \times 1$ GHz
Ω_{12}	$2\pi \times 1.2$ GHz
$\dot{\phi}_a = 2\dot{\phi}_b$	$2\pi \times 20$ MHz

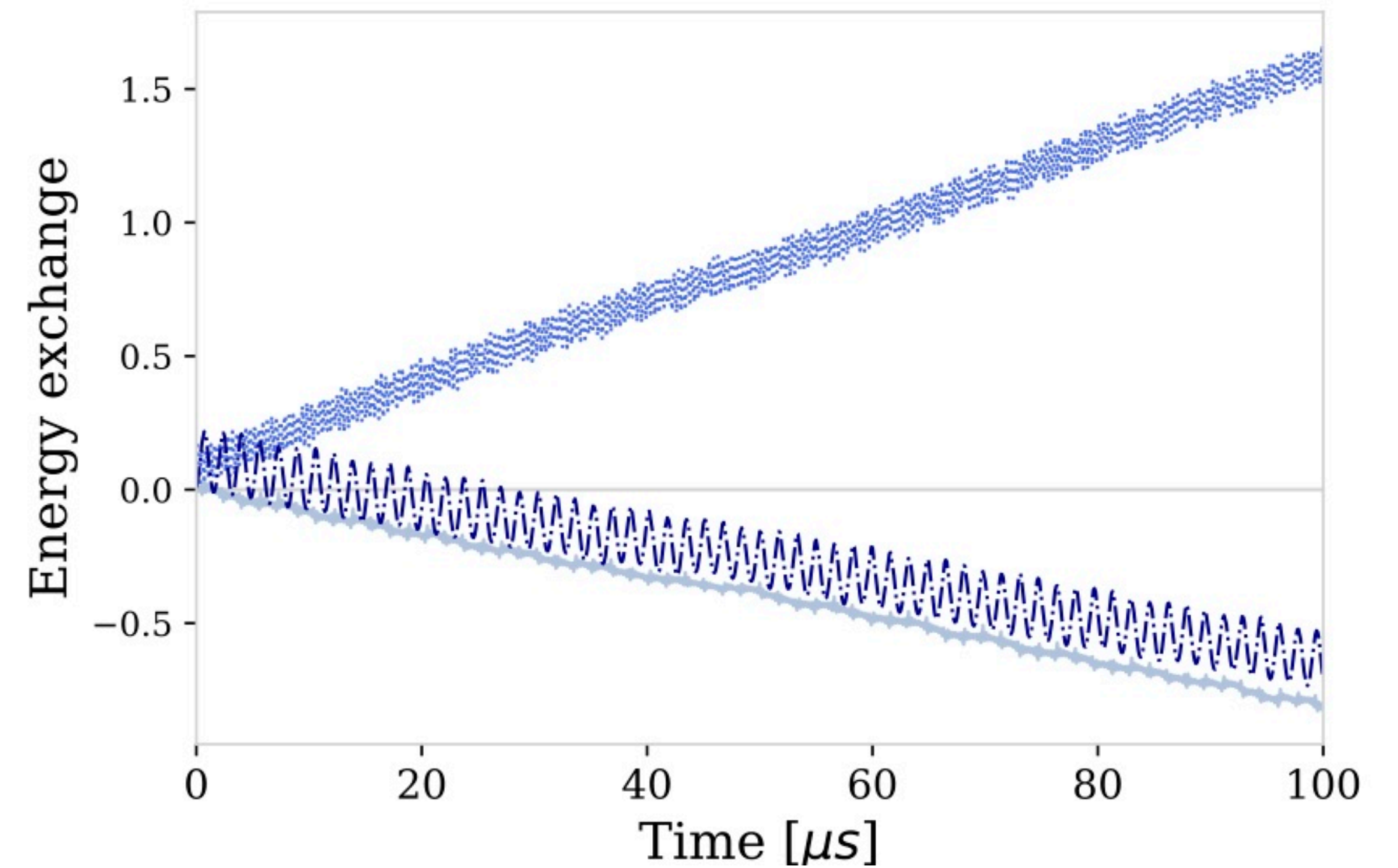
Topological transfer :

$$\dot{n}_I = \dot{n}_1 + \dot{n}_3 = \hbar F^{12} \omega_2$$

$$\dot{n}_{II} = \dot{n}_2 - \dot{n}_3 = -\hbar F^{12} \omega_1$$

Energy exchange : $\Delta \mathcal{E}_i = \dot{n}_i \omega_i \implies \Delta \mathcal{E}_1 = -\Delta \mathcal{E}_2$

Topological rate : $\frac{\Delta \mathcal{E}_1}{\omega_1} + \frac{\Delta \mathcal{E}_3}{\omega_1 - \omega_2} = \frac{\hbar}{2\pi} c_{12} \omega_2$



Elasticity of non-orientable objects

Non-orientable object

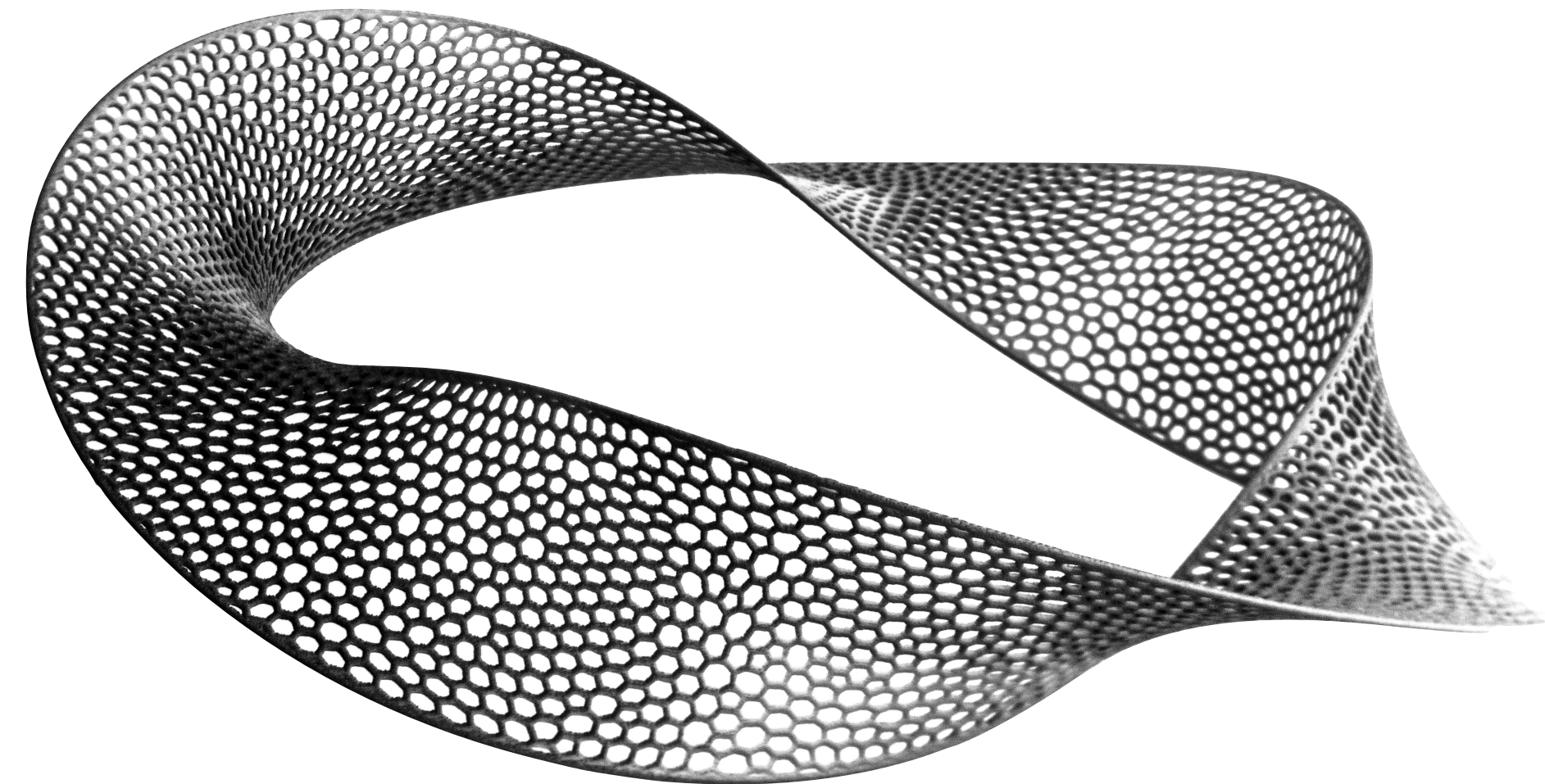
- ▶ cannot define normal vector to surface continuously
- ▶ topological property
- ▶ simplest example : Möbius strip (3D printed)



Denis Bartolo
(ENS Lyon)



Marcelo Guzman
(ENS Lyon)

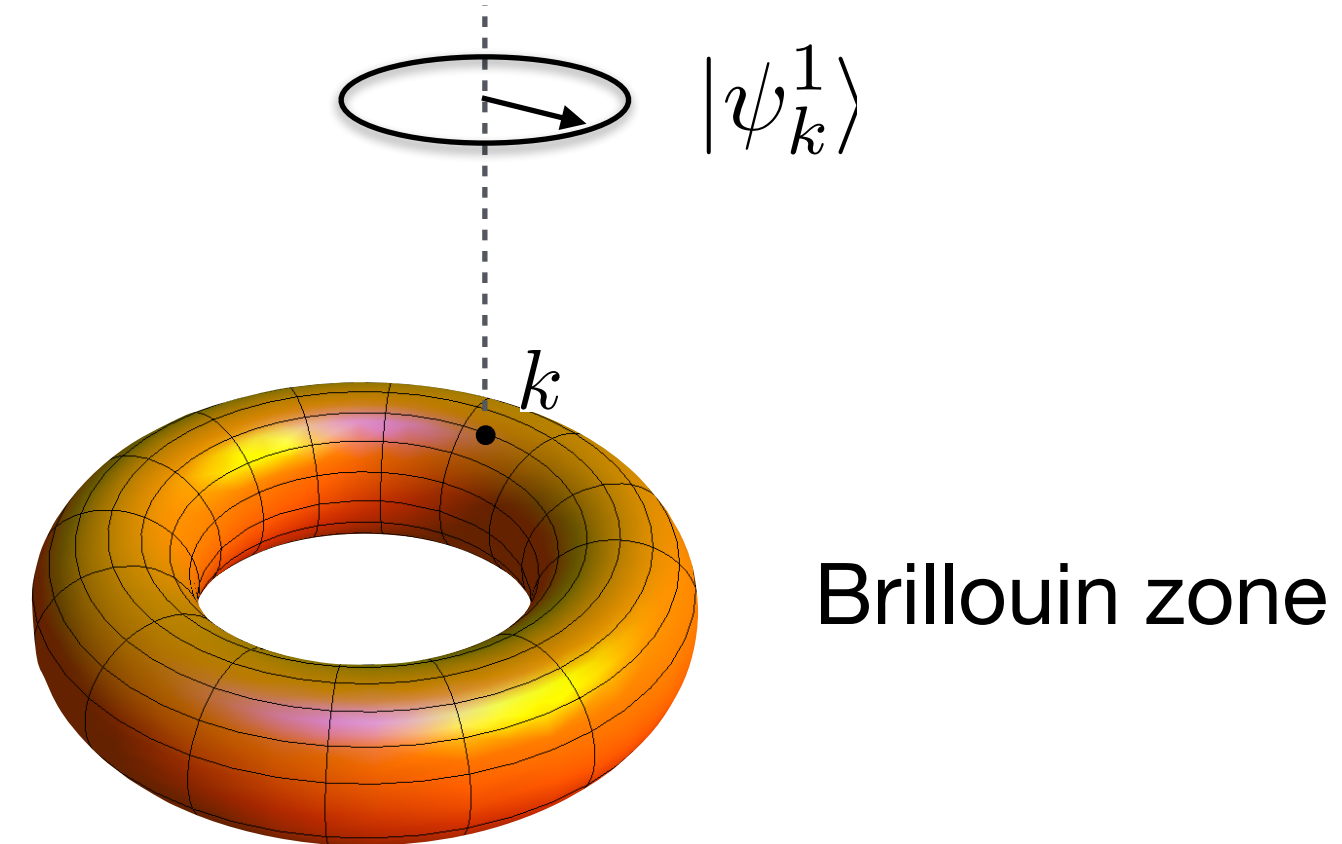
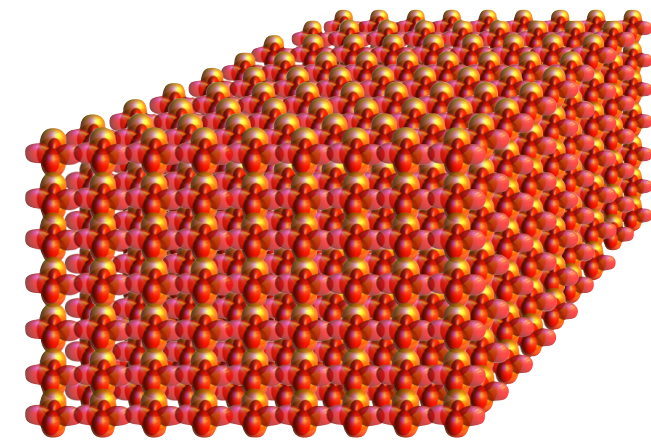


Does the non-orientability (topology) of a Möbius strip manifest itself in its mechanical response ?

Elasticity of non-orientable objects

1. Electronic Properties of Quantum Matter

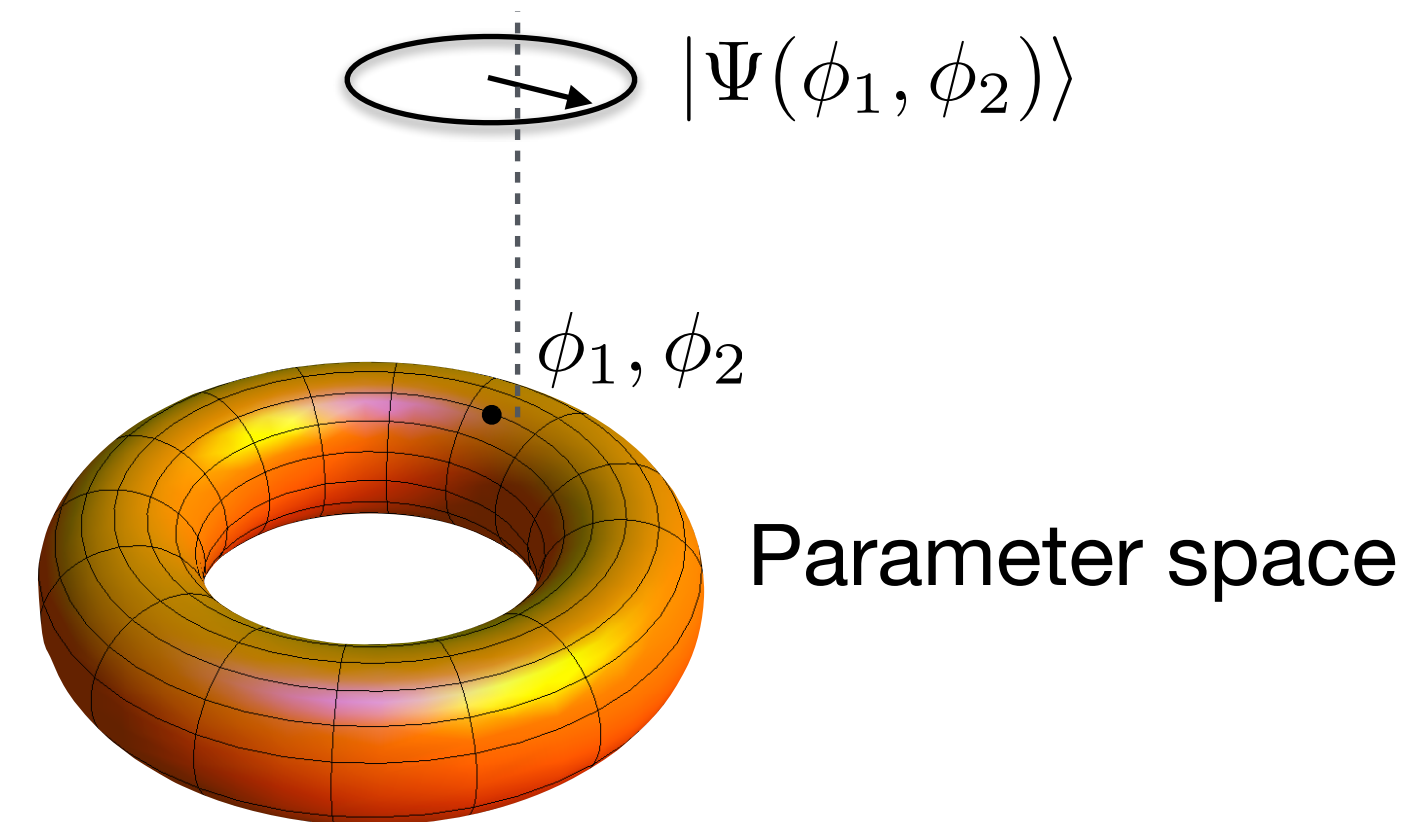
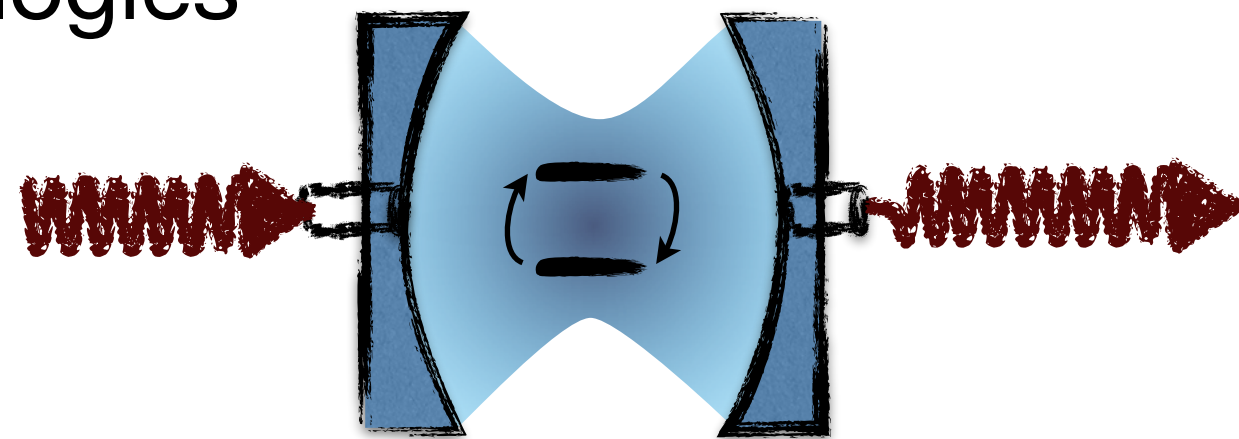
Topological Insulators



Brillouin zone

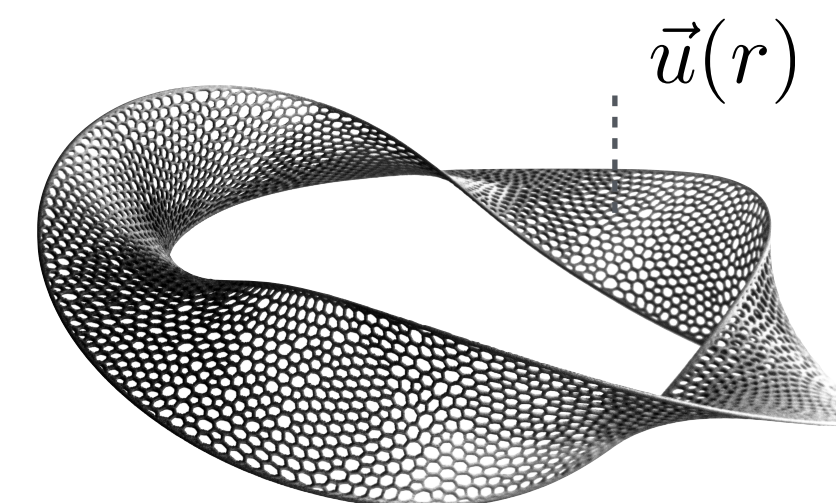
2. Quantum Technologies

Topological pump



Parameter space

3. Mechanics / Metamaterials



$\vec{u}(r)$

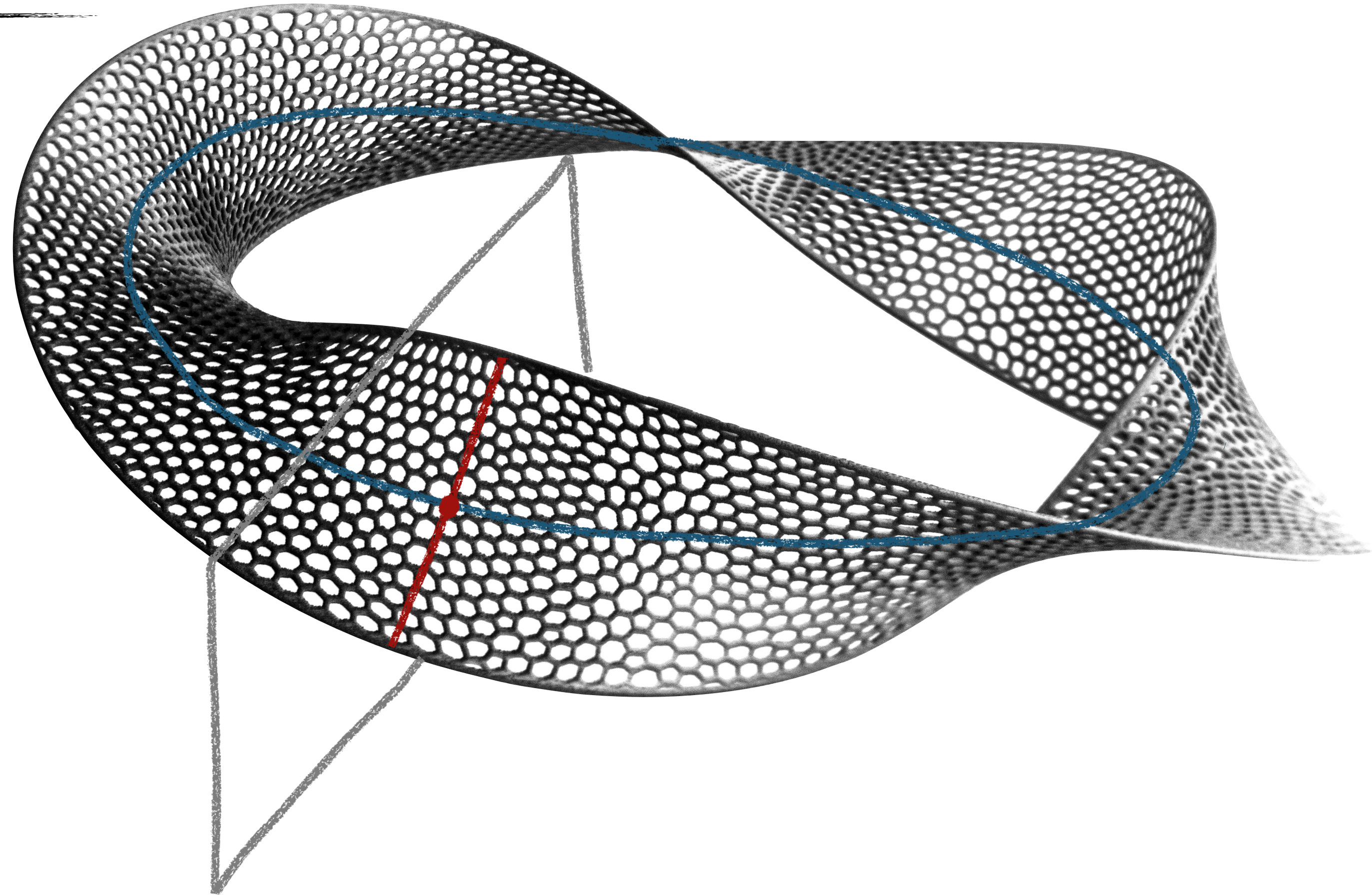
topology by
« engineering »
the vectors

topology by
« engineering »
the base-space

Elasticity of a Möbius strip

Definition of the shape of the strip

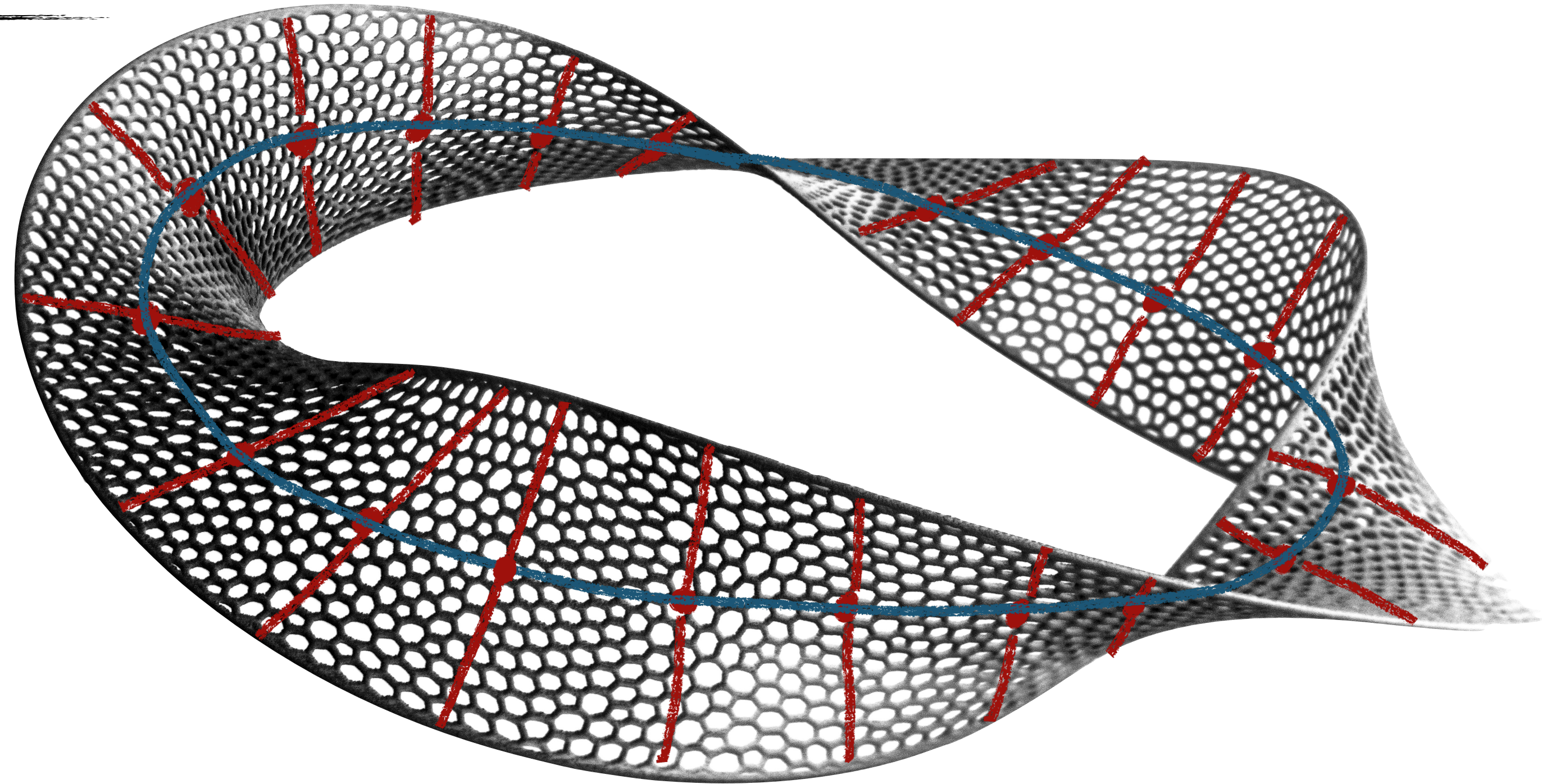
- ▶ discrete version of the strip



Elasticity of a Möbius strip

Definition of the shape of the strip

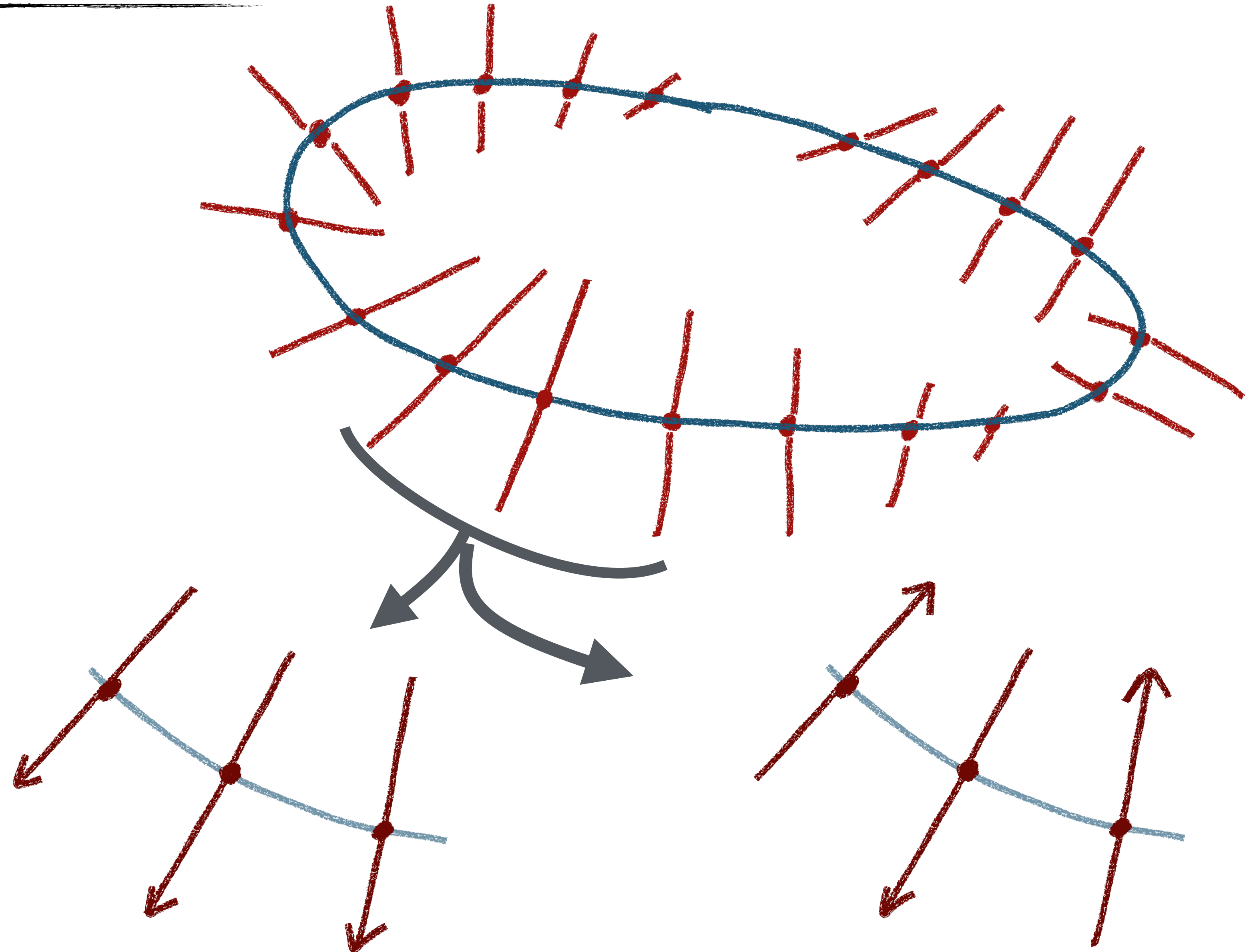
- ▶ discrete version of the strip



Elasticity of a Möbius strip

Definition of the shape of the strip

- ▶ discrete version of the strip
- ▶ lattice of « directions »
- ▶ but elasticity requires vectors !

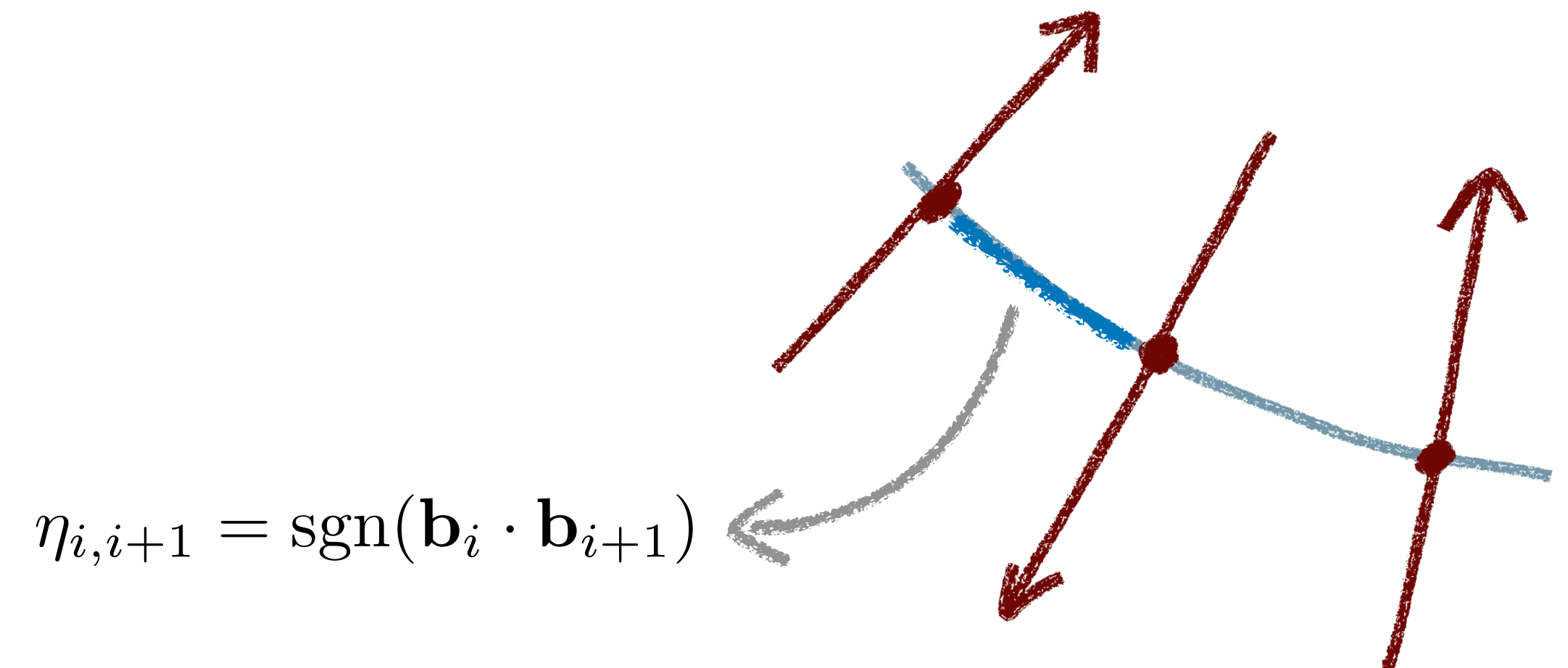
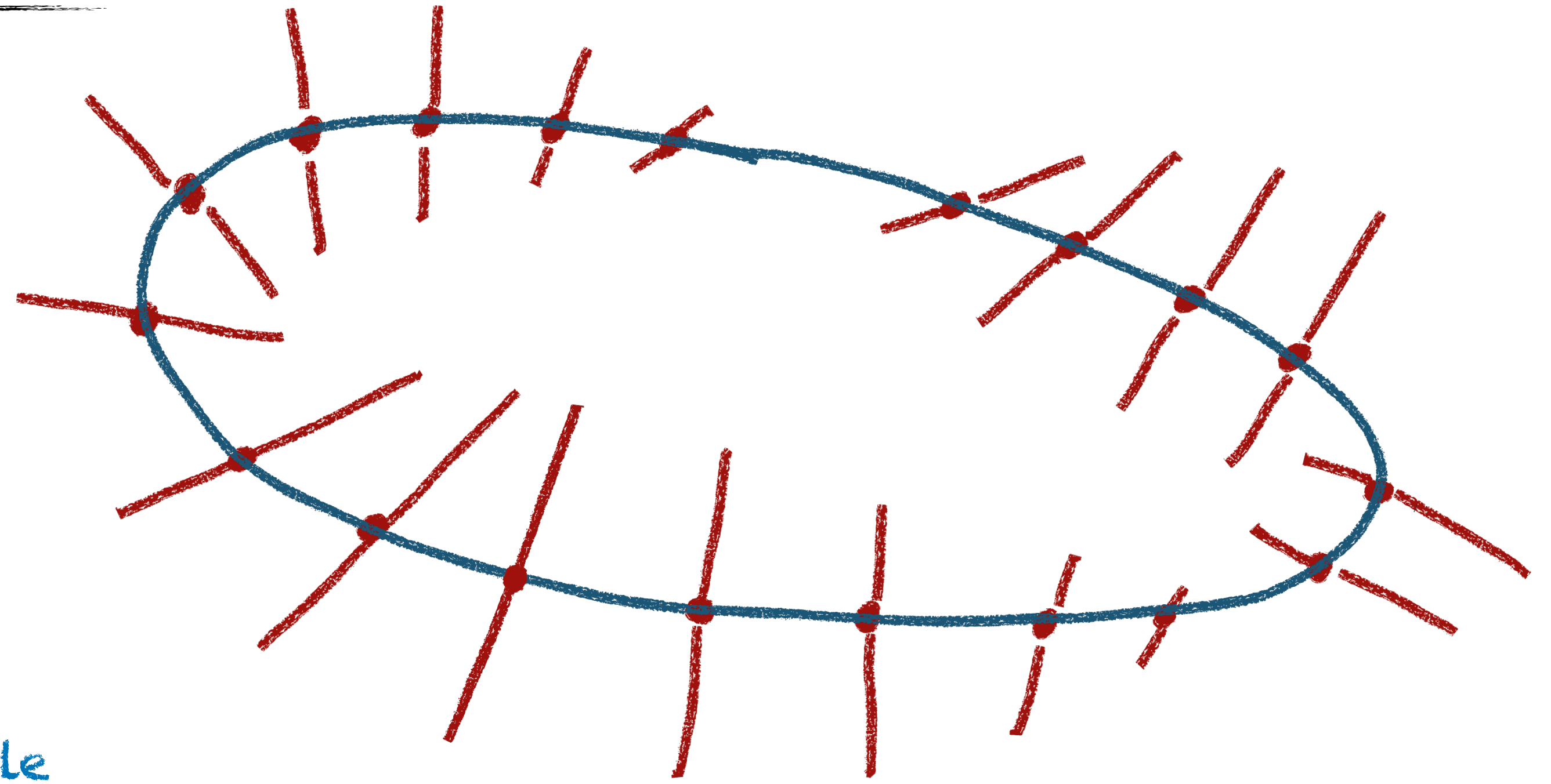


Elasticity of a Möbius strip

Definition of the shape of the strip

- ▶ discrete version of the strip
- ▶ lattice of « directions »
- ▶ but elasticity requires vectors !
- ▶ link variable $\eta_{i,i+1}$: coherence of site orientations
- ▶ orientability:

$$\mathcal{O} = \prod_{i=1}^N \eta_{i,i+1} = \begin{cases} +1 & \text{orientable} \\ -1 & \text{non-orientable} \end{cases}$$



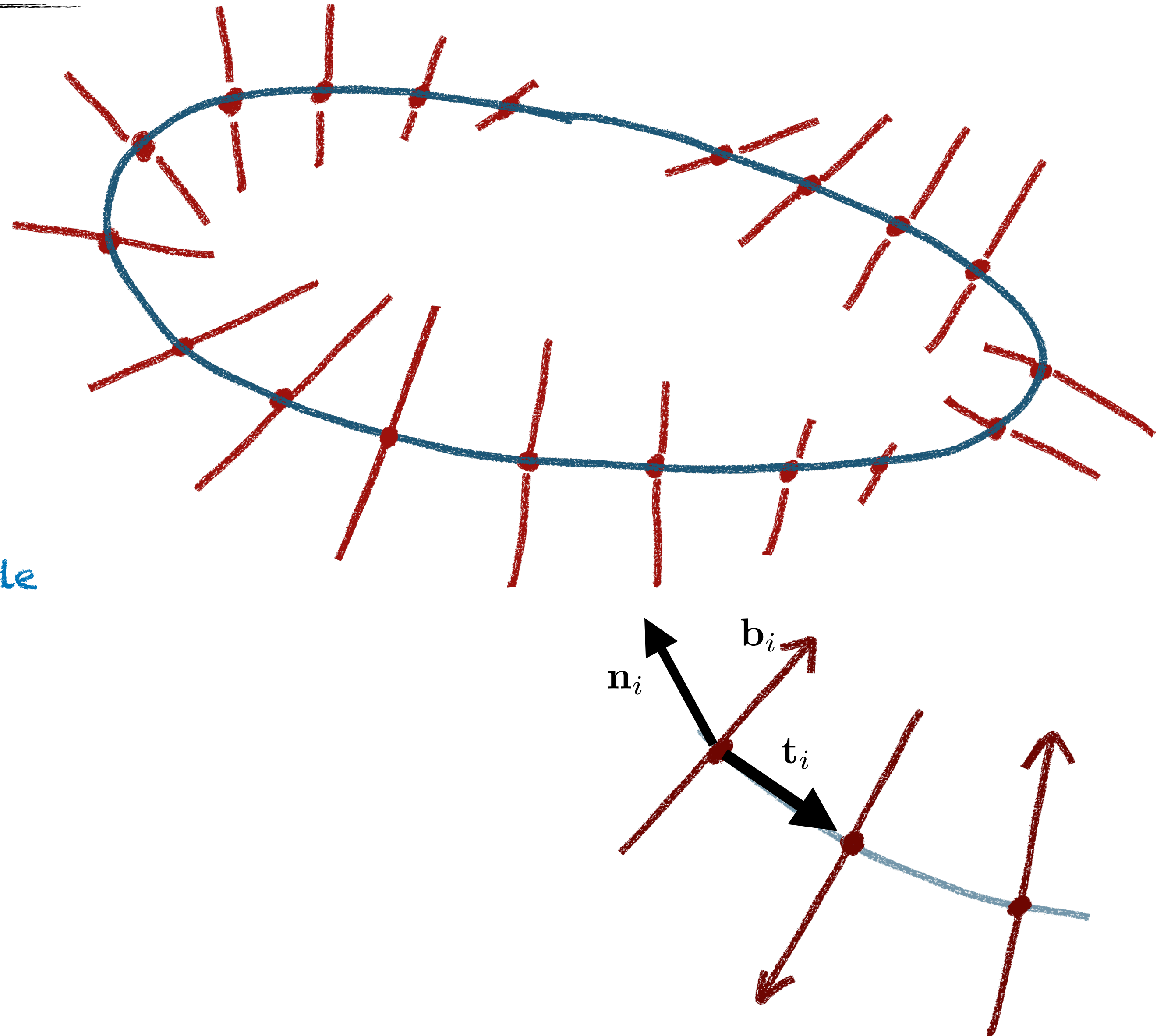
Elasticity of a Möbius strip

Definition of the shape of the strip

- ▶ discrete version of the strip
- ▶ lattice of « directions »
- ▶ but elasticity requires vectors !
- ▶ link variable $\eta_{i,i+1}$: coherence of site orientations
- ▶ orientability:

$$\mathcal{O} = \prod_{i=1}^N \eta_{i,i+1} = \begin{cases} +1 & \text{orientable} \\ -1 & \text{non-orientable} \end{cases}$$

- ▶ local basis $\mathbf{t}_i, (\epsilon_i)\mathbf{b}_i, (\epsilon_i)\mathbf{n}_i$



Elasticity of a Möbius strip

Shear deformations

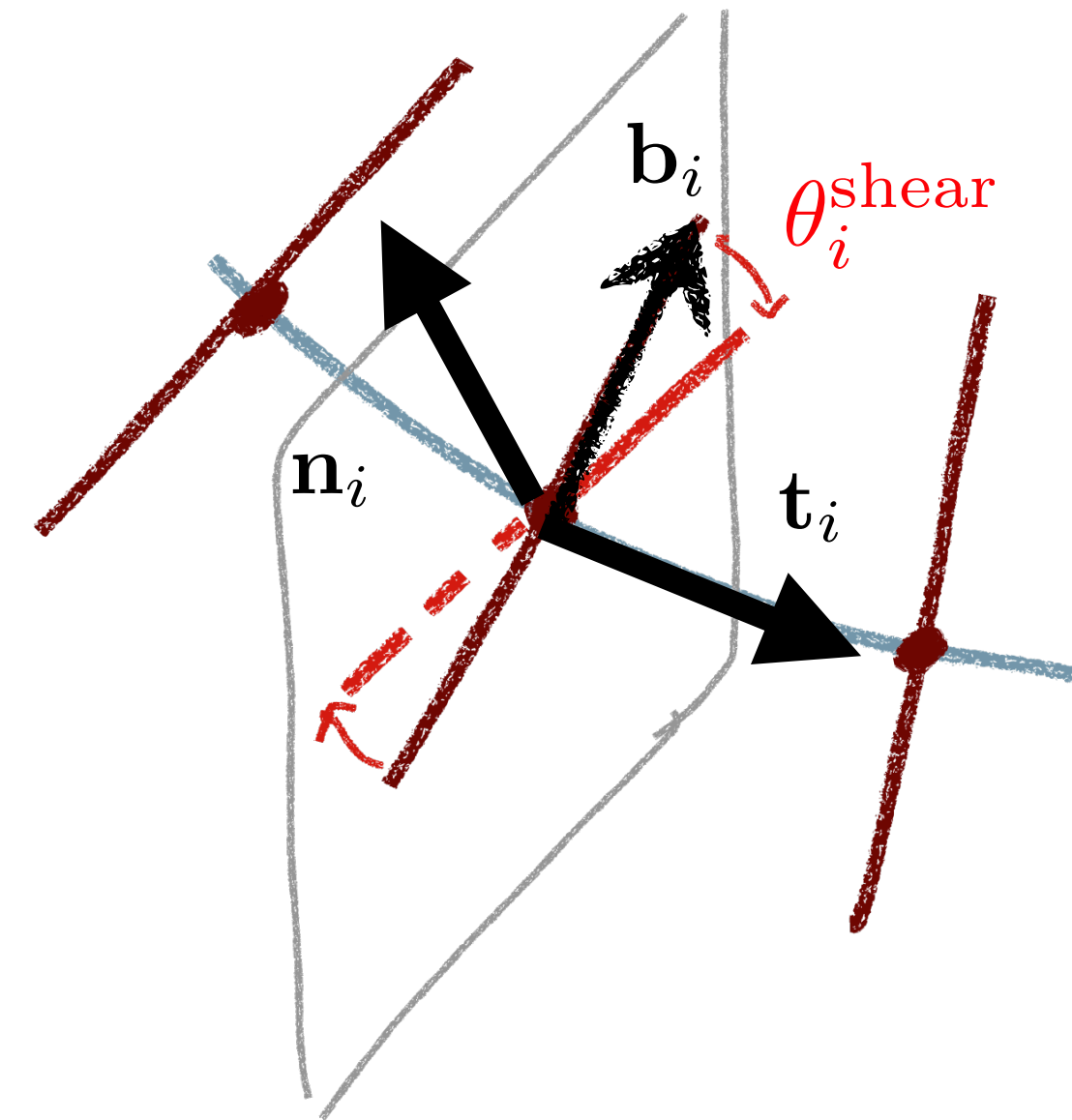
► deformation field: $\epsilon_i \mathbf{u}_i = \epsilon_i \theta_i^{\text{shear}} \mathbf{t}_i^0$.

► elastic energy: $\frac{E}{N} = \sum_i \frac{K_i^s}{2} [\theta_{i+1}^{\text{shear}} - \eta_{i,i+1} \theta_i^{\text{shear}}]^2$

Z_2 gauge theory

► continuous elasticity: $E = \int \frac{K_s}{2} (\partial_s \theta^{\text{shear}})^2$.

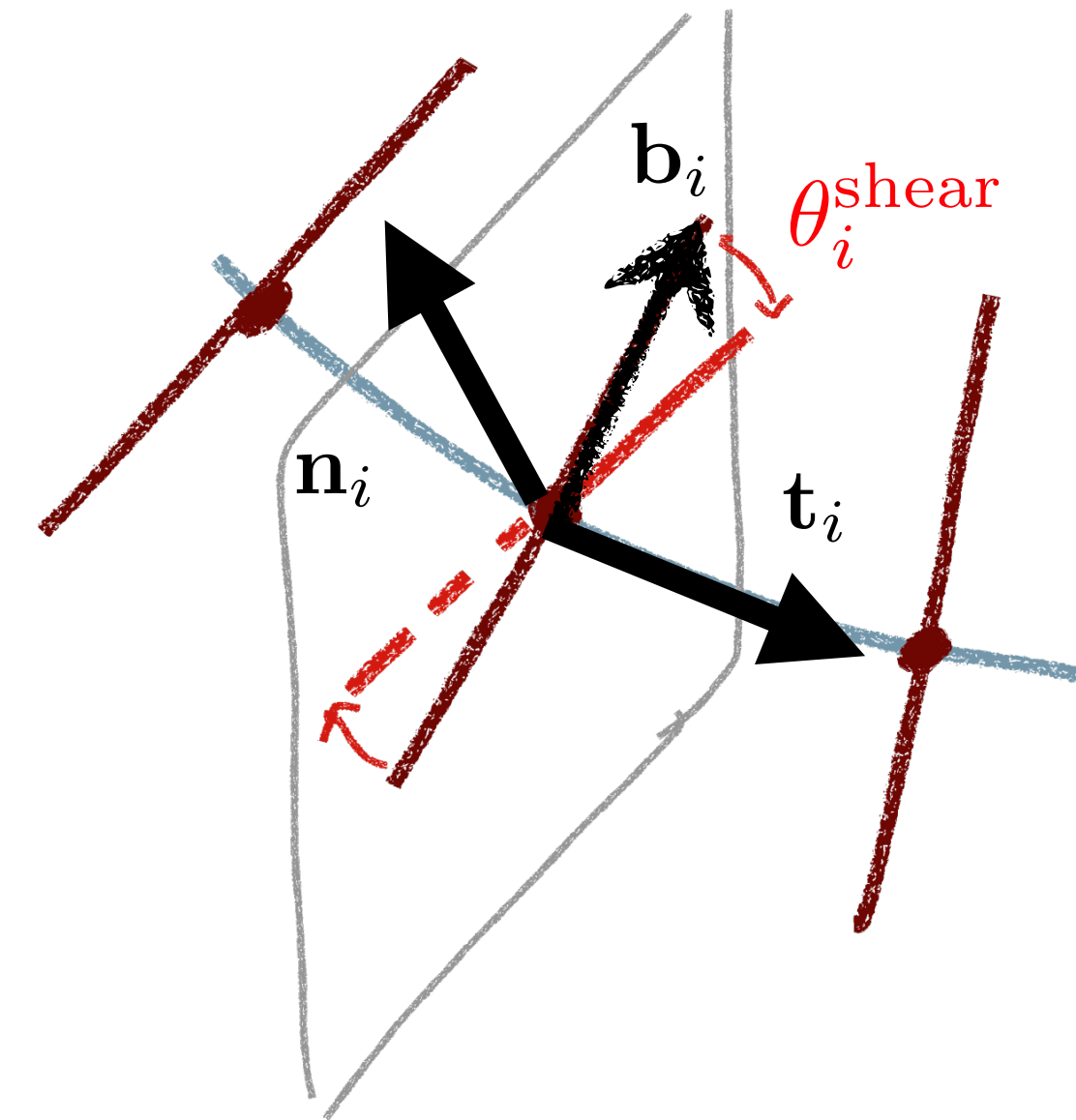
additional condition: $\theta^{\text{shear}}(s_0) = 0$ if $\mathcal{O} = -1$
with s_0 free (degree of freedom)



Elasticity of a Möbius strip

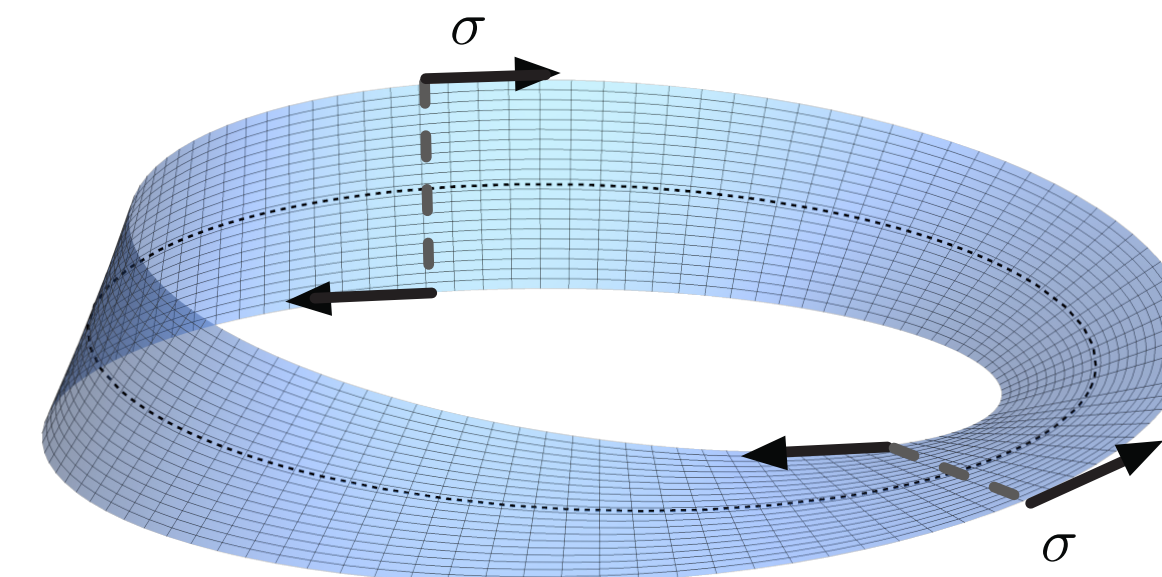
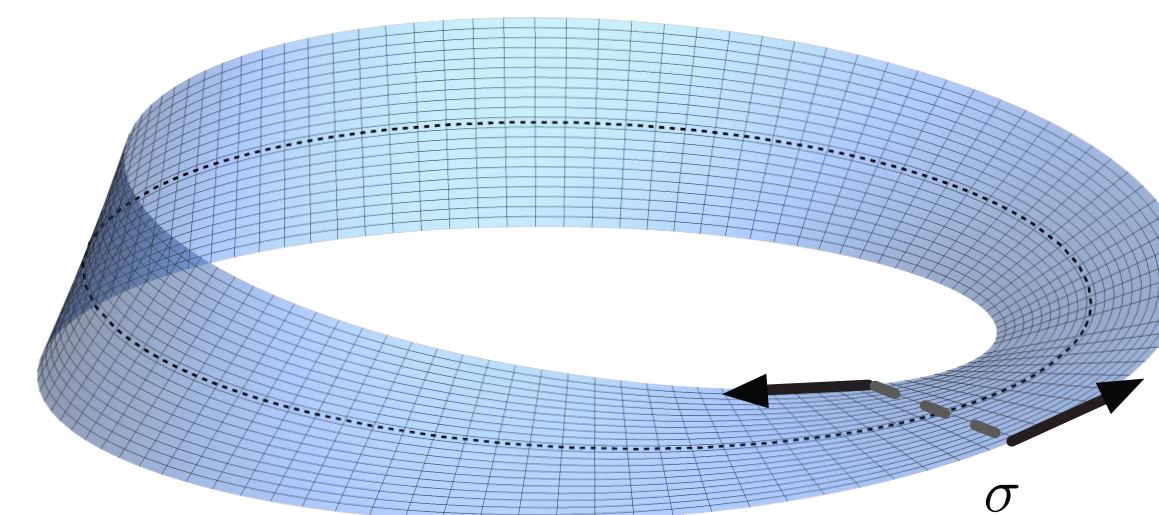
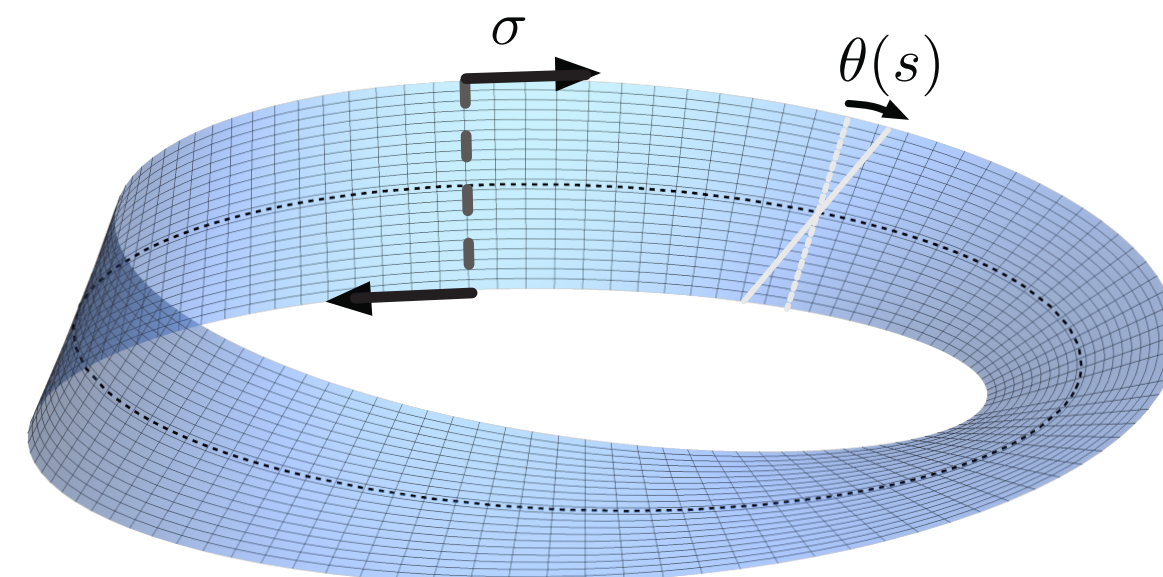
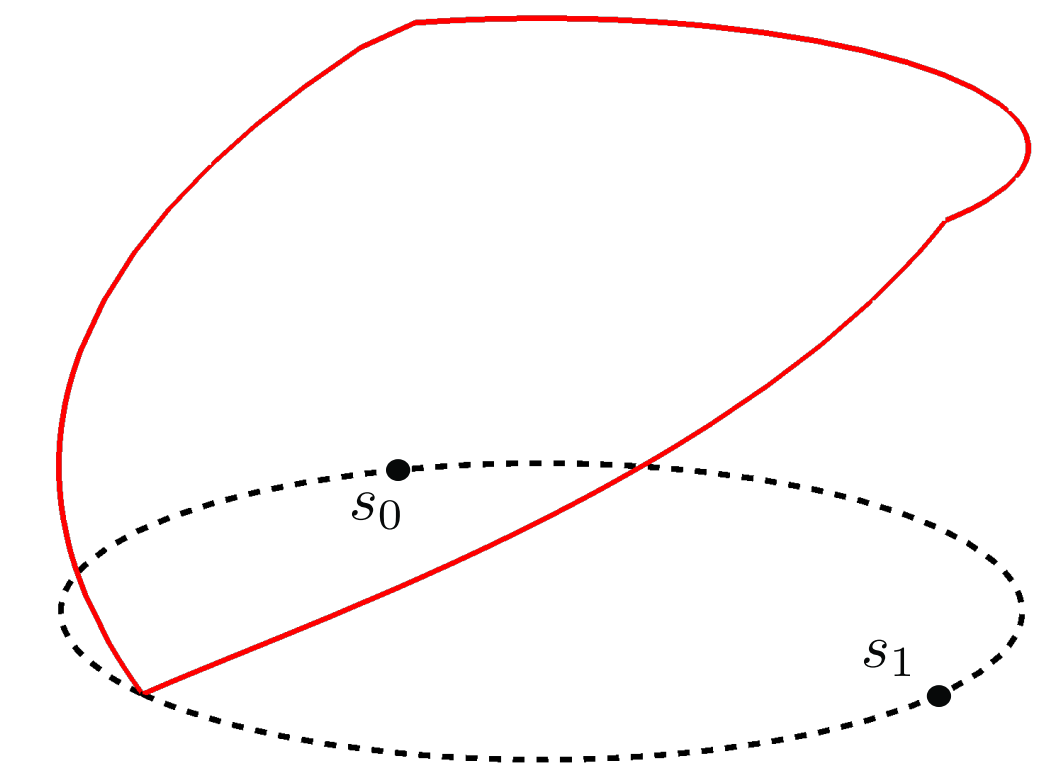
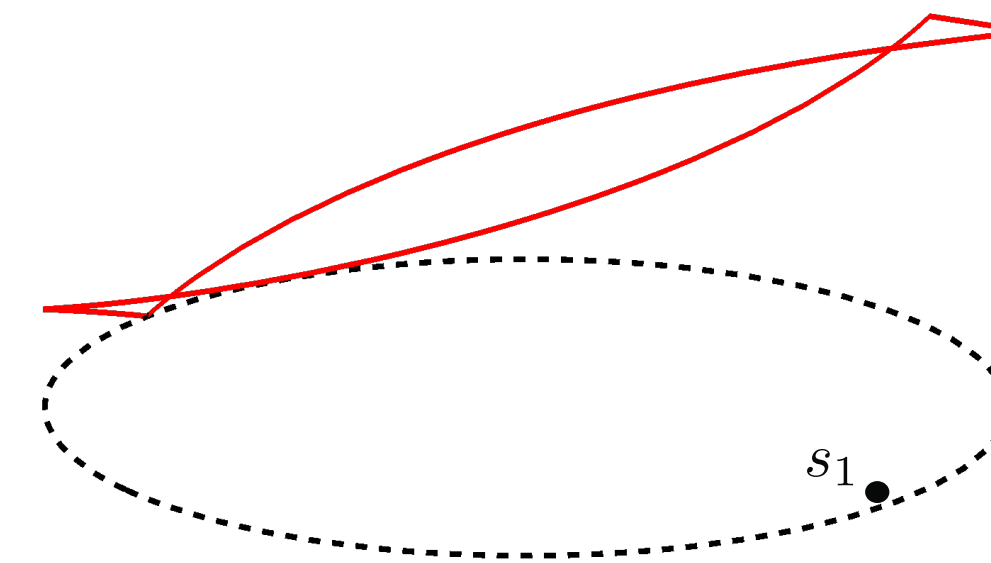
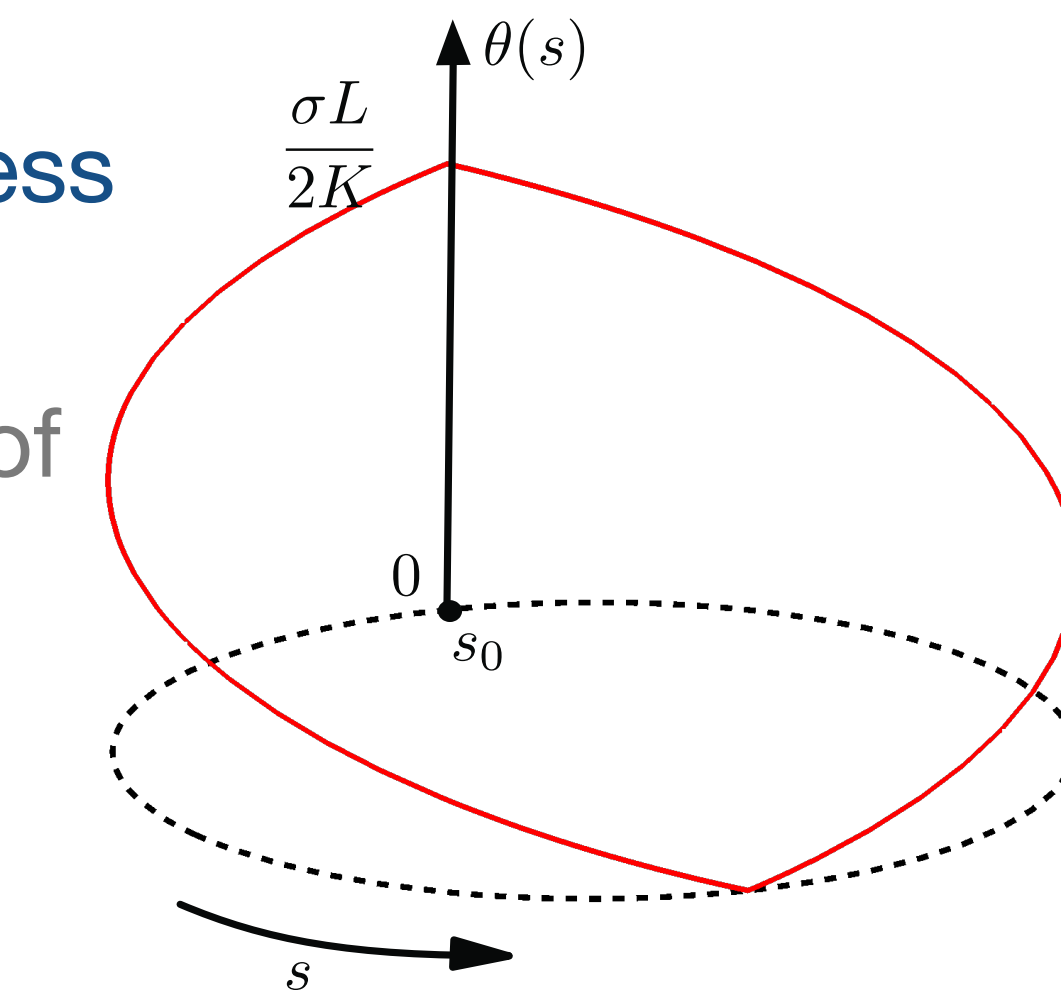
Shear deformations

- ▶ continuous elasticity: $E = \int \frac{K_s}{2} (\partial_s \theta^{\text{shear}})^2$.
- ▶ additional condition: $\theta^{\text{shear}}(s_0) = 0$ if $\mathcal{O} = -1$ with s_0 free (degree of freedom)



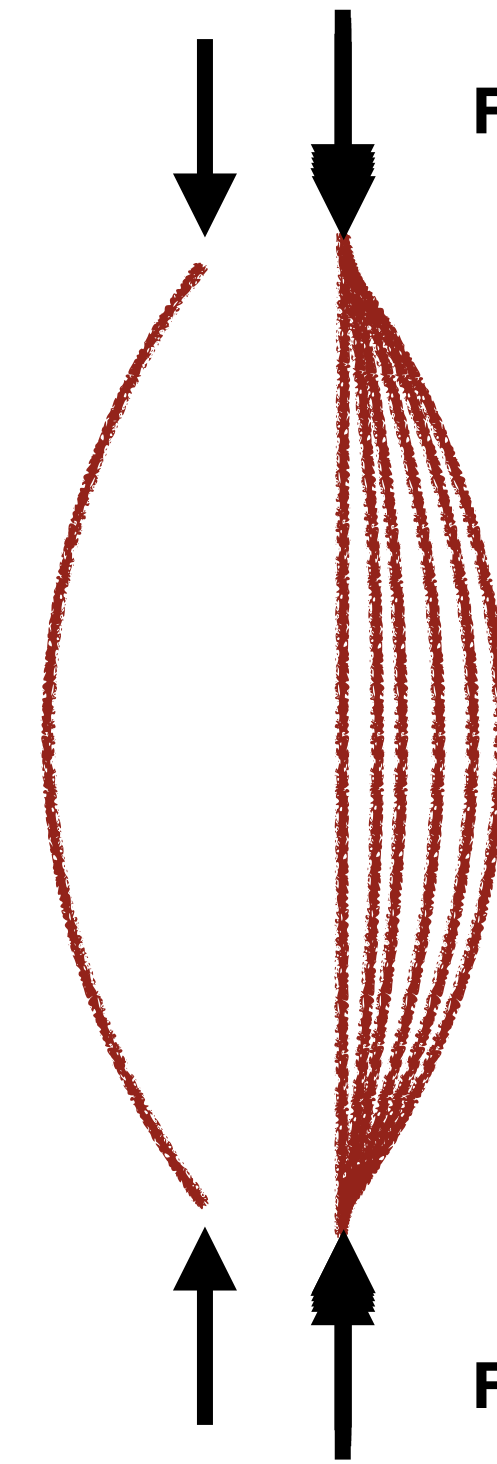
Response to local shear stress

- ▶ non linear response
- ▶ depends on « history » of constraints



Buckling of a Möbius strip

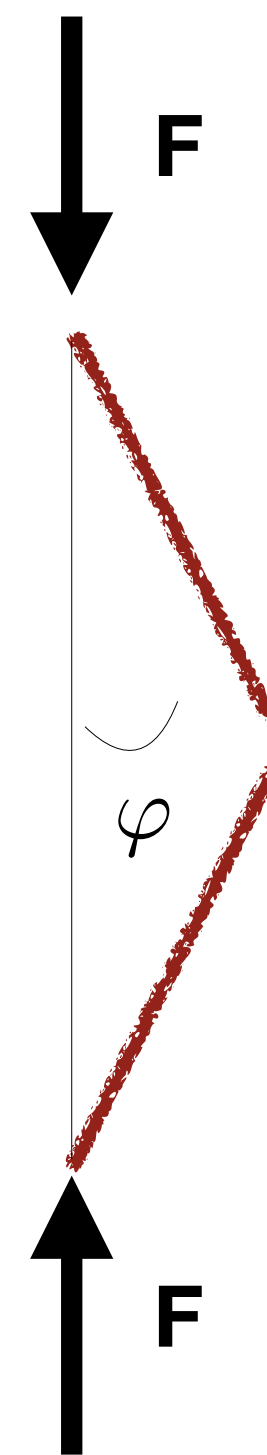
- What is buckling ?
 - ▶ nonlinear relation $F \leftrightarrow$ deformation
 - ▶ can be applied homogeneously on the Möbius strip



Buckling of a Möbius strip

- What is buckling ?
 - ▶ nonlinear relation $F \leftrightarrow$ deformation
 - ▶ can be applied homogeneously on the Möbius strip
- Buckling of the Möbius strip
 - ▶ continuous elasticity

$$\mathcal{E}_B = \frac{1}{2} \int \left[K'_B (\partial_s \varphi)^2 + K_B \varphi^2 + \Sigma l \cos \varphi \right] ds \quad \text{with} \quad \varphi(s_0) = 0$$



Work of F

$$E = \frac{1}{2} K_B \varphi^2 - FL(1 - \cos \varphi)$$

buckling if $FL > K_B$

Buckling of a Möbius strip

- What is buckling ?

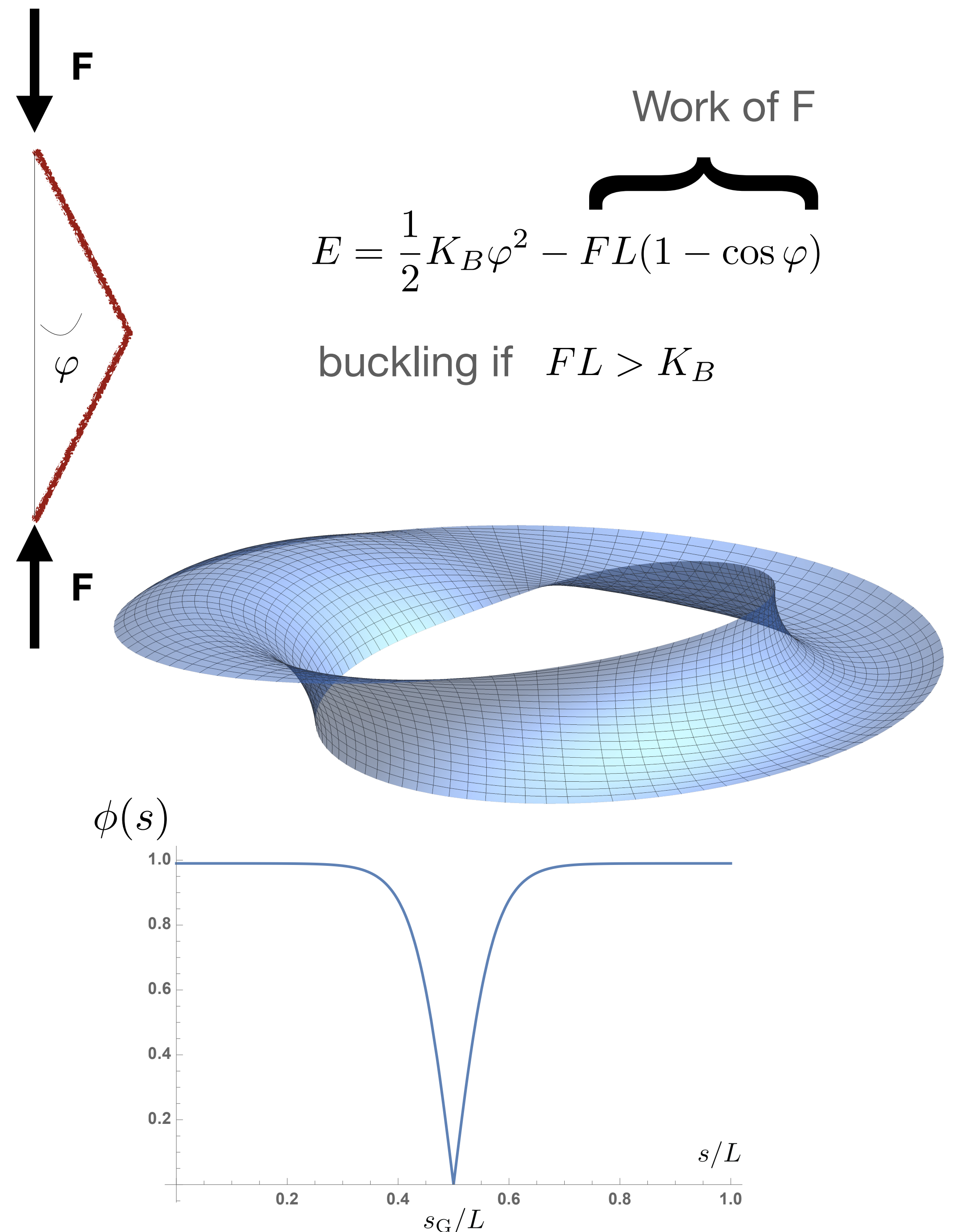
- ▶ nonlinear relation $F \leftrightarrow$ deformation
- ▶ can be applied homogeneously on the Möbius strip

- Buckling of the Möbius strip

- ▶ continuous elasticity

$$\mathcal{E}_B = \frac{1}{2} \int \left[K'_B (\partial_s \varphi)^2 + K_B \varphi^2 + \Sigma l \cos \varphi \right] ds \quad \text{with } \varphi(s_0) = 0$$

- ▶ **1 point** along the strip remains **undistorted** !
- ▶ topologically protected **solitary wave** of distortion along the strip
- ▶ **topological \mathbb{Z}_2 charge** (Witten-Olive 1978)



Buckling of a Möbius strip



Marcelo Guzman

- Buckling of the Möbius strip

- ▶ continuous elasticity

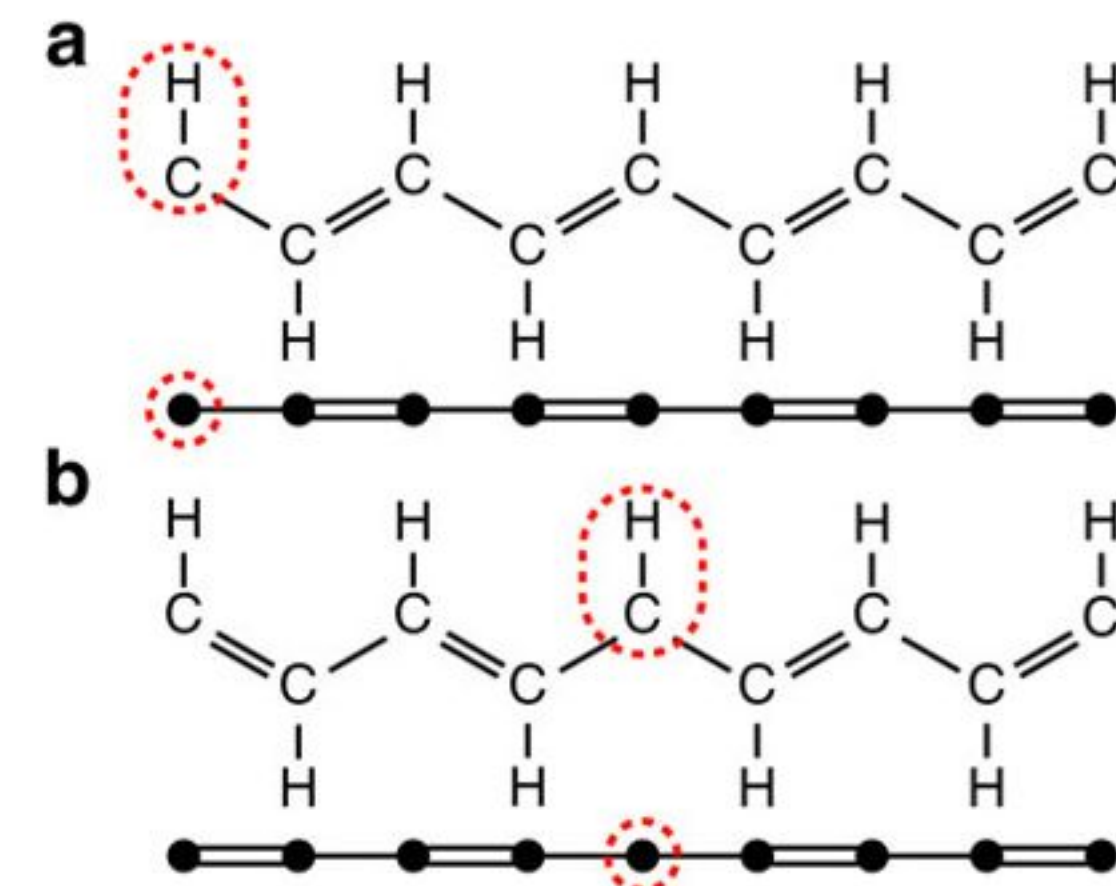
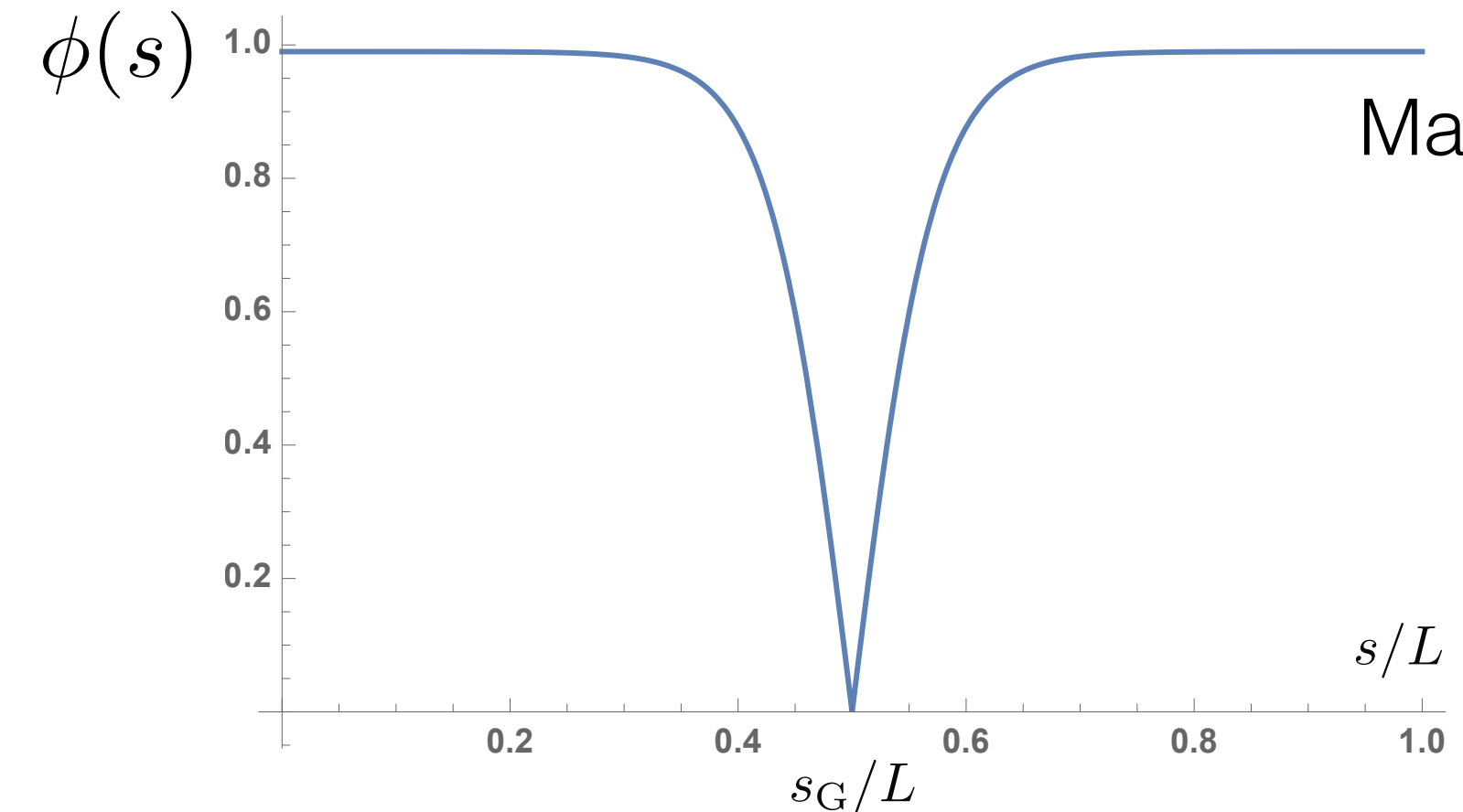
$$\mathcal{E}_B = \frac{1}{2} \int \left[K'_B (\partial_s \varphi)^2 + K_B \varphi^2 + \Sigma l \cos \varphi \right] ds \quad \text{with} \quad \varphi(s_0) = 0$$

- ▶ **1 point** along the strip remains **undistorted** !
- ▶ topologically protected **solitary wave** of distortion along the strip
- ▶ **topological \mathbb{Z}_2 charge** (Witten-Olive 1978)

- Su–Schrieffer–Heeger model of polyacetylene

- ▶ 2 configurations topologically distincts
- ▶ boundary state between them = soliton !
- ▶ same **topological \mathbb{Z}_2 charge**

Möbius soliton : topological edge state without an edge !



Summary

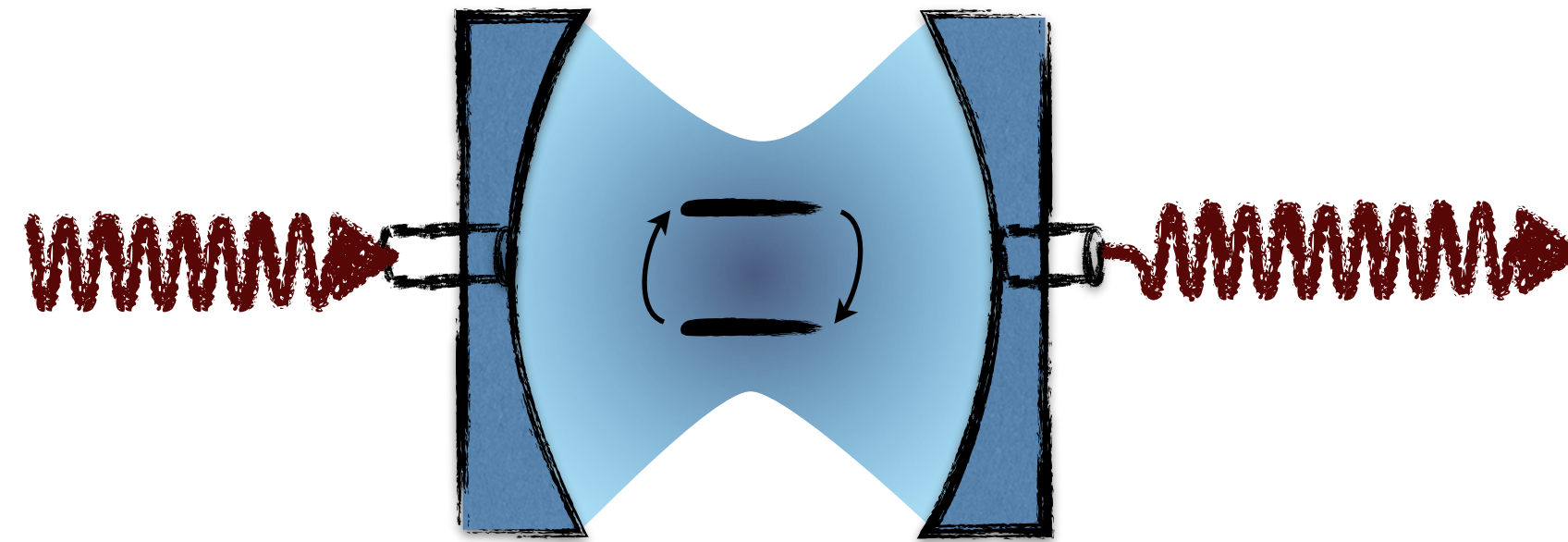
1. Electronic Properties of Quantum Matter

Topological Insulators

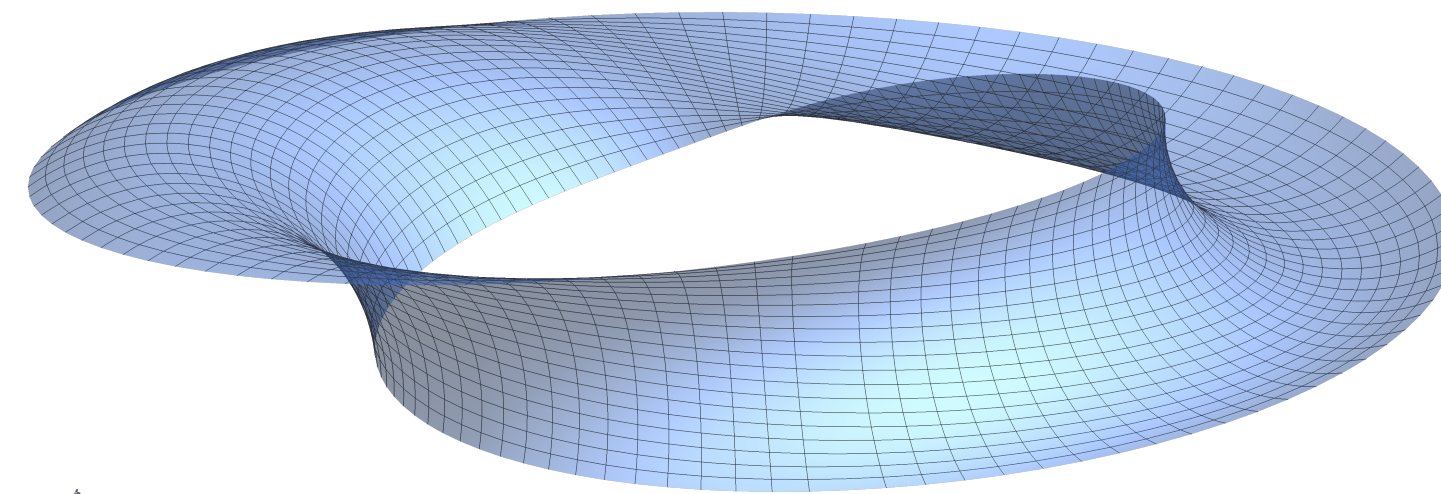


2. Quantum Technologies

Topological pump



3. Mechanics / Metamaterials



What's next ? (any suggestion ?)

Topology and Symmetry (I)

Kitaev (2010)
Schnyder, Ryu, Furusaki, Ludwig (2008)

	Cartan Nomenclature	TRS	PHS	SLS	$d = 1$	$d = 2$	$d = 3$
standard (Wigner-Dyson)	A (unitary)	0	0	0	-	\mathbb{Z}	-
	AI (orthogonal)	+1	0	0	-	-	-
	AII (symplectic)	-1	0	0	-	\mathbb{Z}_2	\mathbb{Z}_2
chiral (sublattice)	AIII (chiral unitary)	0	0	1	\mathbb{Z}	-	\mathbb{Z}
	BDI (chiral orthogonal)	+1	+1	1	\mathbb{Z}	-	-
	CII (chiral symplectic)	-1	-1	1	\mathbb{Z}	-	\mathbb{Z}_2
BdG	D	0	+1	0	\mathbb{Z}_2	\mathbb{Z}	-
	C	0	-1	0	-	\mathbb{Z}	-
Superconductors	DIII	-1	+1	1	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}
	CI	+1	-1	1	-	-	\mathbb{Z}

Time Reversal Symmetry :

- ▶ anti unitary ($\Theta = UK$)
- ▶ states at same E, $[H, \Theta] = 0$
- ▶ $\Theta^2 = \pm \mathbb{I}$

Particle Hole Symmetry :

- ▶ anti unitary ($P = VK$)
- ▶ states at opposite E, $\{H, P\} = 0$
- ▶ $P^2 = \pm \mathbb{I}$

Chiral / Sublattice Symmetry :

- ▶ unitary
- ▶ states at opposite E, $\{H, C\} = 0$
- ▶ $C = \Theta.P$

Symmetries and topology: classification

Kitaev (2010)
Schnyder, Ryu, Furusaki, Ludwig (2008)

Here: only general (ubiquitous) symmetries
Other symmetries: see T. Neupert Lectures

(superconductors)

Time Reversal Symmetry :

- ▶ anti unitary ($\Theta = UK$)
- ▶ states at same E, $[H, \Theta] = 0$
- ▶ $\Theta^2 = \pm \mathbb{I}$

Particle Hole Symmetry :

- ▶ anti unitary ($P = VK$)
- ▶ states at opposite E, $\{H, P\} = 0$
- ▶ $P^2 = \pm \mathbb{I}$

Chiral / Sublattice Symmetry :

- ▶ $C = \Theta.P$
- ▶ unitary
- ▶ states at opposite E, $\{H, C\} = 0$

Symmetries and topology: classification

Kitaev (2010)
 Schnyder, Ryu, Furusaki, Ludwig (2008)

	Cartan Nomenclature	TRS	PHS	SLS	$d = 1$	$d = 2$	$d = 3$
standard (Wigner-Dyson)	A (unitary)	0	0	0	-	\mathbb{Z}	-
	AI (orthogonal)	+1	0	0	-	-	-
	AII (symplectic)	-1	0	0	-	\mathbb{Z}_2	\mathbb{Z}_2
chiral (sublattice)	AIII (chiral unitary)	0	0	1	\mathbb{Z}	-	\mathbb{Z}
	BDI (chiral orthogonal)	+1	+1	1	\mathbb{Z}	-	-
	CII (chiral symplectic)	-1	-1	1	\mathbb{Z}	-	\mathbb{Z}_2
BdG	D	0	+1	0	\mathbb{Z}_2	\mathbb{Z}	-
	C	0	-1	0	-	\mathbb{Z}	-
Superconductors	DIII	-1	+1	1	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}
	CI	+1	-1	1	-	-	\mathbb{Z}

Symmetries and topology : classification

		Cartan Nomenclature	TRS	PHS	SLS	dimension		
			$d = 1$	$d = 2$	$d = 3$			
standard (Wigner-Dyson)	A (unitary)	0	0	0	-	\mathbb{Z}	-	
	AI (orthogonal)	+1	0	0	-	-	-	
	AII (symplectic)	-1	0	0	-	\mathbb{Z}_2	\mathbb{Z}_2	
chiral (sublattice)	AIII (chiral unitary)	0	0	1	\mathbb{Z}	-	\mathbb{Z}	
	BDI (chiral orthogonal)	+1	+1	1	\mathbb{Z}	-	-	
	CII (chiral symplectic)	-1	-1	1	\mathbb{Z}	-	\mathbb{Z}_2	
BdG	D	0	+1	0	\mathbb{Z}_2	\mathbb{Z}	-	
	C	0	-1	0	-	\mathbb{Z}	-	
Superconductors	DIII	-1	+1	1	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	
	CI	+1	-1	1	-	-	\mathbb{Z}	

Symmetry Operator : U

- 0** : no symmetry
- +1** : symmetry with $U^2 = \mathbb{I}$
- 1** : symmetry with $U^2 = -\mathbb{I}$

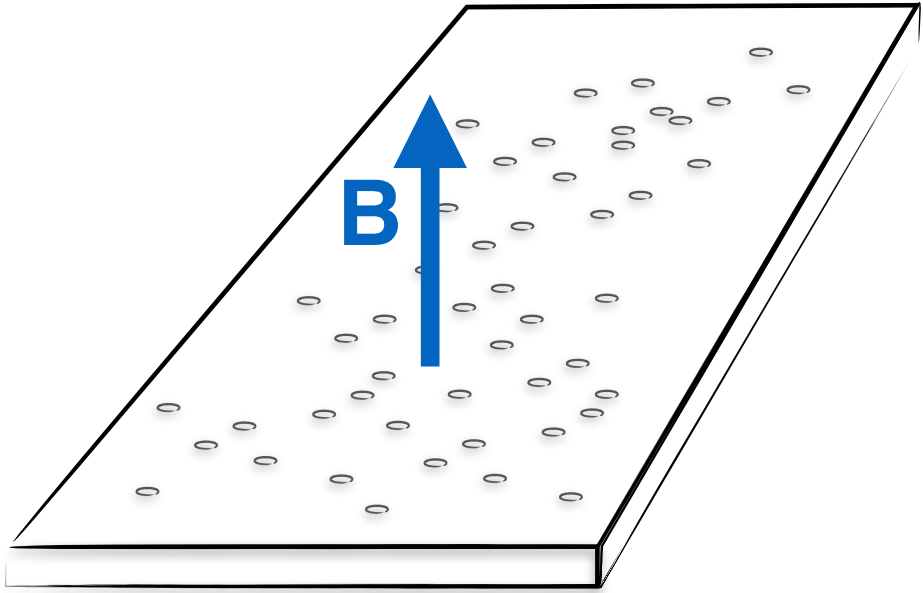
10 classes :

- TRS : **x3** (0,+1,-1)
- PHS : **x3** (0,+1,-1)
- SLS = TRS . PHS (**+1**)

Symmetries and topology : classification

Chern insulators : ex. Quantum Hall Effect

- ▶ $d=2$
- ▶ breaks all symmetries (TRS)
- ▶ Topological index : Chern number



2DEG (Heterojunction GaAs/AlGaAs)

Thouless, Kohmoto, Nightingale and den Nijs (1982)
 Niu, Thouless, and Wu (1985)
 Haldane (1985)

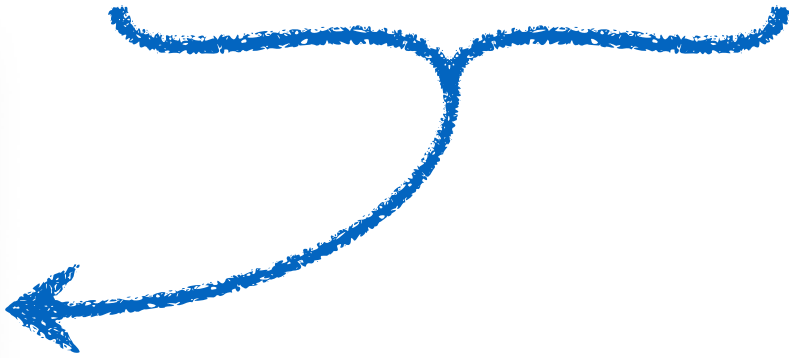
	Cartan Nomenclature	TRS	PHS	SLS	$d = 1$	$d = 2$	$d = 3$
standard (Wigner-Dyson)	A (unitary)	0	0	0	-	\mathbb{Z}	-
	AI (orthogonal)	+1	0	0	-	-	-
	AII (symplectic)	-1	0	0	-	\mathbb{Z}_2	\mathbb{Z}_2
chiral (sublattice)	AIII (chiral unitary)	0	0	1	\mathbb{Z}	-	\mathbb{Z}
	BDI (chiral orthogonal)	+1	+1	1	\mathbb{Z}	-	-
	CII (chiral symplectic)	-1	-1	1	\mathbb{Z}	-	\mathbb{Z}_2
BdG	D	0	+1	0	\mathbb{Z}_2	\mathbb{Z}	-
	C	0	-1	0	-	\mathbb{Z}	-
Superconductors	DIII	-1	+1	1	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}
	CI	+1	-1	1	-	-	\mathbb{Z}

Symmetry Operator : U

0 : no symmetry

+1 : symmetry with $U^2 = \mathbb{I}$

-1 : symmetry with $U^2 = -\mathbb{I}$



Symmetries and topology : classification

Topological insulators :

- ▶ d=2 and d=3
- ▶ TRS with spin 1/2 : spin-orbit
- ▶ Topological index : Kane-Mele

Kane and Mele (2005)
 Bernevig, Hughes, and Zhang (2006)
 Fu, Kane et Mele (2007)
 Moore and Balents (2007)
 Roy (2009)
 Fu and Kane (2007)

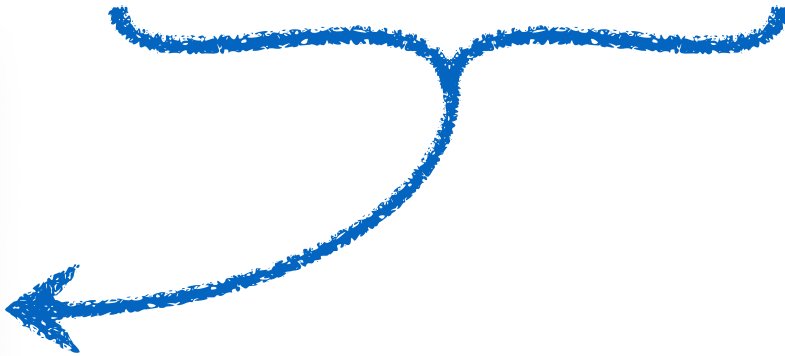
	Cartan Nomenclature	TRS	PHS	SLS	d = 1	d = 2	d = 3
standard (Wigner-Dyson)	A (unitary)	0	0	0	-	\mathbb{Z}	-
	AI (orthogonal)	+1	0	0	-	-	-
	AII (symplectic)	-1	0	0	-	\mathbb{Z}_2	\mathbb{Z}_2
chiral (sublattice)	AIII (chiral unitary)	0	0	1	\mathbb{Z}	-	\mathbb{Z}
	BDI (chiral orthogonal)	+1	+1	1	\mathbb{Z}	-	-
	CII (chiral symplectic)	-1	-1	1	\mathbb{Z}	-	\mathbb{Z}_2
BdG	D	0	+1	0	\mathbb{Z}_2	\mathbb{Z}	-
	C	0	-1	0	-	\mathbb{Z}	-
Superconductors	DIII	-1	+1	1	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}
	CI	+1	-1	1	-	-	\mathbb{Z}

Symmetry Operator : U

0 : no symmetry

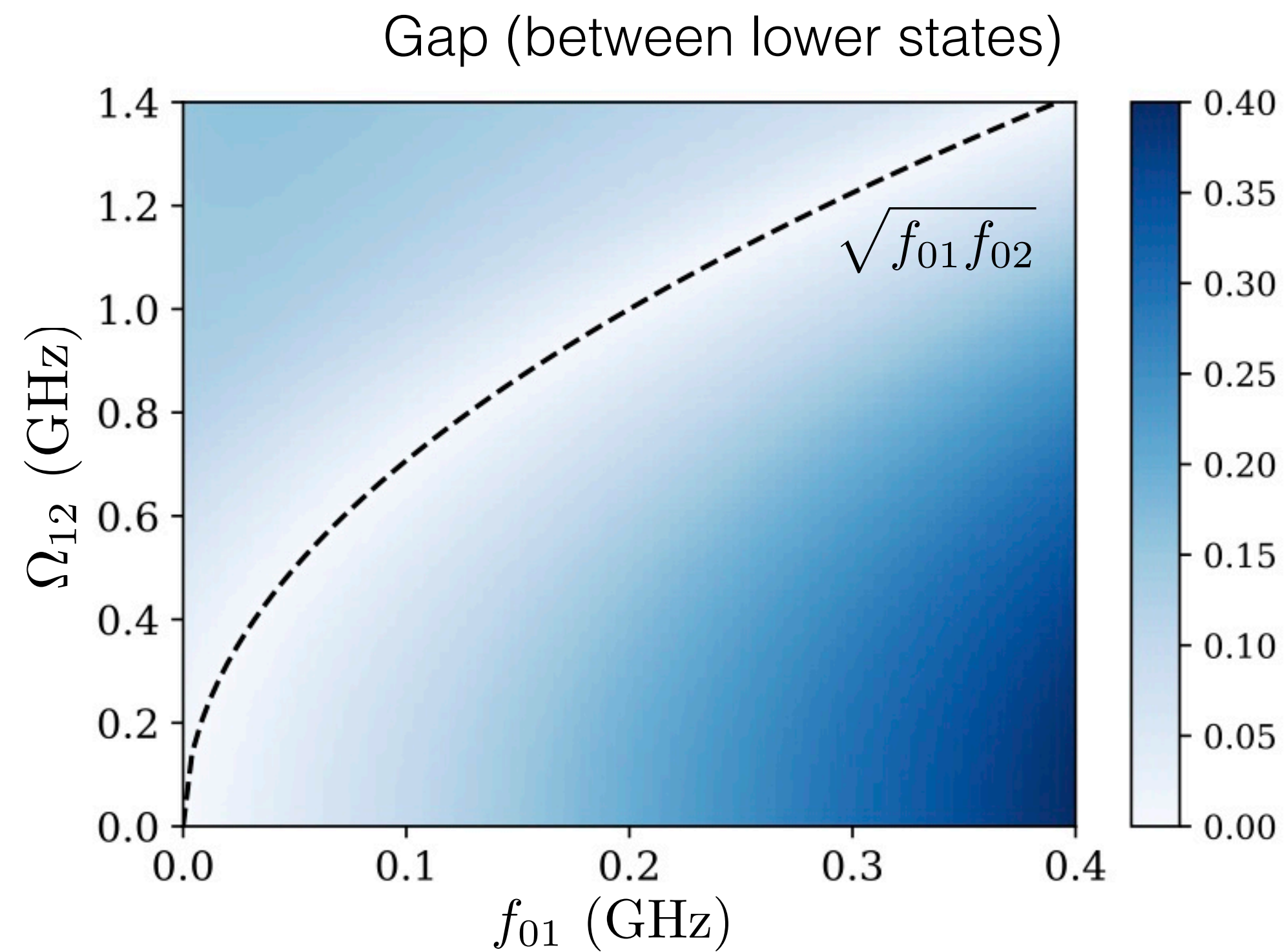
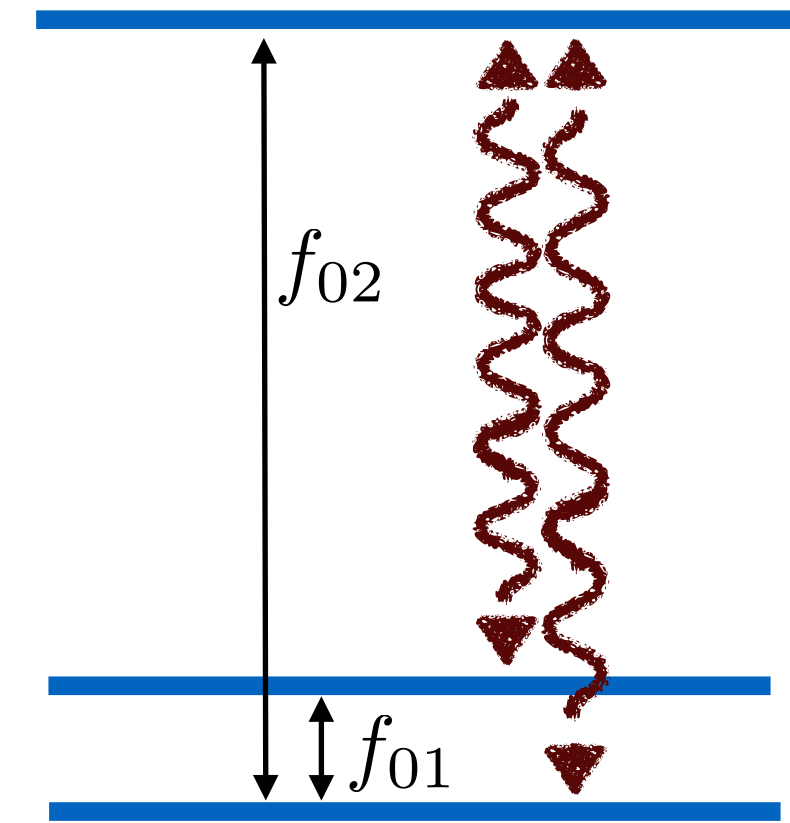
+1 : symmetry with $U^2 = \mathbb{I}$

-1 : symmetry with $U^2 = -\mathbb{I}$



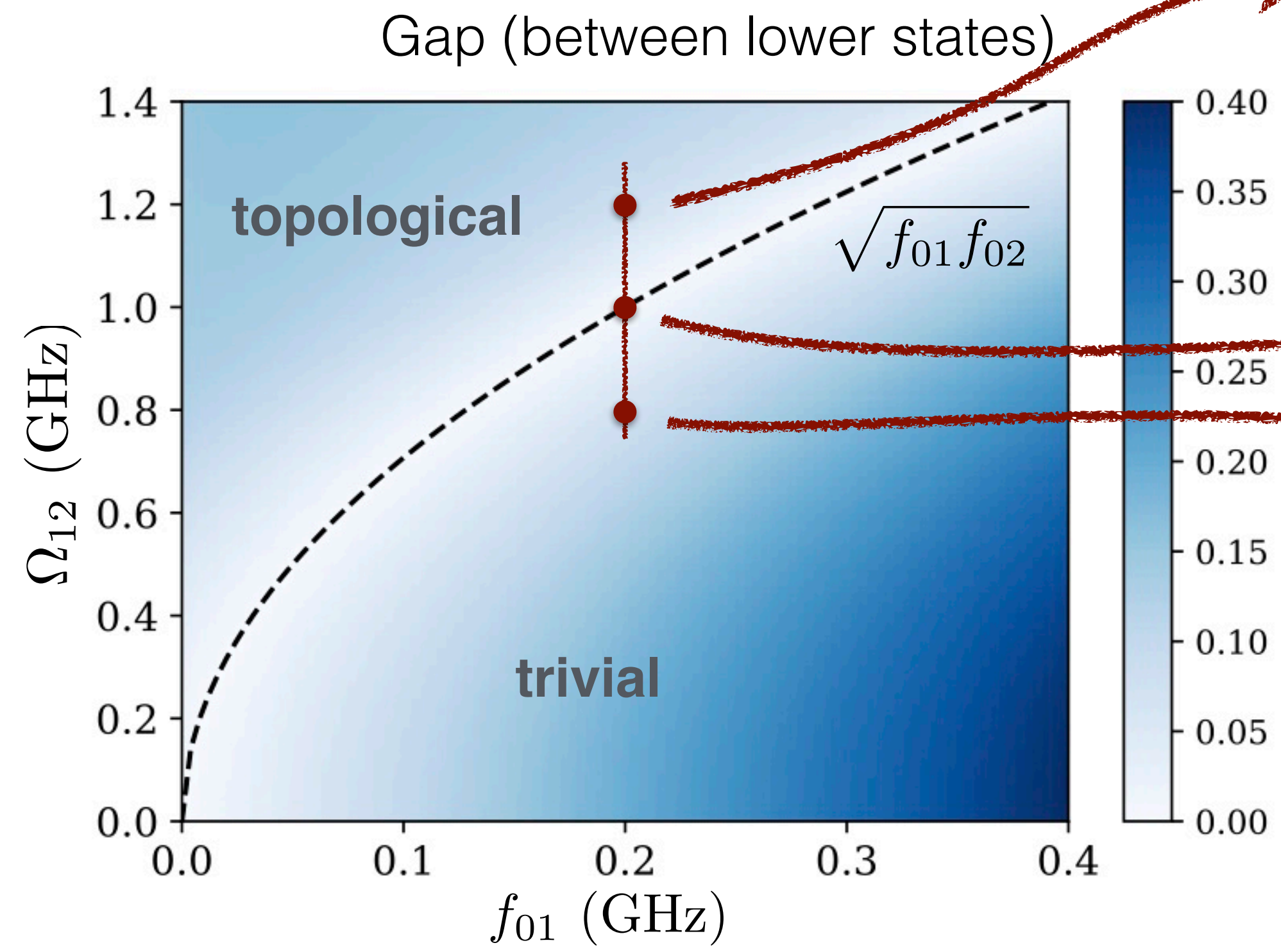
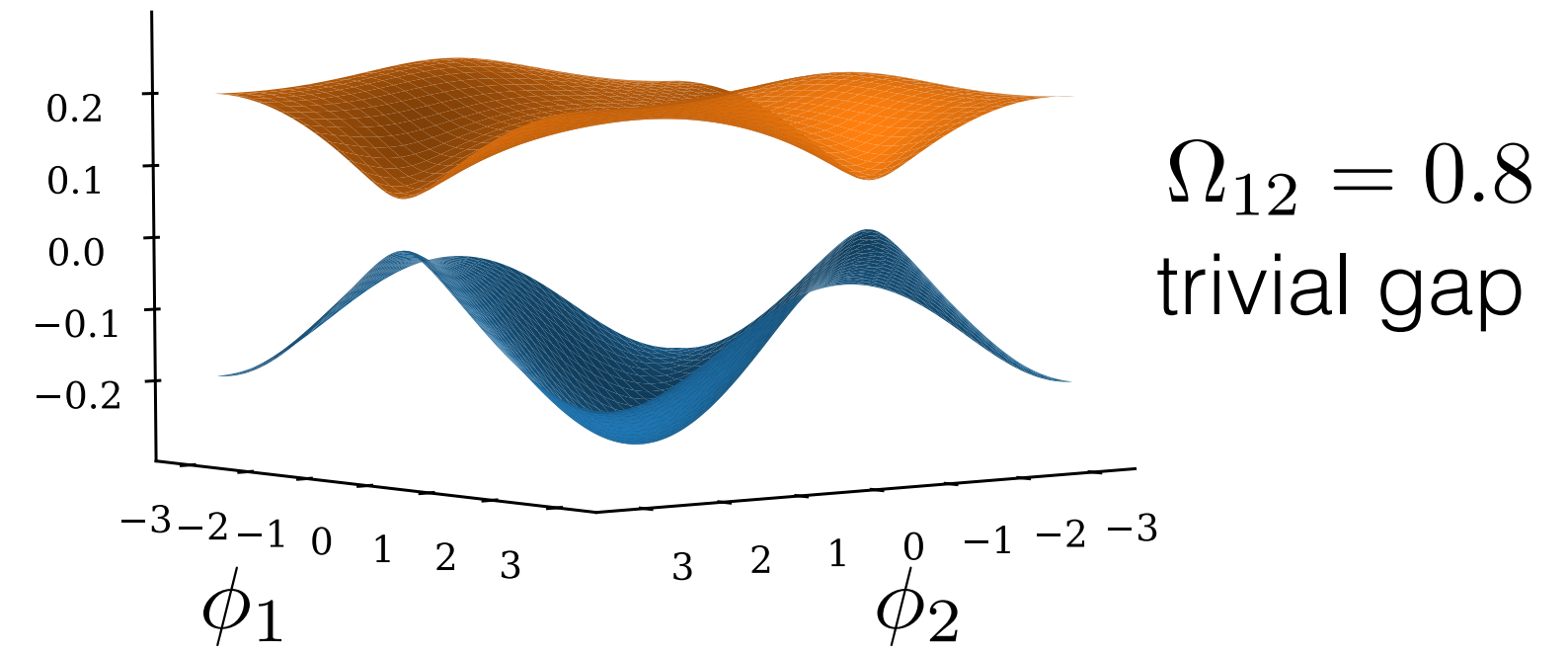
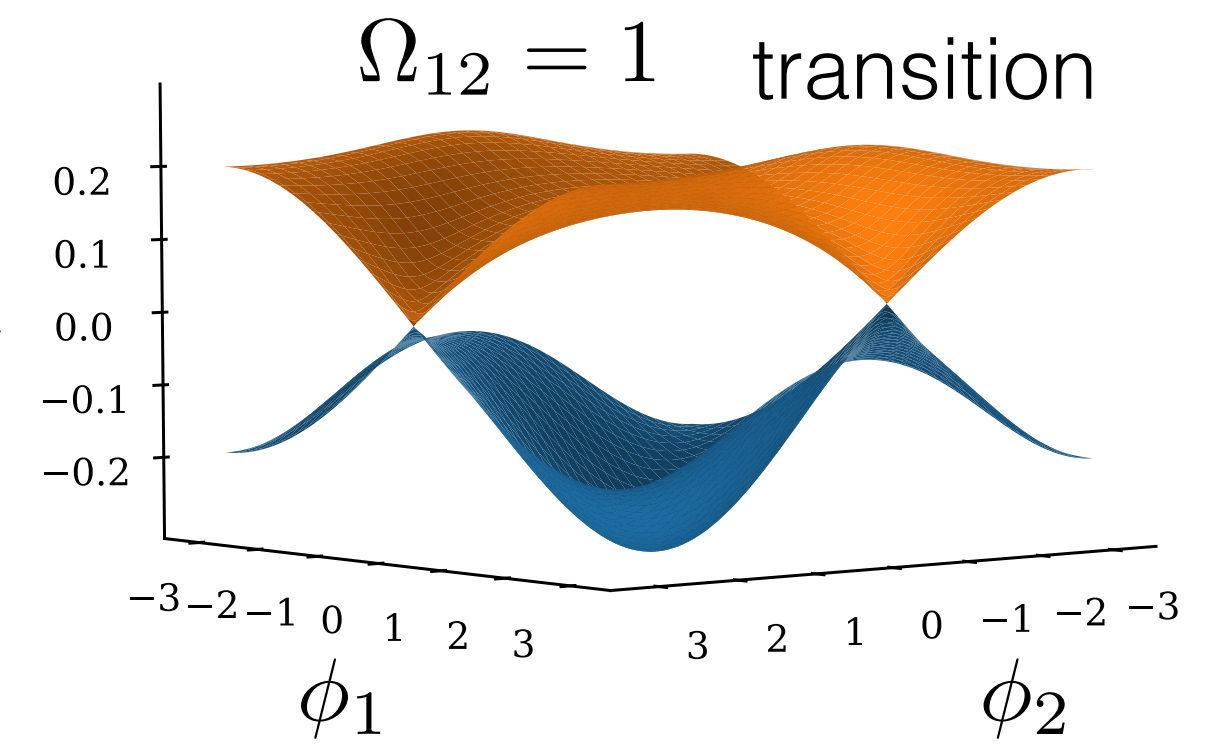
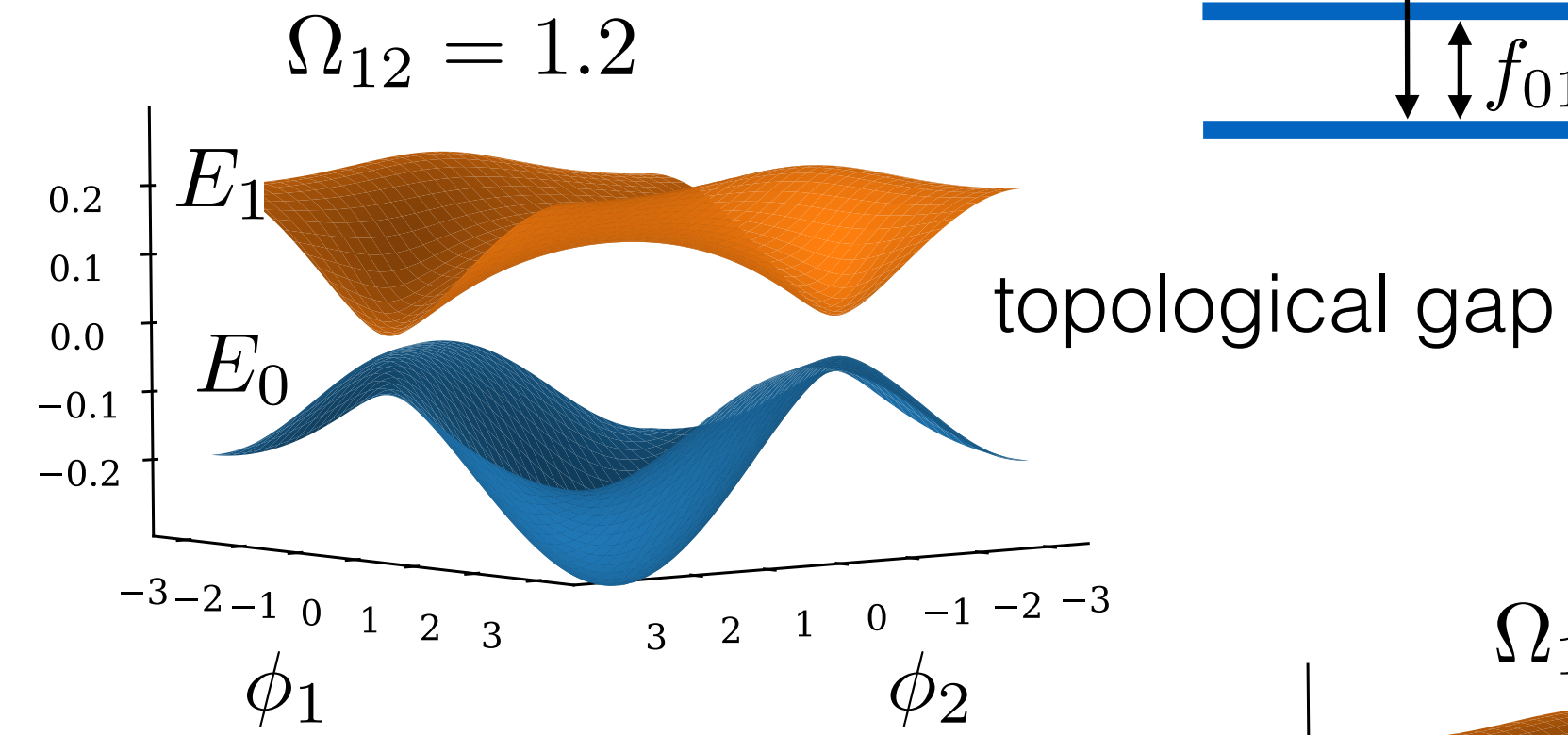
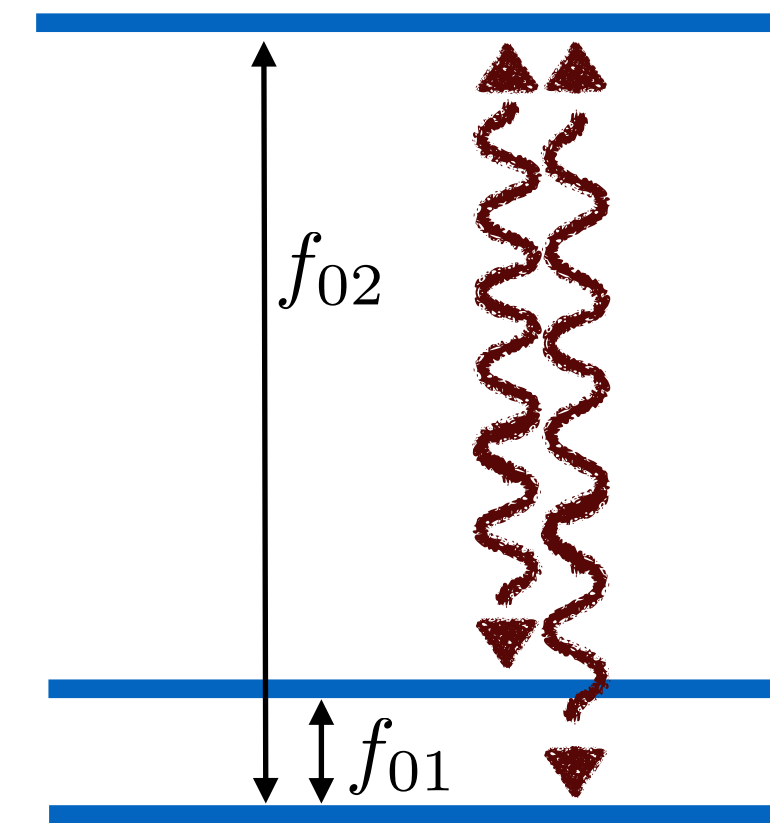
Topological Pump : effective 2 levels (qubit)

ω_{01}	$2\pi \times 200$ MHz
ω_{02}	$2\pi \times 5$ GHz
Ω_{01}	$2\pi \times 100$ MHz
Ω_{02}	$2\pi \times 1$ GHz
Ω_{12}	$2\pi \times 1.2$ GHz
$\dot{\phi}_a = 2\dot{\phi}_b$	$2\pi \times 20$ MHz

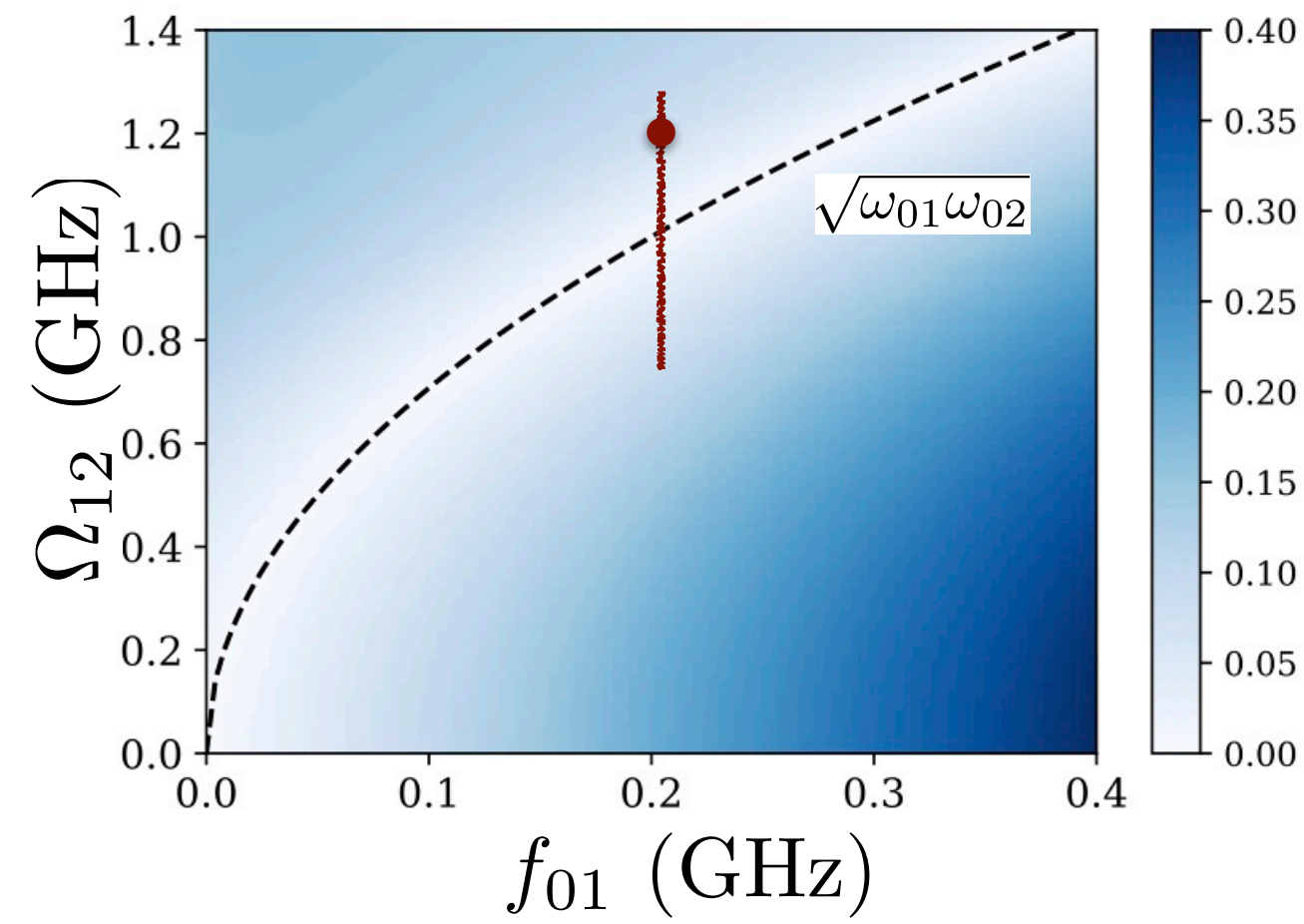


Topological Pump : effective 2 levels (qubit)

ω_{01}	$2\pi \times 200$ MHz
ω_{02}	$2\pi \times 5$ GHz
Ω_{01}	$2\pi \times 100$ MHz
Ω_{02}	$2\pi \times 1$ GHz
Ω_{12}	$2\pi \times 1.2$ GHz
$\dot{\phi}_a = 2\dot{\phi}_b$	$2\pi \times 20$ MHz



Topological Pump : effective 2 levels (qubit)

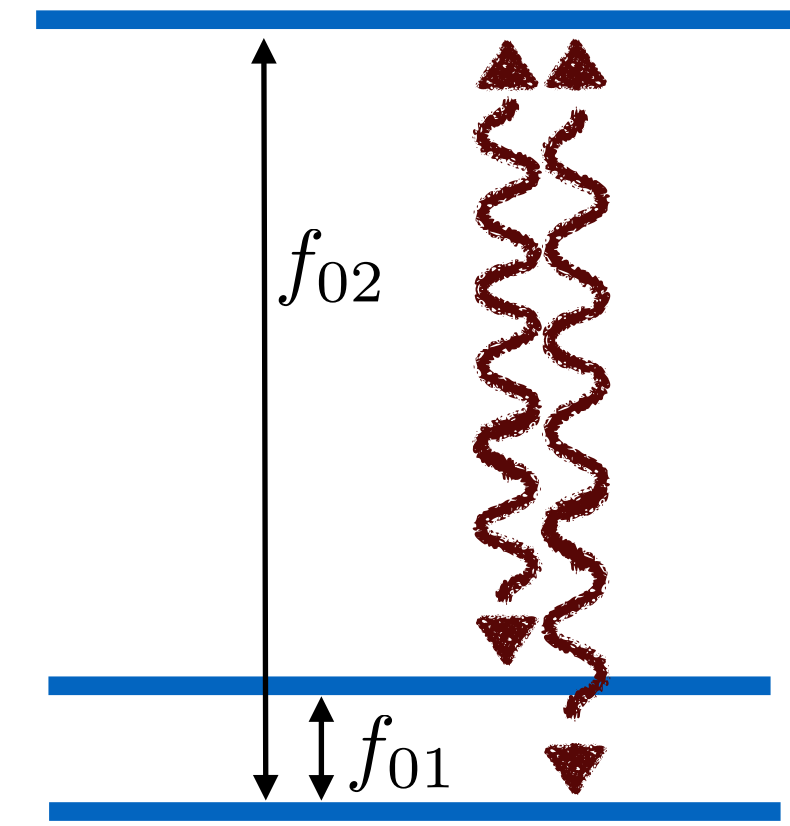


$$\omega_1 = 10 \text{ MHz}$$

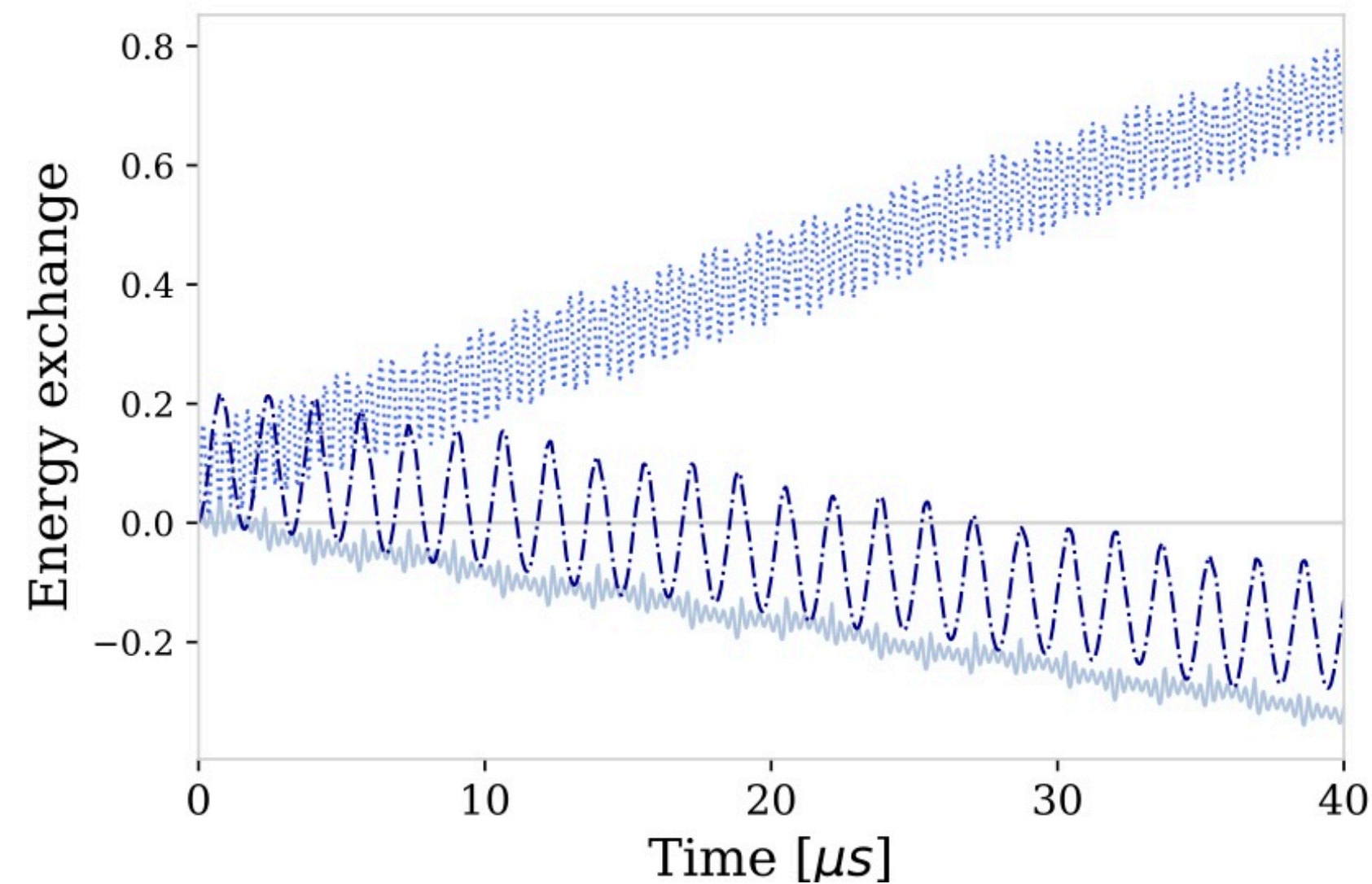
$$\omega_2 \simeq 8 \text{ MHz}$$

$$\Omega_{12} = 1.2$$

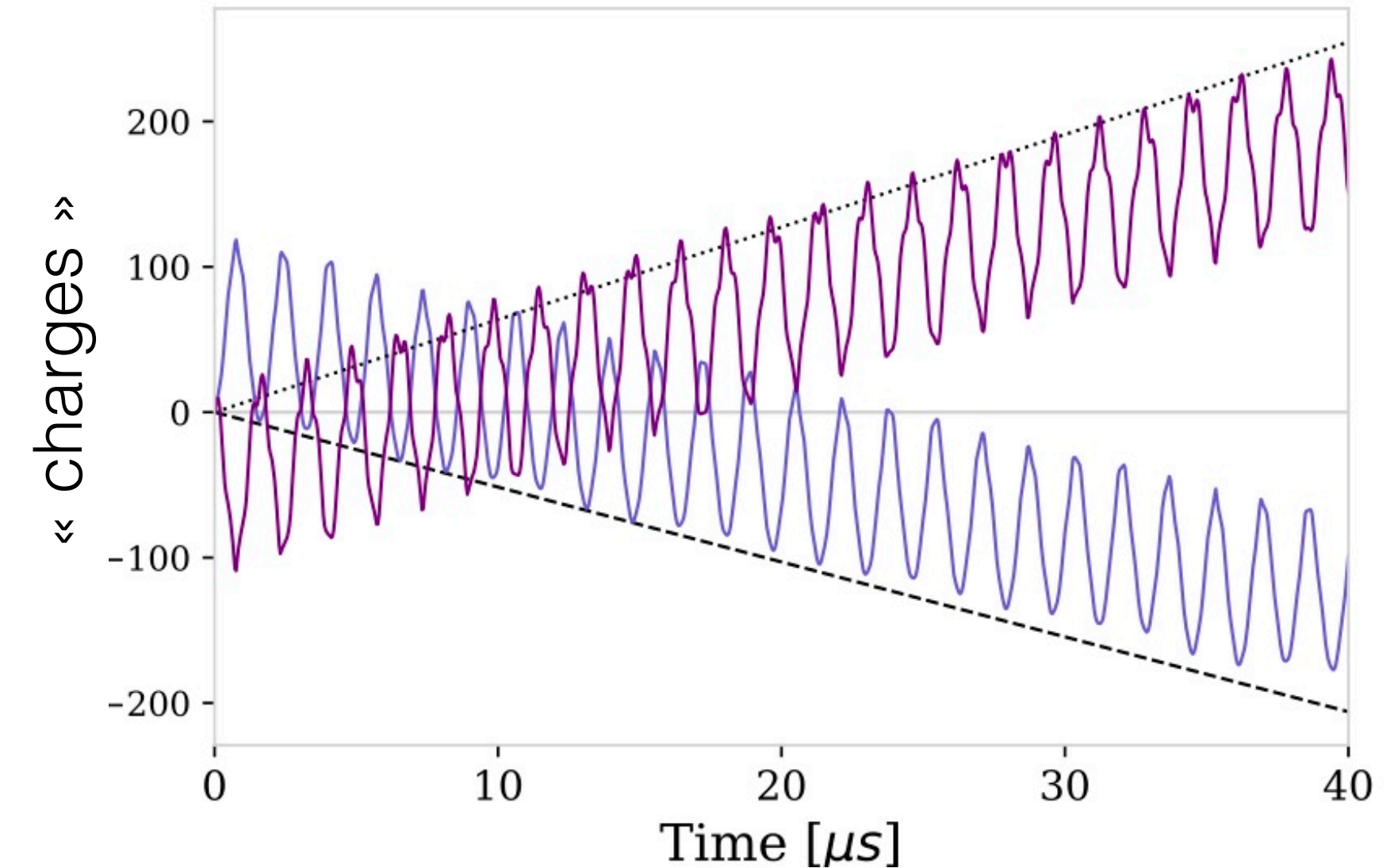
Prepare qutrit in ground state $|\psi_0\rangle$



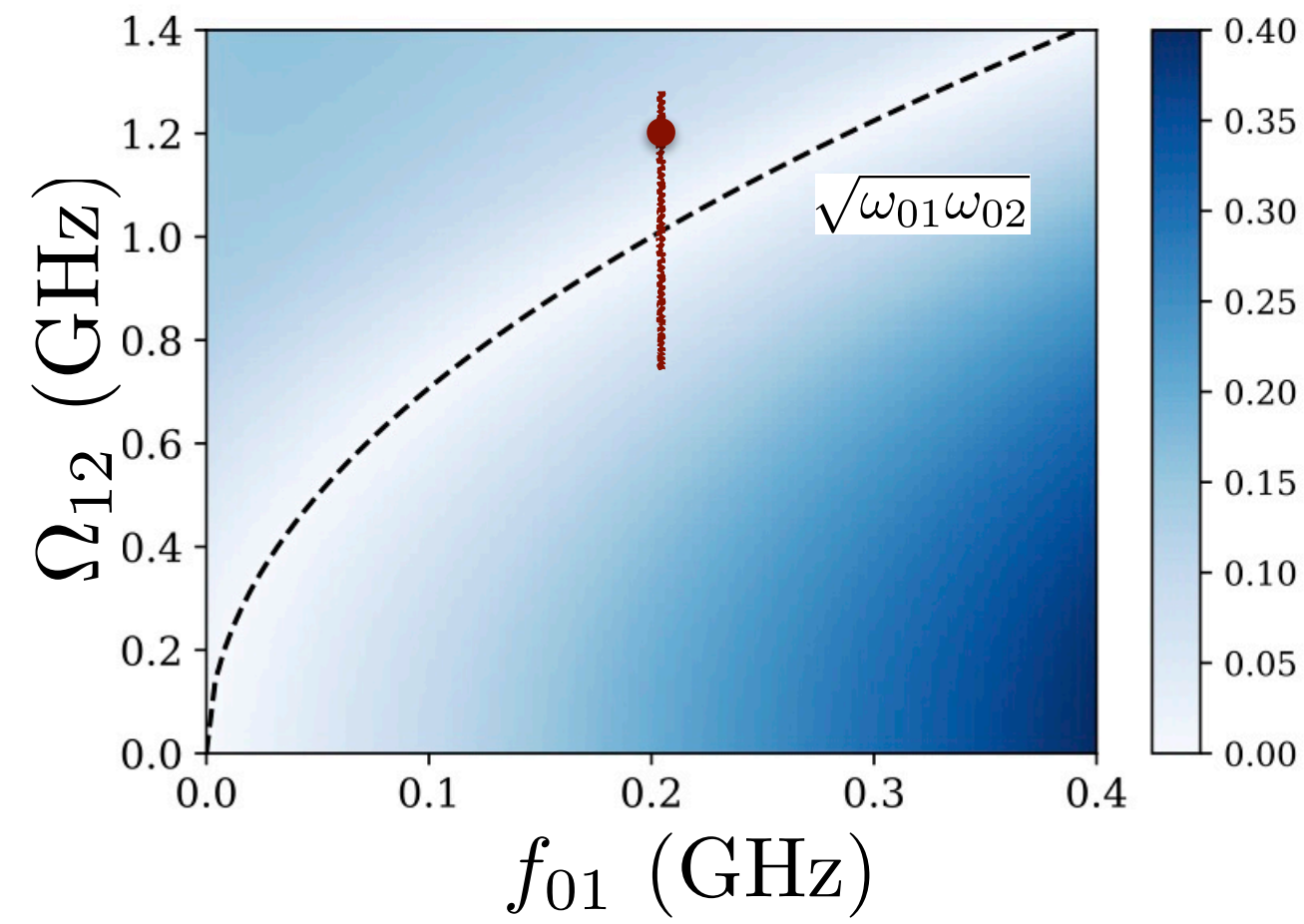
Energy exchange $\Delta\mathcal{E}_i$



« charges » $n_1 + n_3, n_2 - n_3$



Topological Pump : effective 2 levels (qubit)

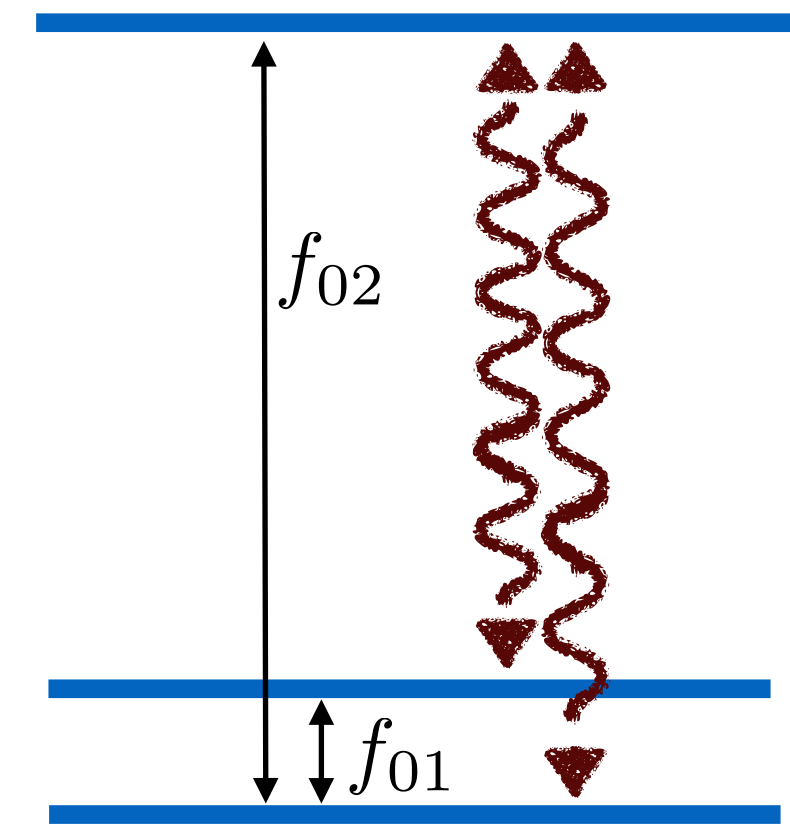


$$\omega_1 = 10 \text{ MHz}$$

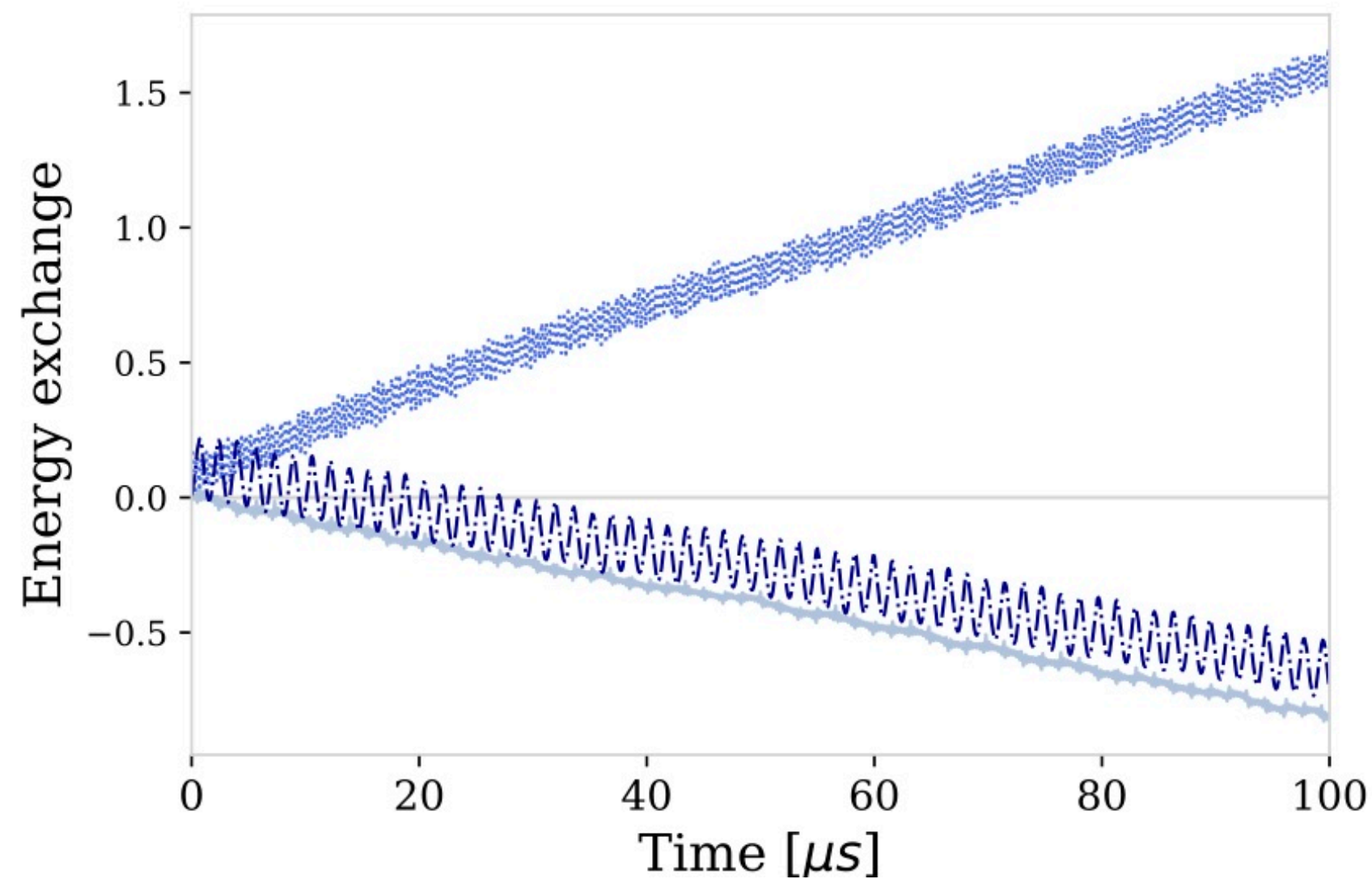
$$\omega_2 \simeq 8 \text{ MHz}$$

$$\Omega_{12} = 1.2$$

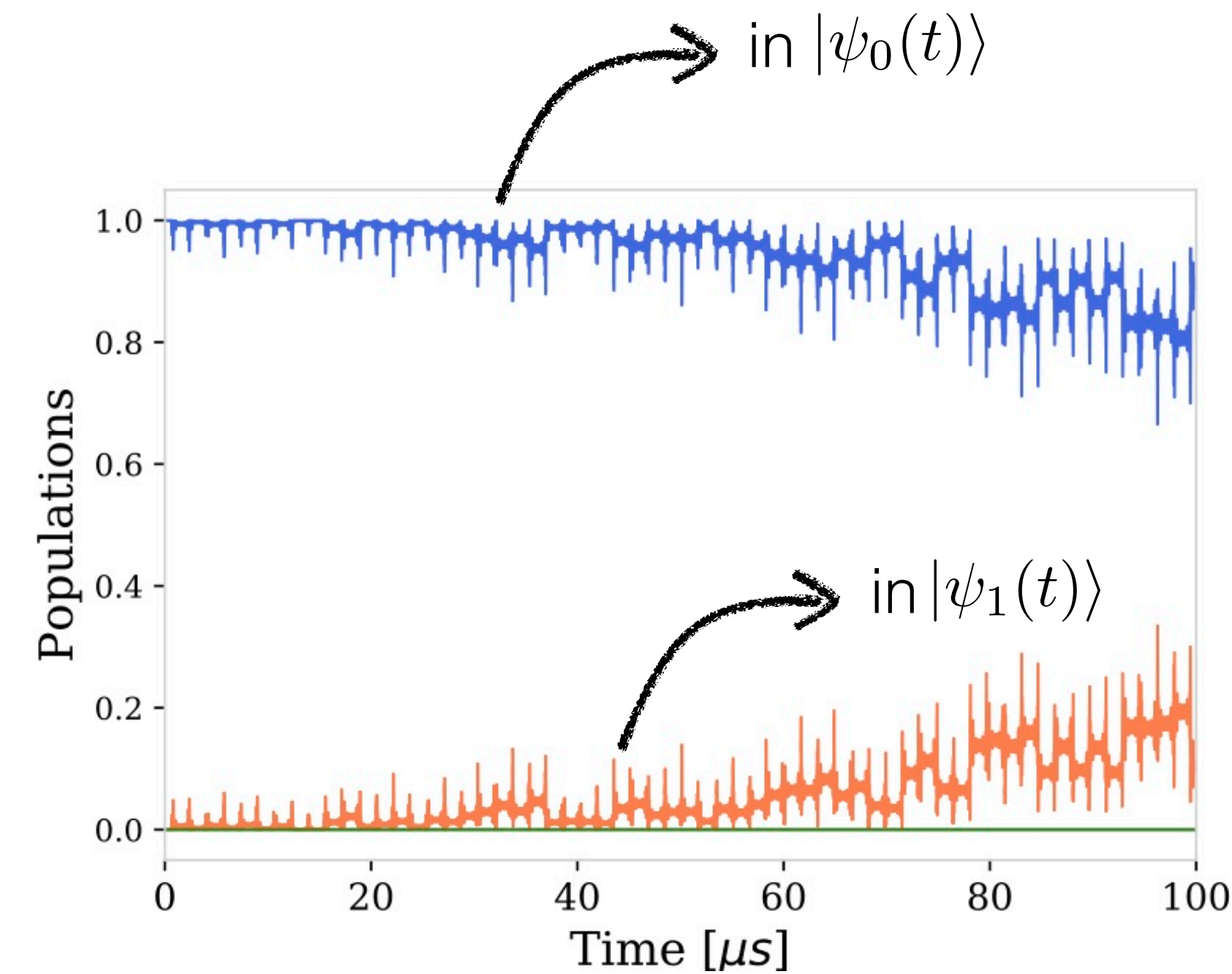
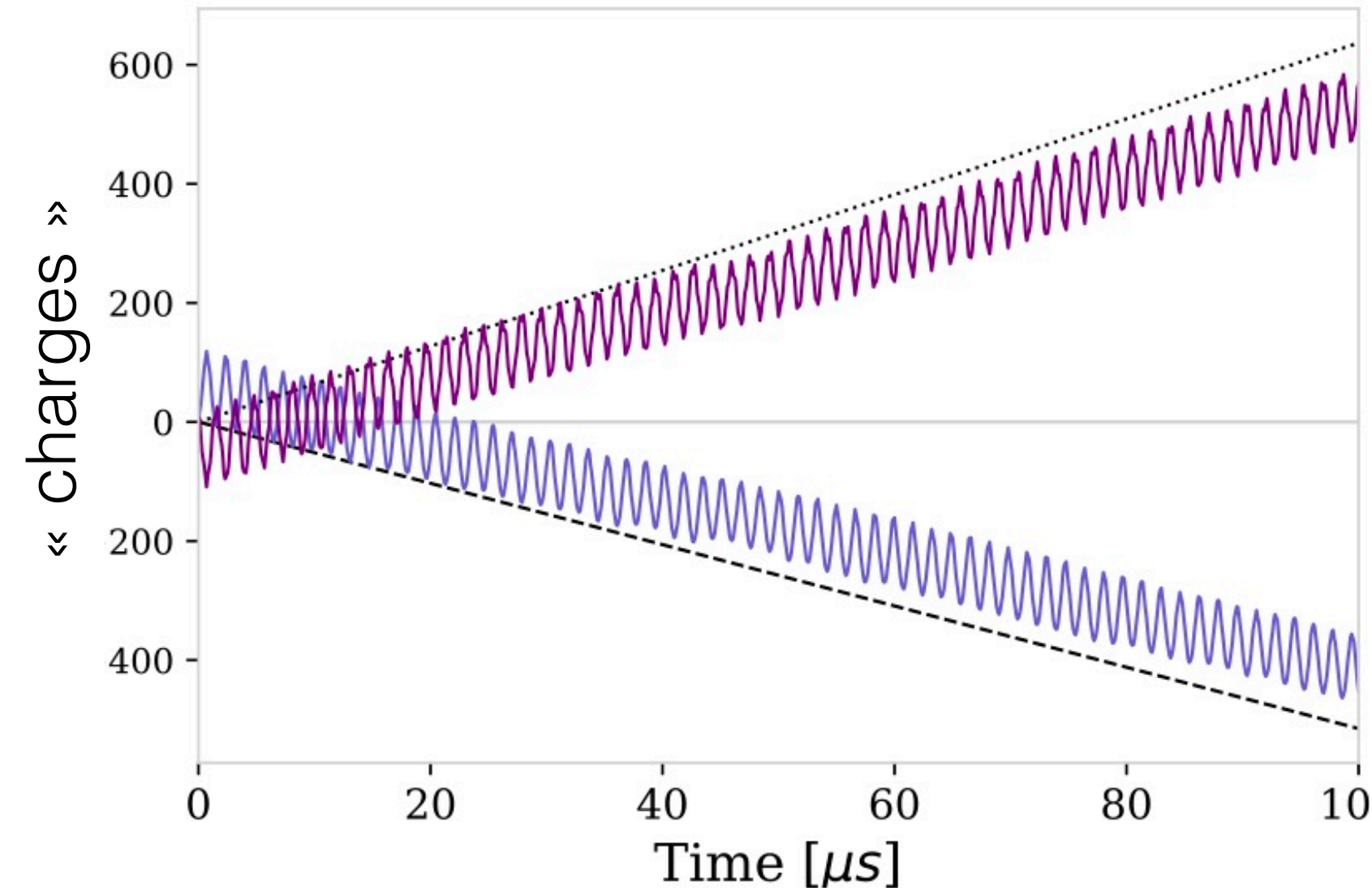
Prepare qutrit in ground state $|\psi_0\rangle$



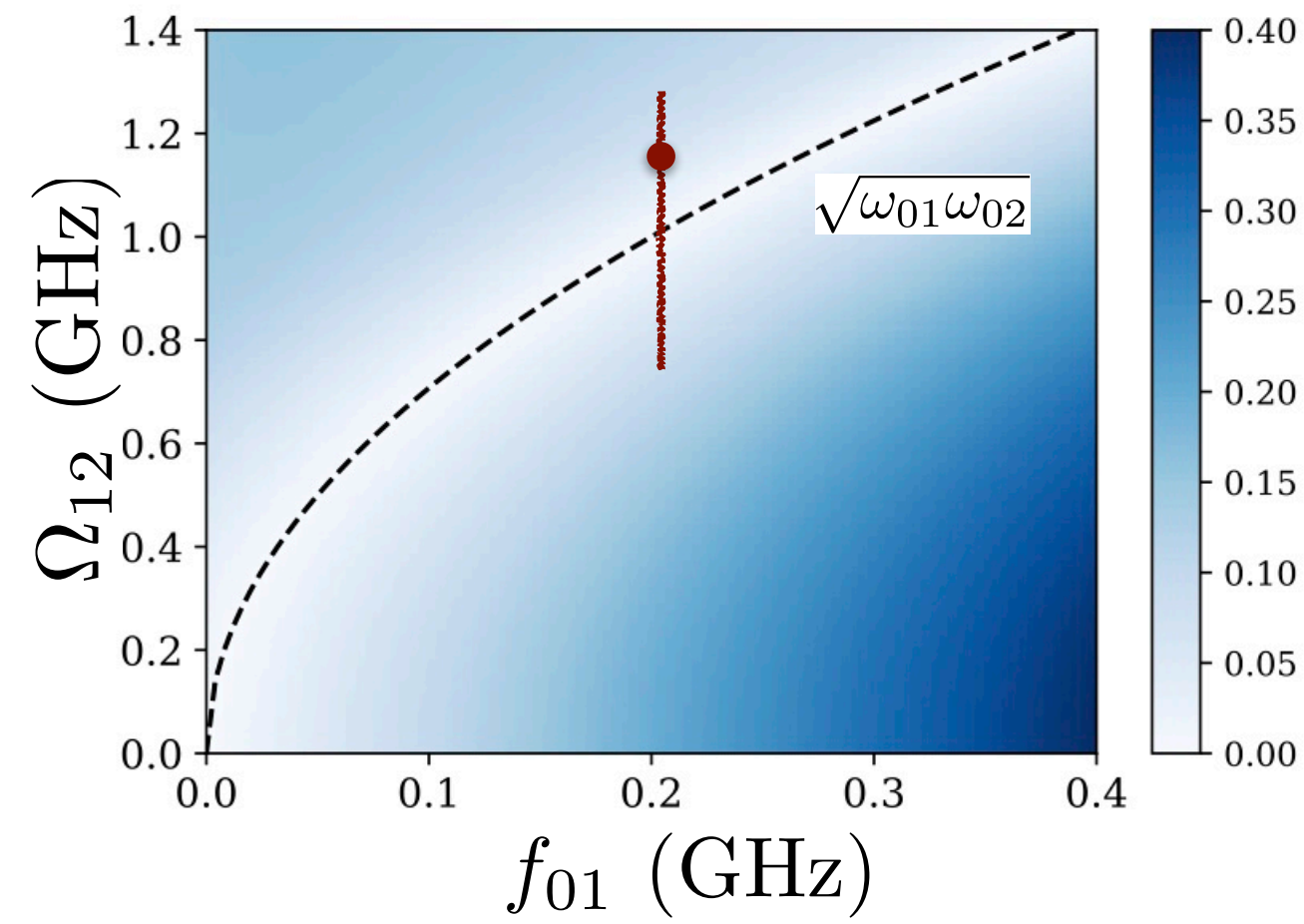
Energy exchange $\Delta\mathcal{E}_i$



« charges » $n_1 + n_3, n_2 - n_3$



Topological Pump : effective 2 levels (qubit)

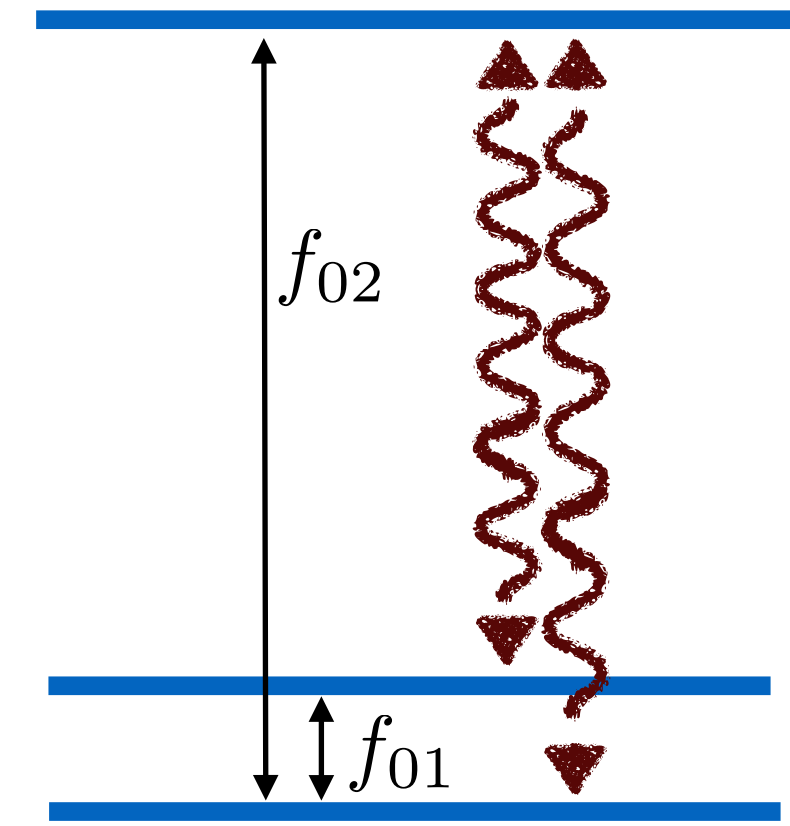


$$\omega_1 = 10 \text{ MHz}$$

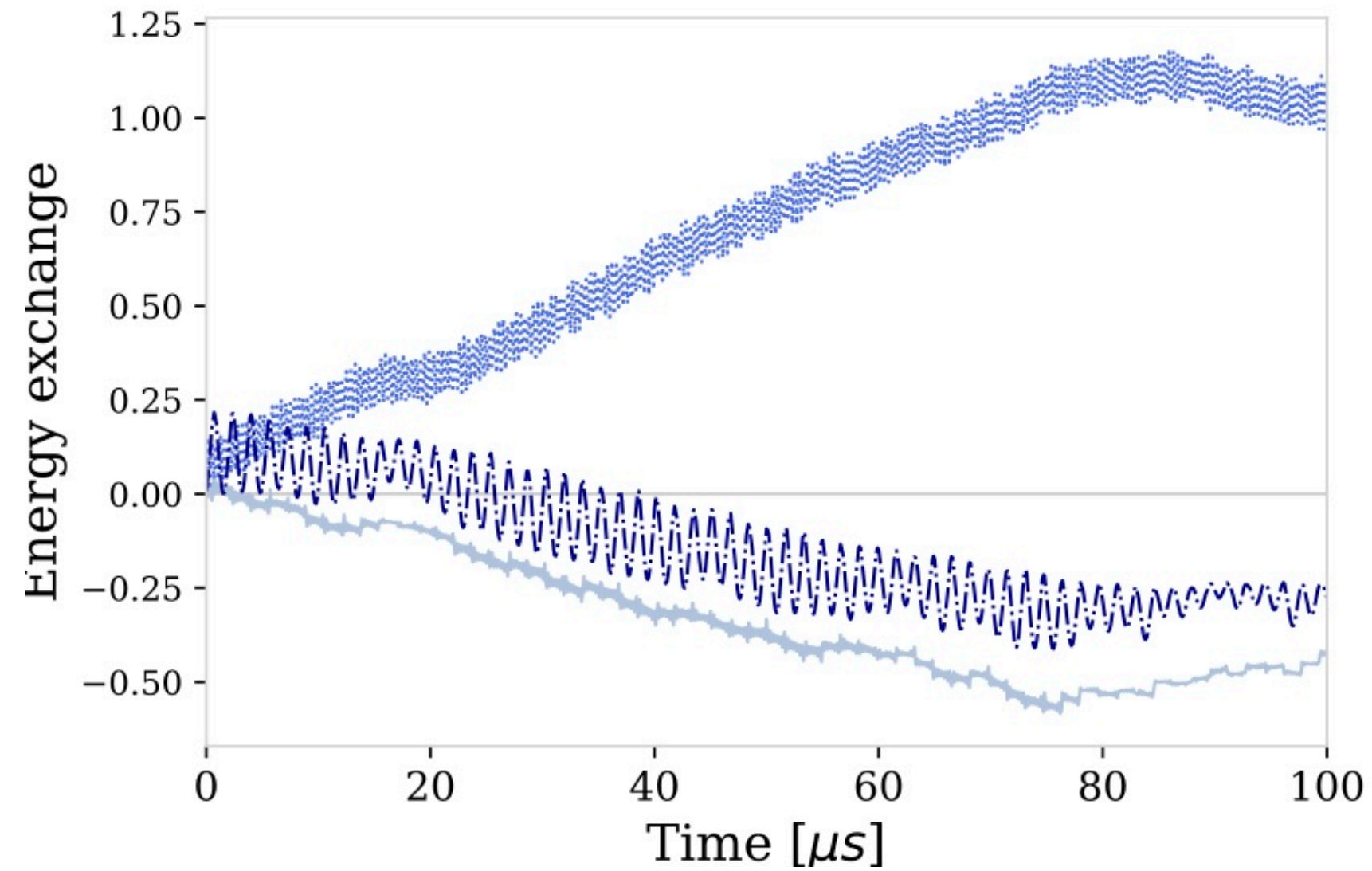
$$\omega_2 \simeq 8 \text{ MHz}$$

$$\Omega_{12} = 1.15$$

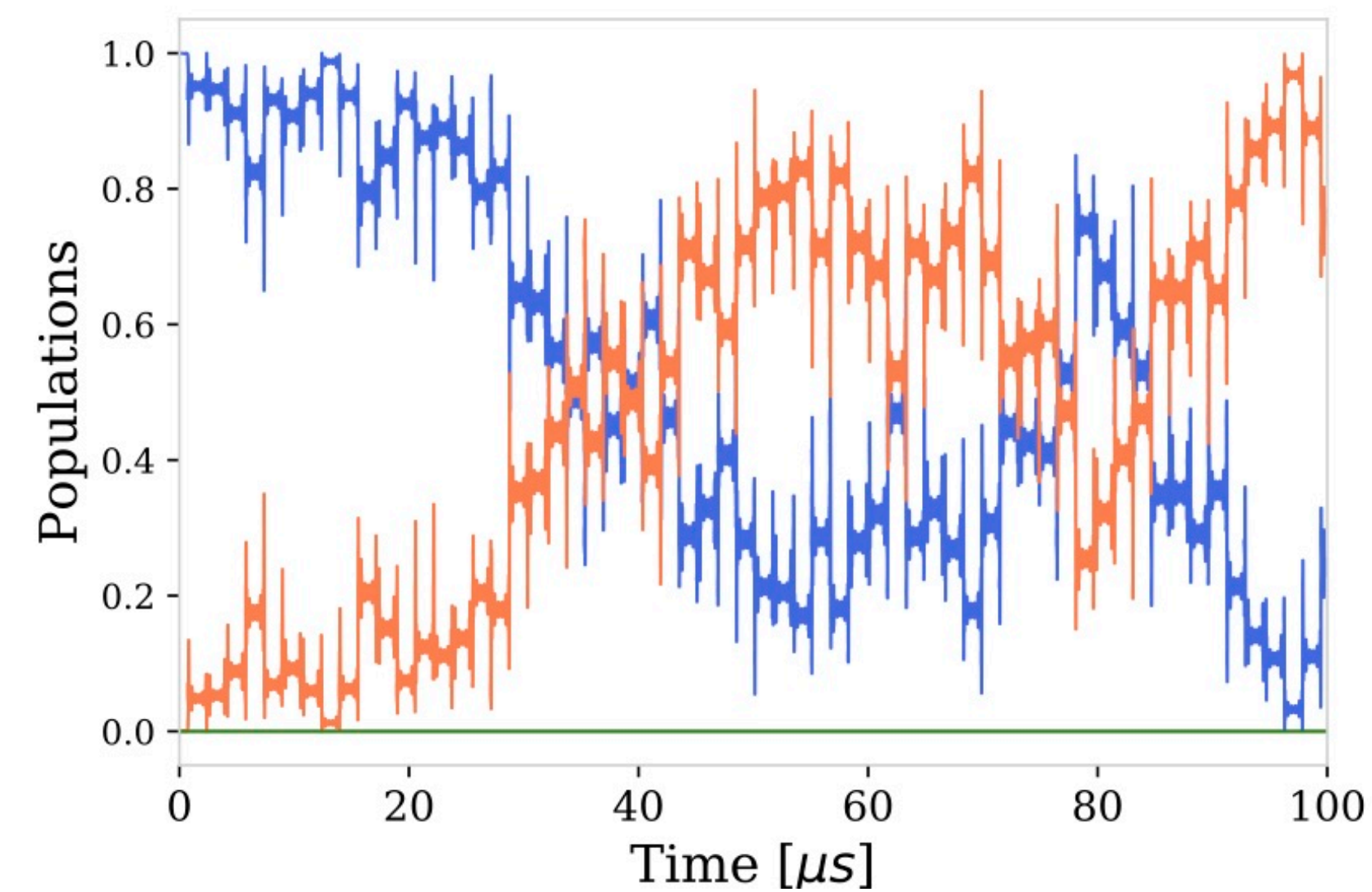
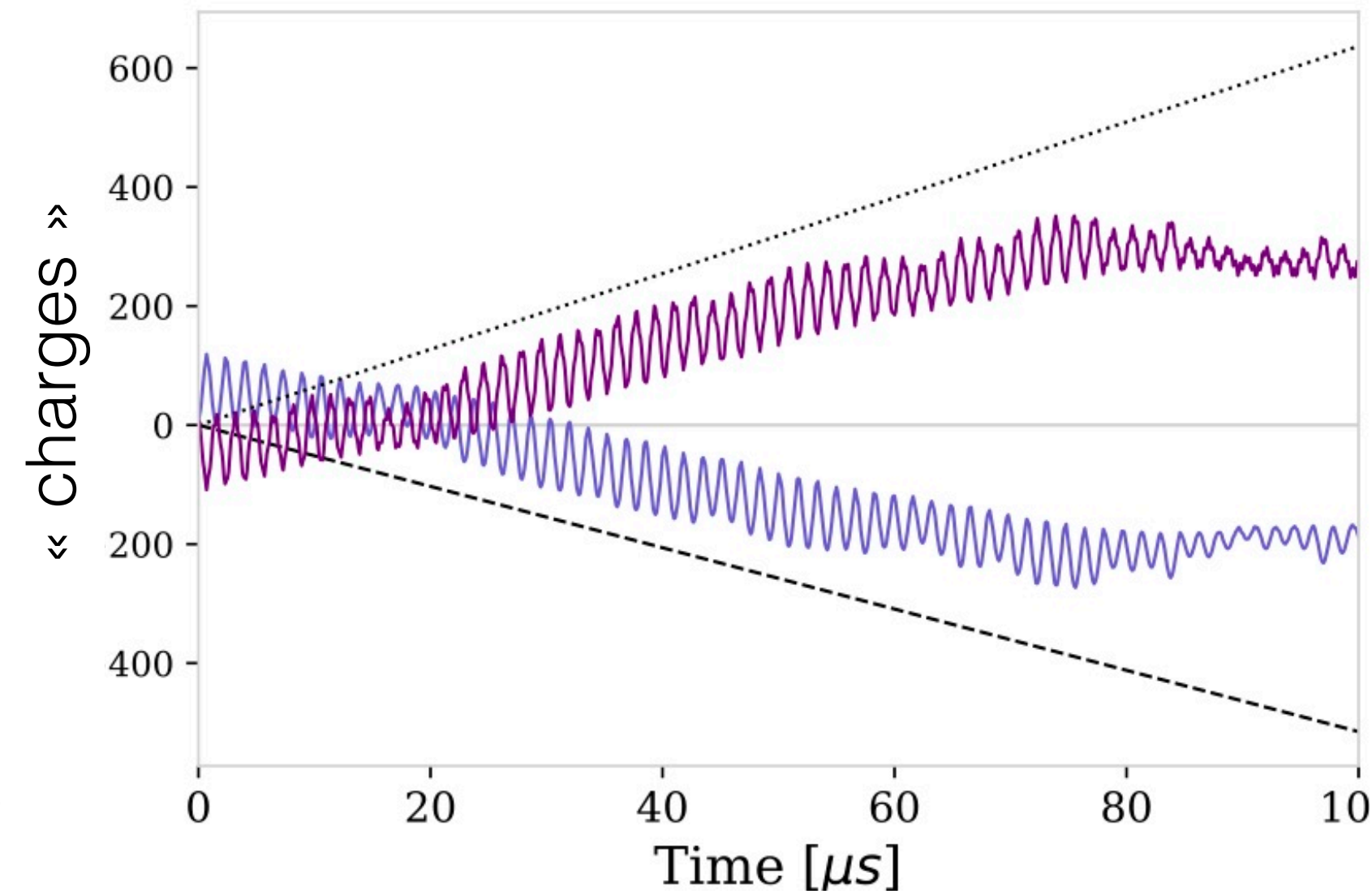
Prepare qutrit in ground state $|\psi_0\rangle$



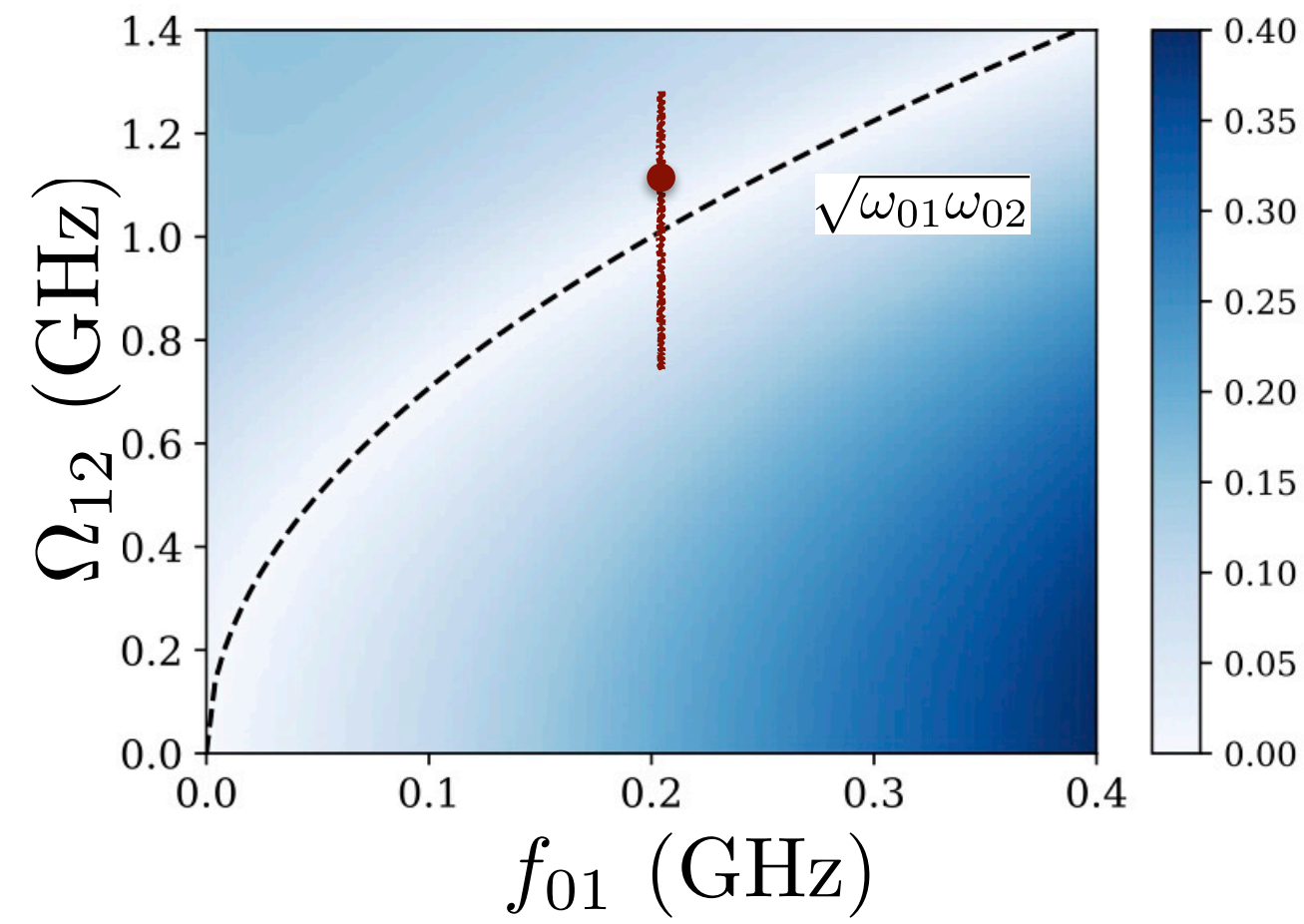
Energy exchange $\Delta\mathcal{E}_i$



« charges » $n_1 + n_3, n_2 - n_3$



Topological Pump : effective 2 levels (qubit)

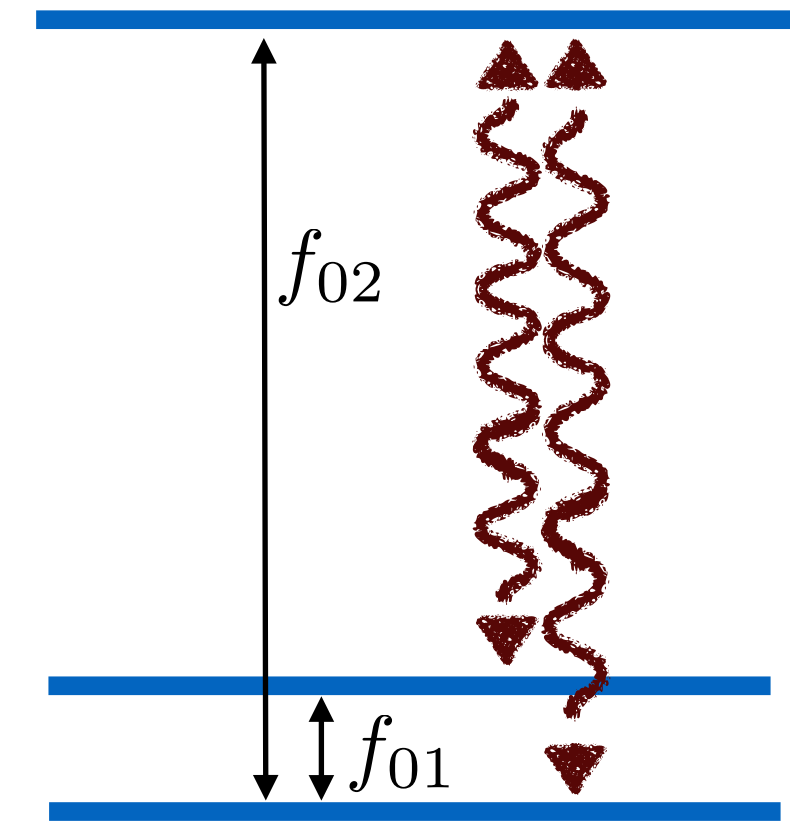


$$\omega_1 = 10 \text{ MHz}$$

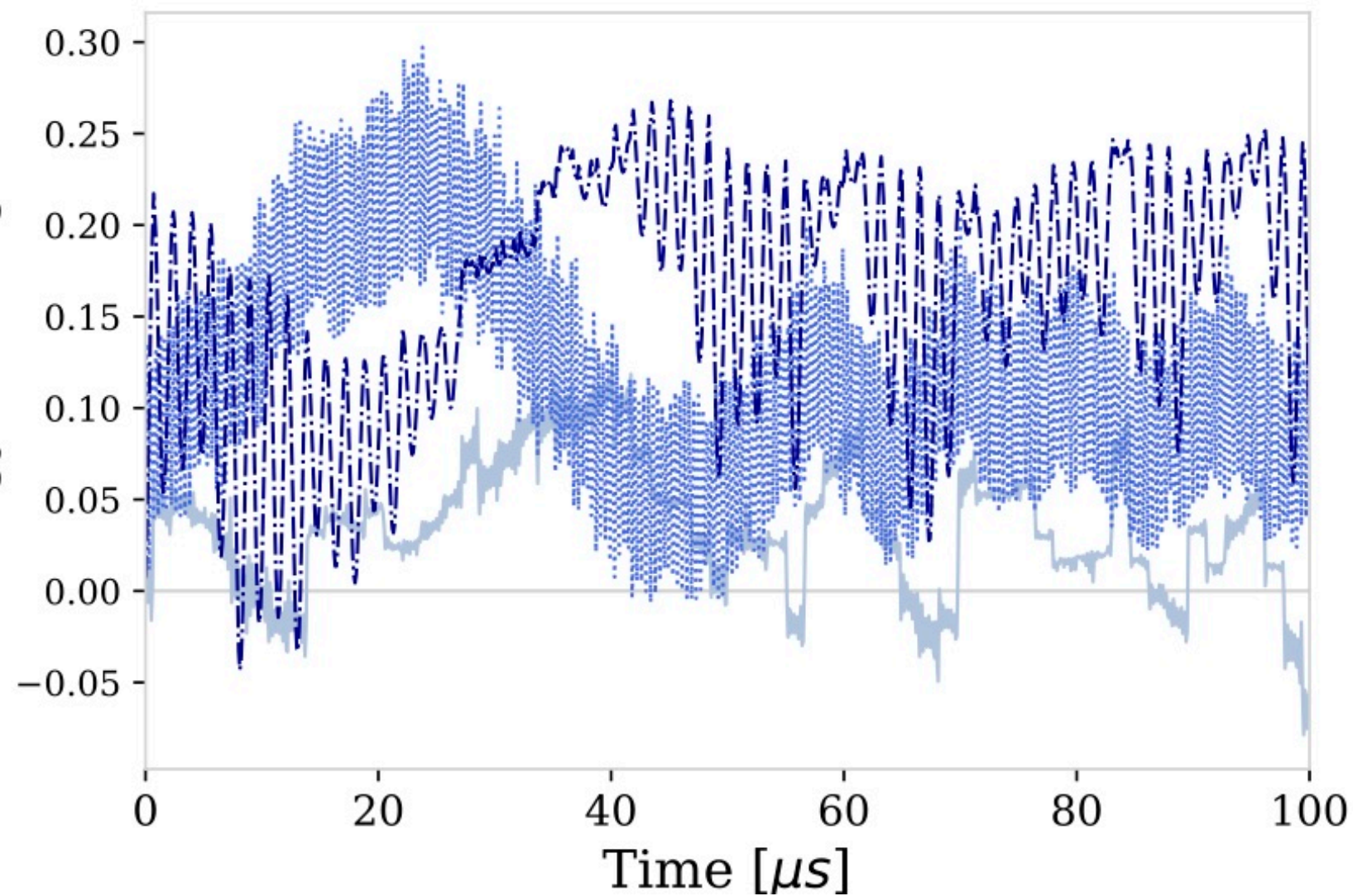
$$\omega_2 \simeq 8 \text{ MHz}$$

$$\Omega_{12} = 1.10$$

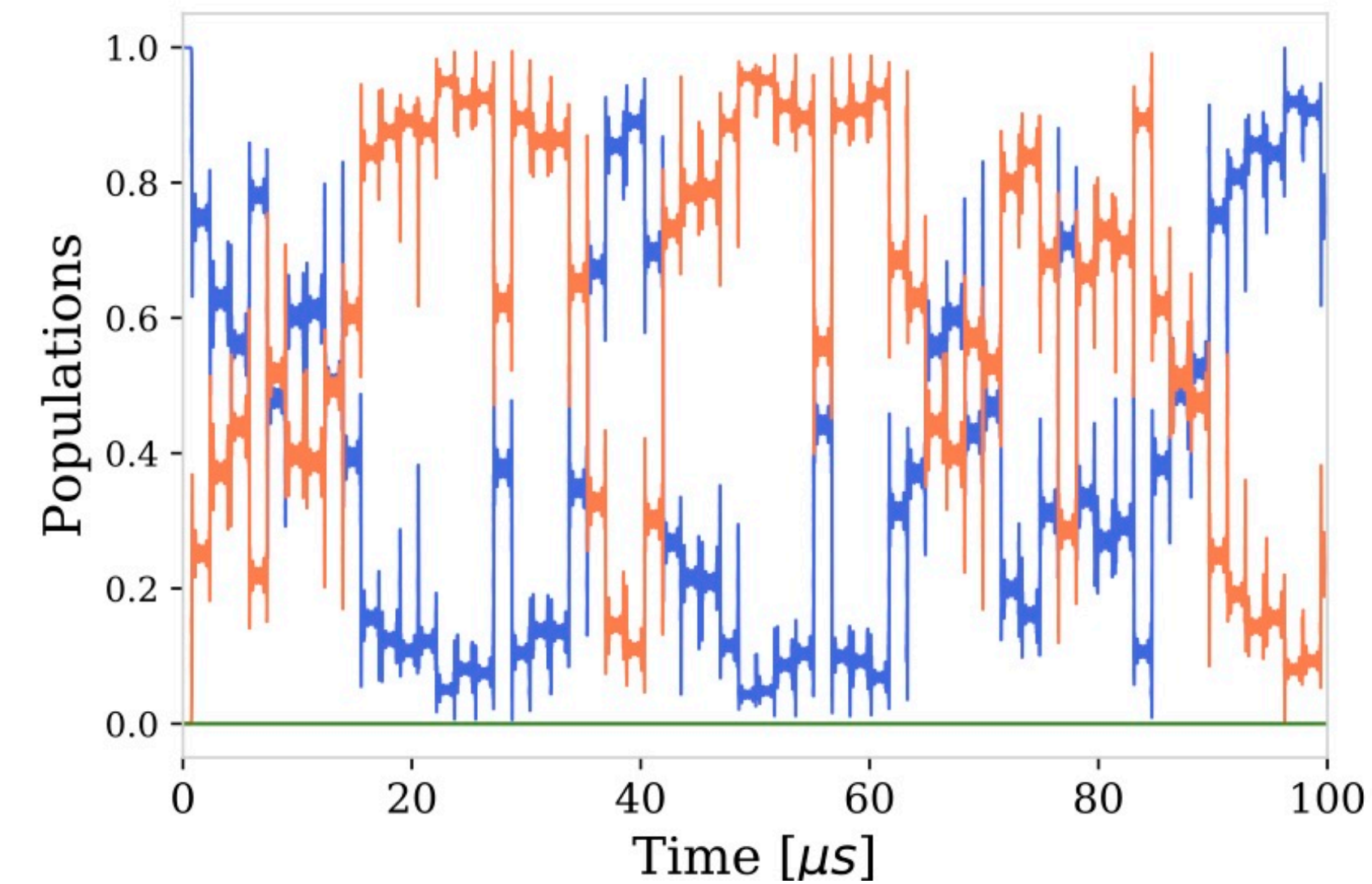
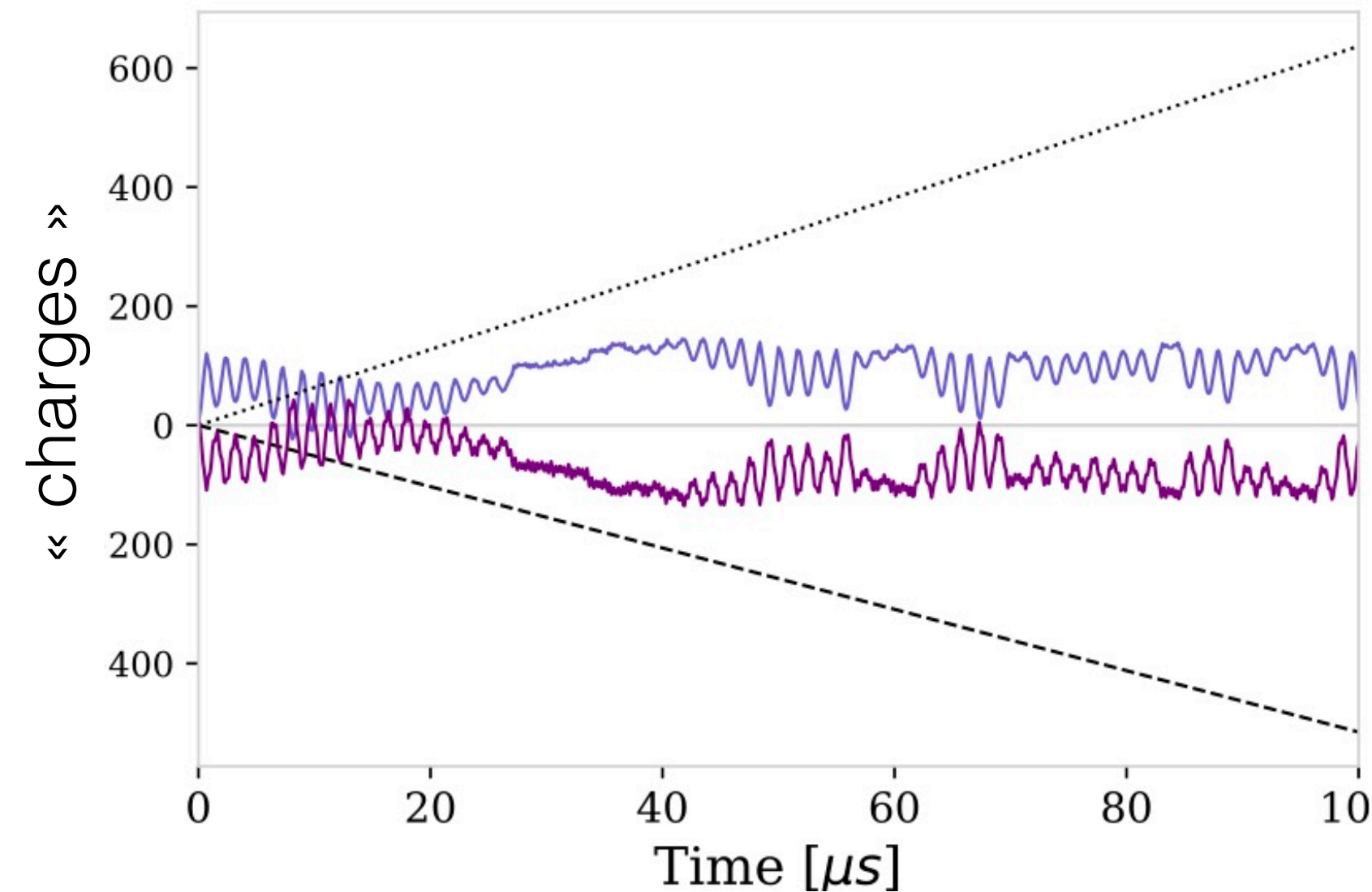
Prepare qutrit in ground state $|\psi_0\rangle$



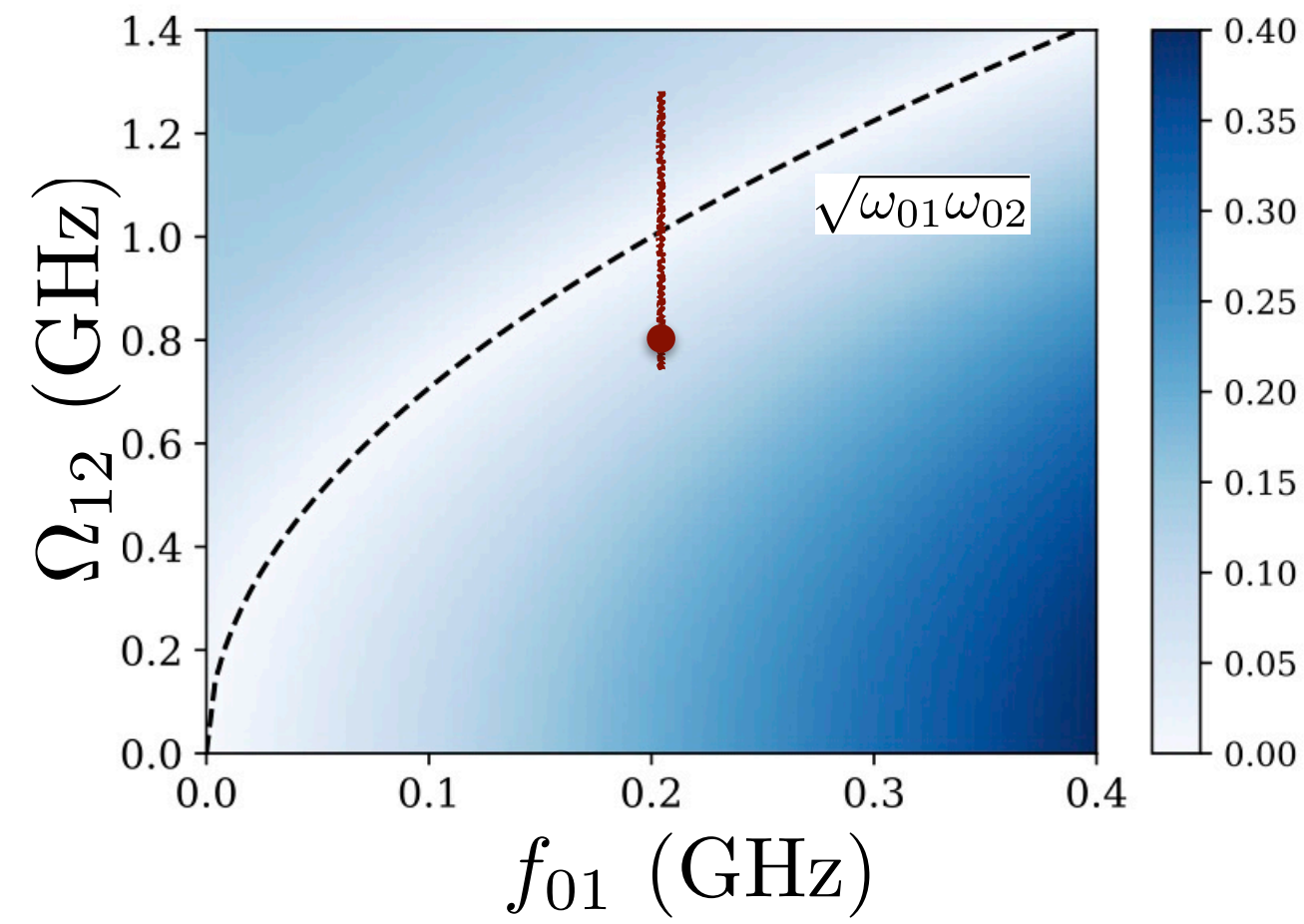
Energy exchange $\Delta\mathcal{E}_i$



« charges » $n_1 + n_3, n_2 - n_3$



Topological Pump : effective 2 levels (qubit)

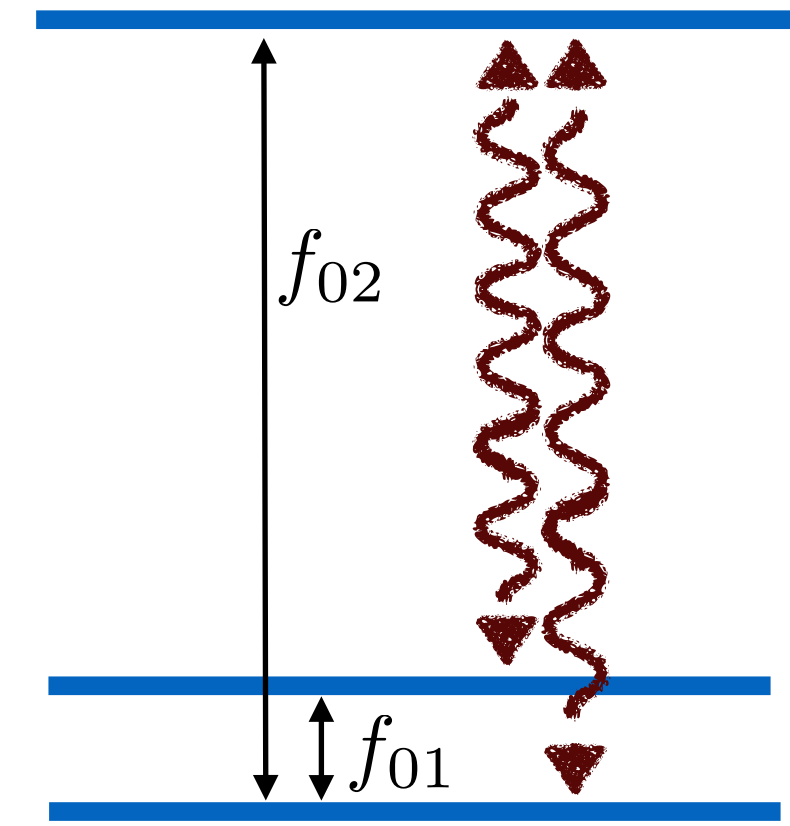


$$\omega_1 = 10 \text{ MHz}$$

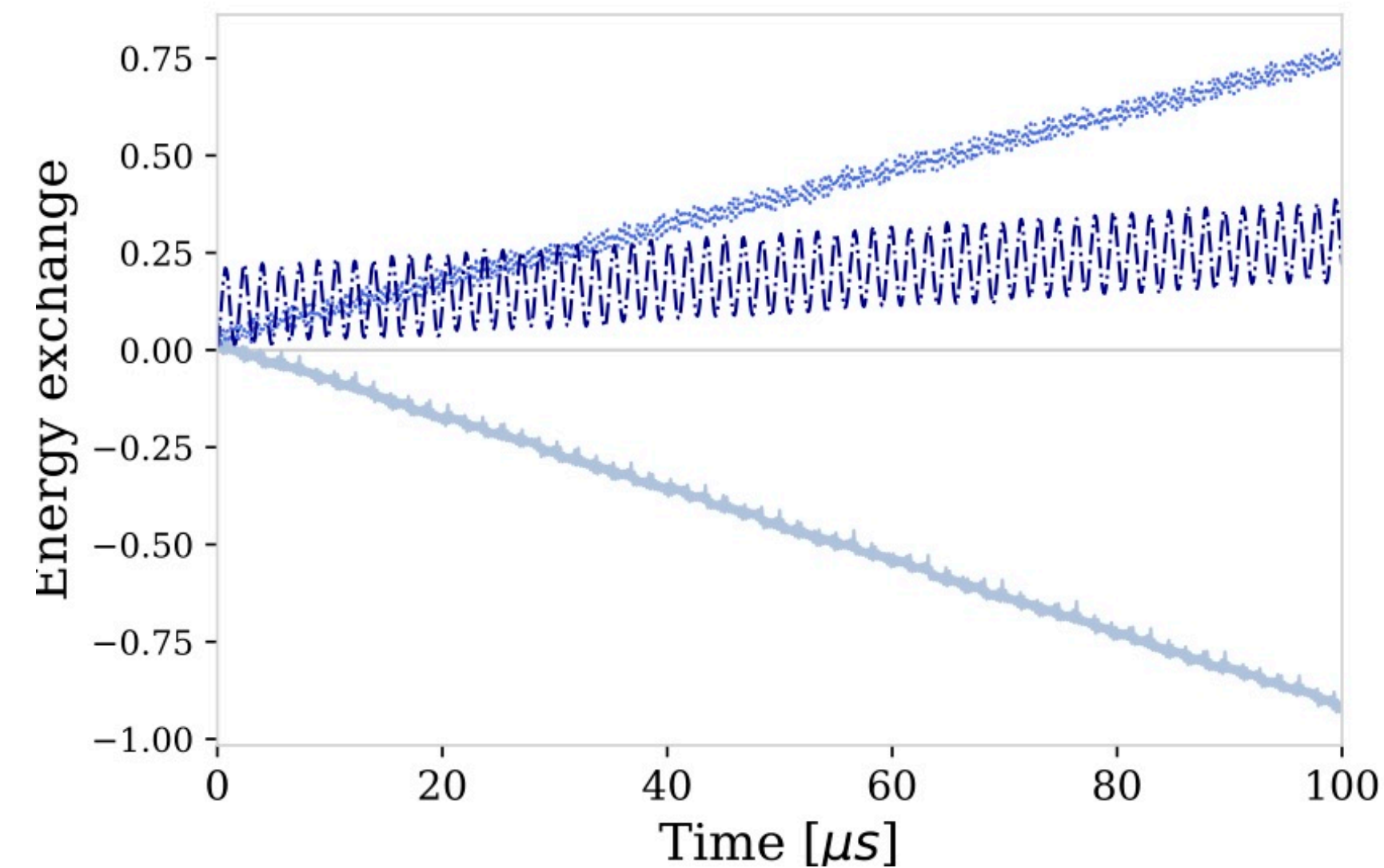
$$\omega_2 \simeq 8 \text{ MHz}$$

$$\Omega_{12} = 0.80$$

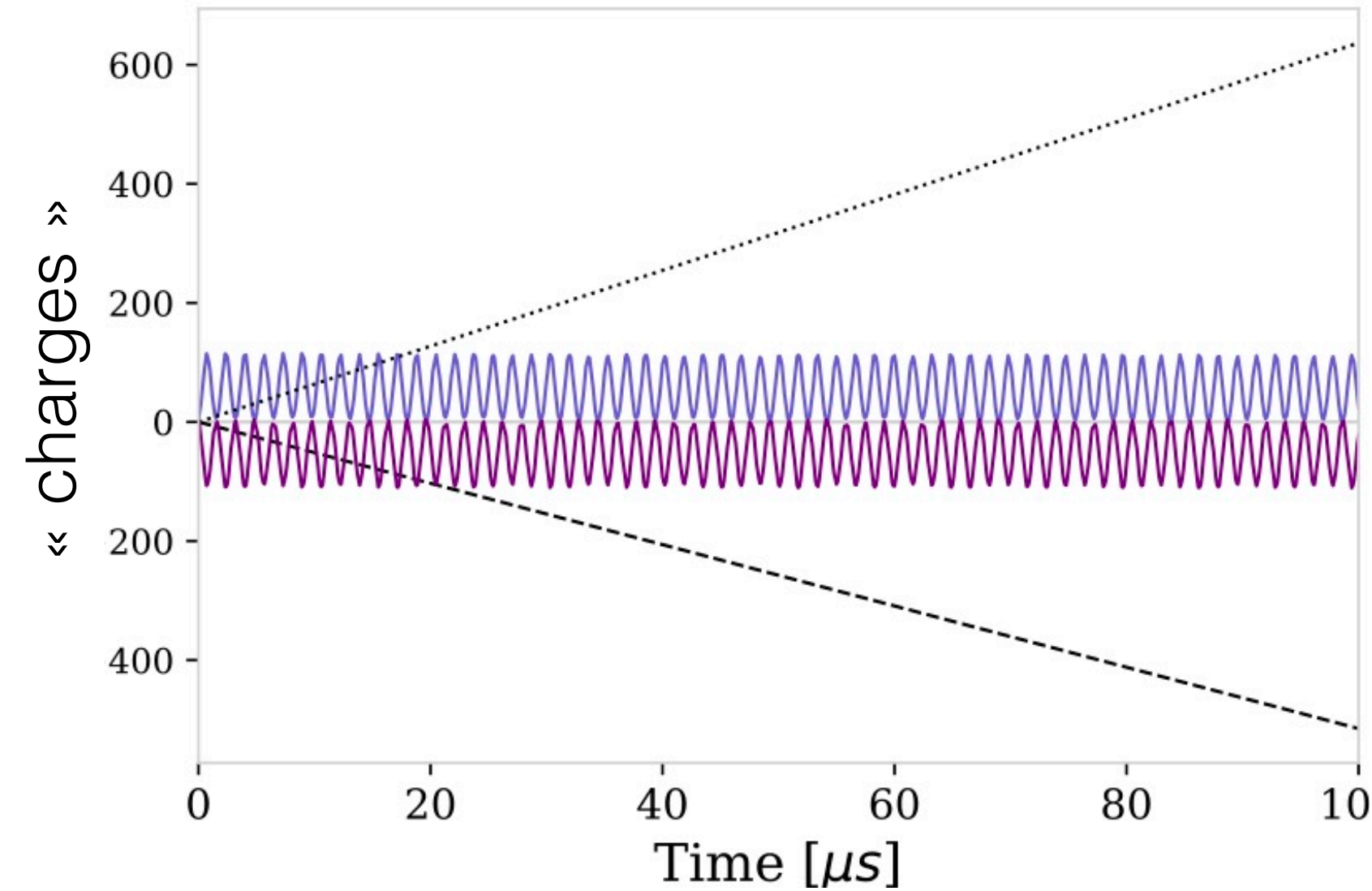
Prepare qutrit in ground state $|\psi_0\rangle$



Energy exchange $\Delta\mathcal{E}_i$



« charges » $n_1 + n_3, n_2 - n_3$



Populations

